

Component PERCUSSION

episode II

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1

 TI INFORMATION
Selective Disclosure

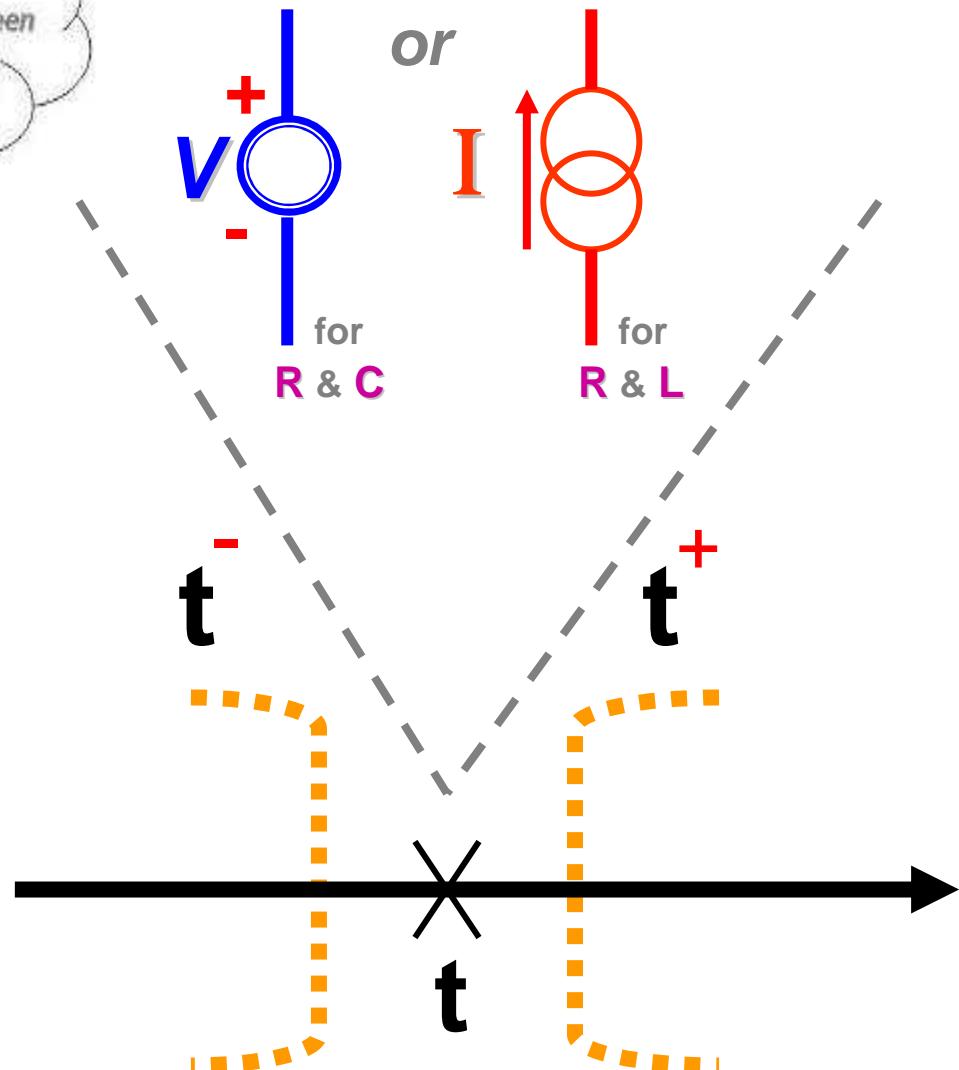




*In the previous
episode ...*



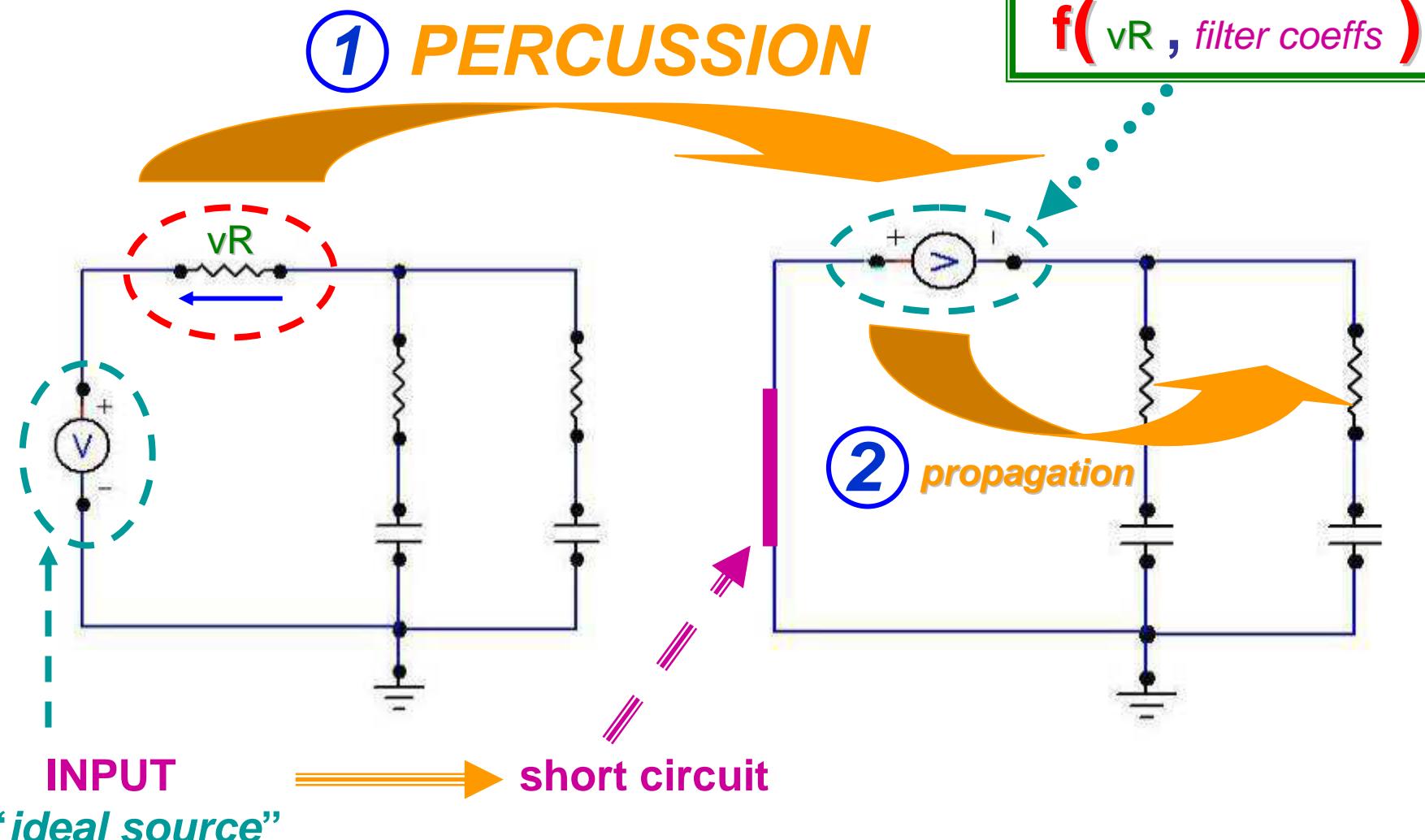
percuted component

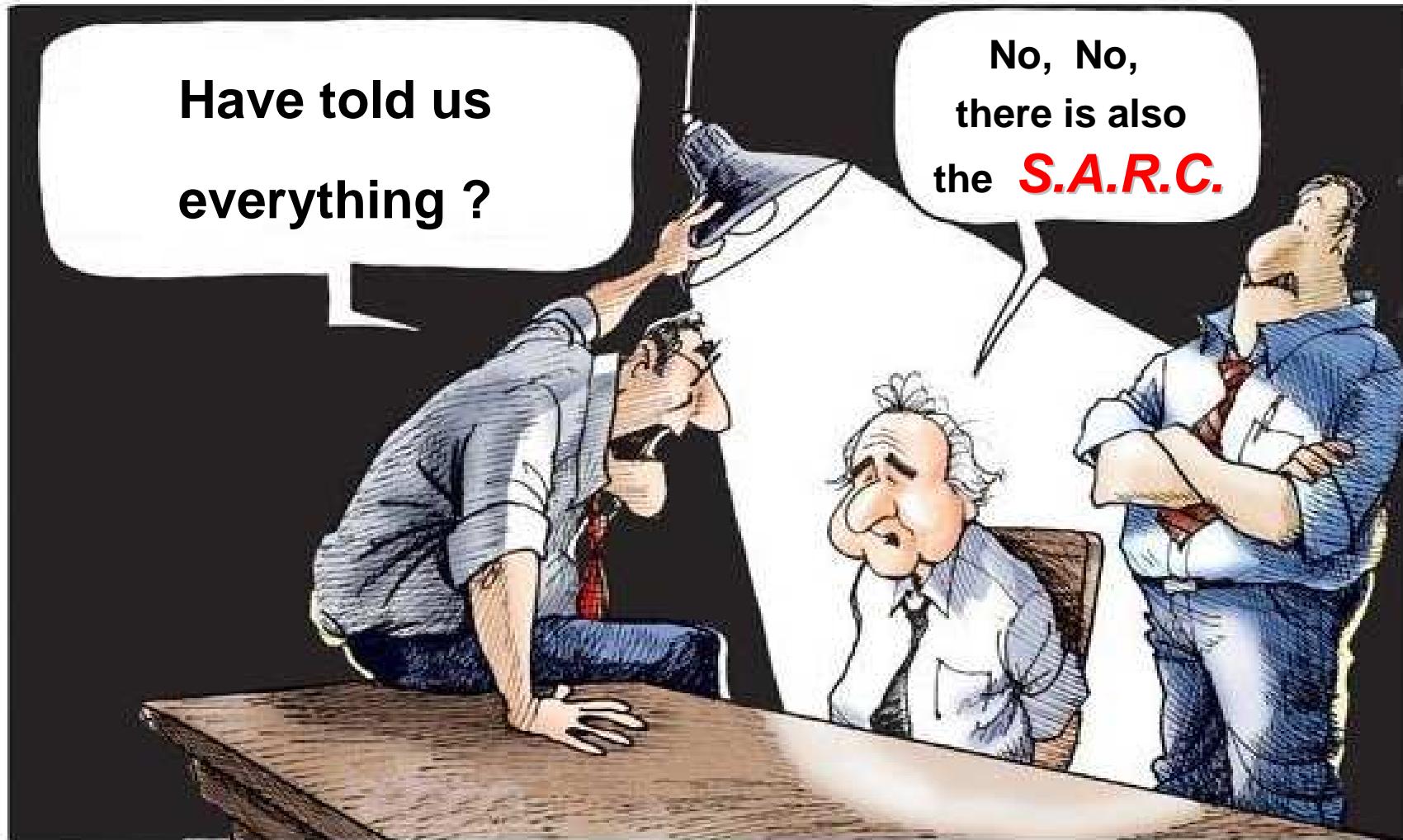


Differential equations
by means of **Distributions**

2 stages mystery . . .

"ideal source"
but for a null time !!!





Semi-analytical recursive algorithms for convolution calculations

W. Janke G. Blakiewicz

Department of Electronics, Technical University of Gdansk, Poland

IEE Proc. Circuits Devices Syst., Vol. 142 , No. 2, April 1995

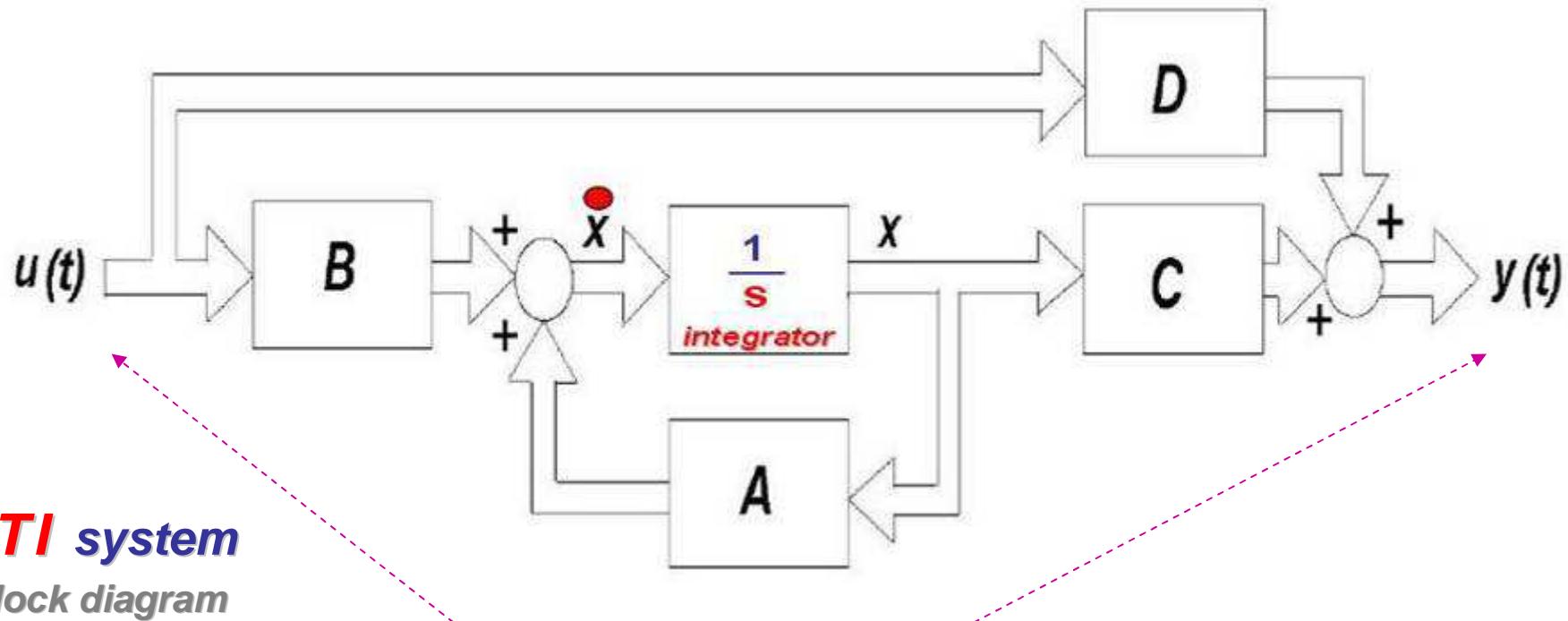
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STATE MODEL

MATRICES & VECTORS

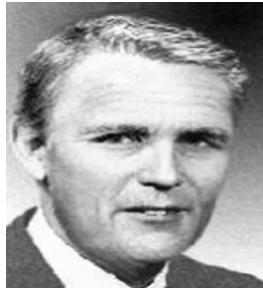
$$\dot{x}(t) = A x(t) + B u(t)$$

$$y(t) = C x(t) + D u(t)$$



LTI system
block diagram

multiple inputs / multiple outputs (MIMO)



Rudolf (Rudy) Emil Kálmán
born May 19, 1930
awarded in 2009 by
President Barack Obama with the
National Medal of Science

$$\dot{x}(t) = A x(t) + B u(t)$$
$$y(t) = C x(t) + D u(t)$$

*In the 60's, based on Kálmán research,
the State Model approach has been fully extended
and used, for examples, on the Apollo and Polaris projects*

For a given system, the **transfer function** output / input is unique
but multiple set of **A, B, C, D** matrices can be expressed
possibly corresponding to different **orders** (and thus different number of **poles** !!!)

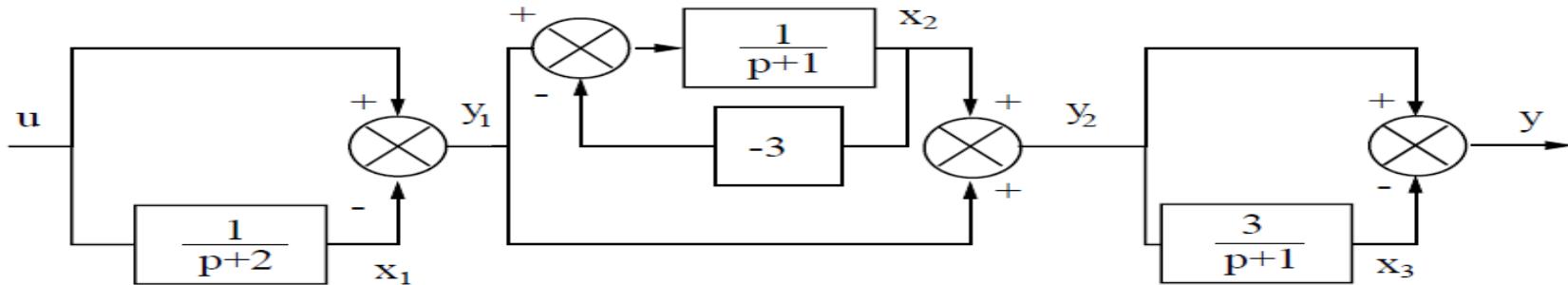
But, the **minimal order** only corresponds to the **A matrices** that enforce both
Controllable and **Observable** Kálmán criteria

It also exists, for singular systems, a “**generalized state representation**”

$$\dot{E} X = A X + B U$$
$$Y = C X + D U$$

E is called the “**derivative**” matrix

for **non-singular** systems **E = Identity**



system “*external*” VIEW

the ***transfer function*** expression is **unique**

1 pole: -2

$$\frac{Y(p)}{U(p)} = \frac{p - 1}{p + 2}$$

pole-zero cancellations
has occurred

The **output / input differential equation**

2 poles: -2 , -1

$$\ddot{y} + 3\dot{y} + 2y = \ddot{u} - u$$

system “*internal*” VIEW

state equation

3 poles: -2 , -1 , 2

$$\dot{x} = \begin{bmatrix} -2 & 0 & 0 \\ -1 & 2 & 0 \\ -3 & 3 & -1 \end{bmatrix} x + \begin{bmatrix} 1 \\ 1 \\ 3 \end{bmatrix} u$$

matrix A

1 pole: -2

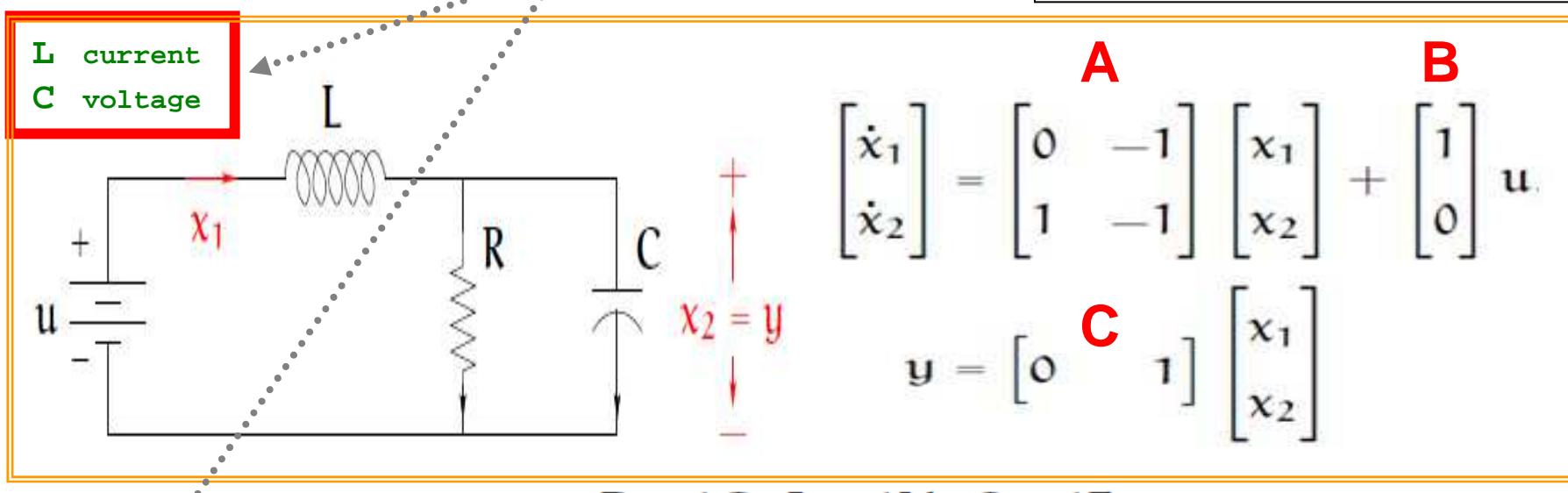
$$\dot{x} = [-2] x + [1] u$$

one of the possible “*not minimal*” state equation

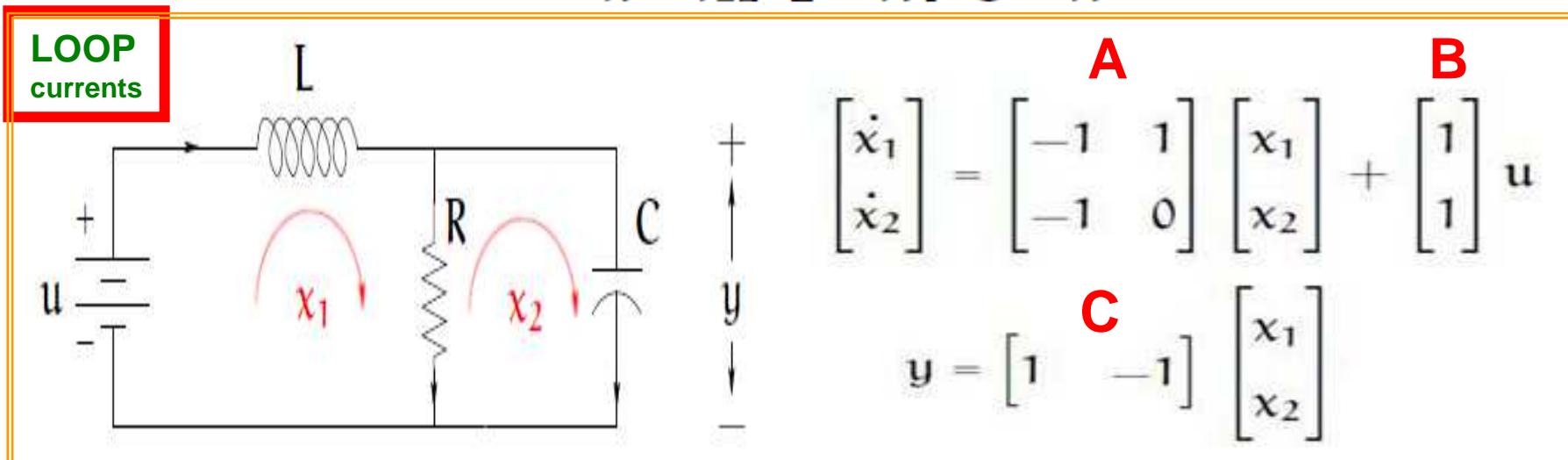
“*minimal*” state equation

Equivalent STATE equations

STATEs can be:
 L current, C voltage, LOOP current,
 NODE voltage, L flux, C charge, etc ...



$$R = 1\Omega \quad L = 1H \quad C = 1F$$



STATE

equivalence transformation

$$z = T x$$

$$x = T^{-1} z$$

Algebraic Equivalence

$$\dot{x} = A x + B u$$

$$y = C x + D u$$

?

$$\dot{z} = \tilde{A} z + \tilde{B} u$$

$$y = \tilde{C} z + \tilde{D} u$$

$$\tilde{A} = T A T^{-1}$$

$$\tilde{B} = T B$$

$$\tilde{C} = C T^{-1}$$

$$\tilde{D} = D$$

STATE MODEL

STATE equation →

OUTPUT equation →

$$\dot{x}(t) = Ax(t) + Bu(t)$$

$$y(t) = Cx(t) + Du(t)$$

WITH

vector of **signals**
having **derivative**
due to **inductance**
or **capacitance**

A
B
C
D

X
U
Y

(or feedthrough , feedforward)

state MATRIX
input MATRIX
output MATRIX
direct transmission MATRIX
state VECTOR
input (or **control**) VECTOR
output VECTOR

This formalism really highlights the *Linear Time Invariant* terminology

It is mostly used by the “automatism” world (“control command”)

$$\dot{\mathbf{x}}(t) = A \mathbf{x}(t) + B \mathbf{u}(t)$$

$$\mathbf{y}(t) = C \mathbf{x}(t) + D \mathbf{u}(t)$$

DIMENSIONS

| Variable | Dimension | Name |
|--------------|---------------------|----------------------------|
| s | | Number of states |
| i | | Number of inputs |
| o | | Number of outputs |
| A | $s \times s$ matrix | State matrix |
| B | $s \times i$ matrix | Input matrix |
| C | $o \times s$ matrix | Output matrix |
| D | $o \times i$ matrix | Direct transmission matrix |
| \mathbf{x} | s vector | State vector |
| \mathbf{u} | i vector | Input vector |
| \mathbf{y} | o vector | Output vector |

MATRICES

VECTORS

STATE equation

$$\dot{\mathbf{X}} = \mathbf{A} \mathbf{X} + \mathbf{B} \mathbf{U}$$

not easy to compute (searching \mathbf{X} involves a “*convolution*”)

order of matrix \mathbf{A} is **nearly** related to the number of *Inductances* and *Capacitances*

When *order* \neq *number energy-storing elements* the network is called “*degenerate network*”

The *order* of complexity of any network equals the total number of *energy-storing elements* minus the *number of all-capacitor “loops”* and the *number of all-inductor “cut-sets”*

OUTPUT equation

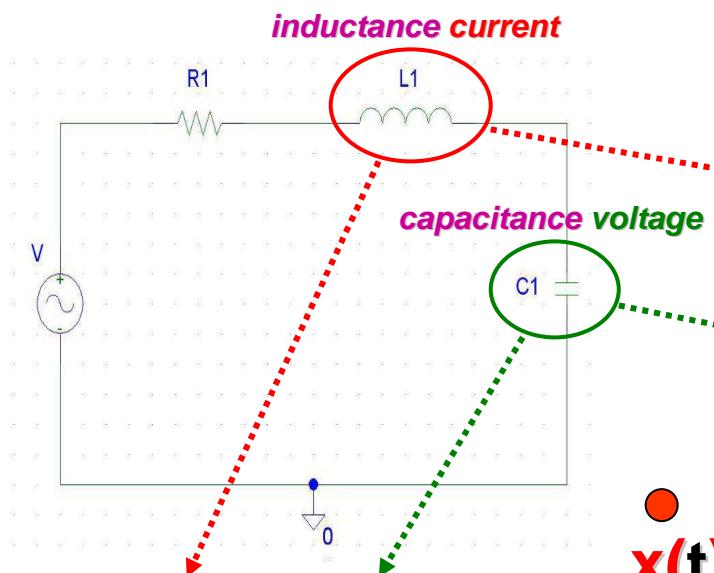
$$\mathbf{Y} = \mathbf{C} \mathbf{X} + \mathbf{D} \mathbf{U}$$

trivial to compute as soon as \mathbf{X} is known

able to return any *voltage* or *current* node similarly to *Kirchhoff / Ohm*

order of vector \mathbf{Y} depends on the number of *expected observed values*

STATE MODEL



$$\dot{x}(t) = Ax(t) + Bu(t)$$

$$y(t) = Cx(t) + Du(t)$$

$$L1 \frac{d iL1(t)}{dt} = -R1 iL1(t) - vC1(t) + u(t)$$

$$C1 \frac{d vC1(t)}{dt} = iL1(t)$$

$x(t)$

A

$x(t)$

B

2 impedance components

state dimension 2

$$\begin{pmatrix} \frac{d iL1(t)}{dt} \\ \frac{d vC1(t)}{dt} \end{pmatrix} = \begin{pmatrix} -\frac{R1}{L1} & \frac{1}{L1} \\ \frac{1}{C1} & 0 \end{pmatrix} \begin{pmatrix} iL1(t) \\ vC1(t) \end{pmatrix} + \begin{pmatrix} \frac{1}{L1} \\ 0 \end{pmatrix} u(t)$$

Now, let's suppose that
for the $y(t)$ output vector
we want to observe $vR1(t)$

from the schematic

$$vR1(t) = R1 \ iL1(t)$$

$$vR1(t) = \begin{pmatrix} R1 \\ 0 \end{pmatrix} \begin{pmatrix} iL1(t) \\ vC1(t) \end{pmatrix} + \begin{pmatrix} 0 \end{pmatrix} u(t)$$

$y(t)$

C

$x(t)$

D $u(t)$

The resulting **4** matrices **A B C D** are *intrinsically characteristic* of our **circuit**

$$A = \begin{pmatrix} -\frac{R1}{L1} & -\frac{1}{L1} \\ \frac{1}{C1} & 0 \end{pmatrix}$$

$$B = \begin{pmatrix} \frac{1}{L1} \\ 0 \end{pmatrix}$$

$$C = \begin{pmatrix} R1 \\ 0 \end{pmatrix}$$

$$D = \begin{pmatrix} 0 \end{pmatrix}$$

the STATE MODEL

*is giving us the opportunity to
answer to the frightening question :*



***What is the rational behind
transfer function denominator ???***

We claimed that the 4 matrices A B C D are characteristic of the circuit
does that means that they are independent of the “processing domain” ???

In other words, are they meaningful as well, for example, to the Laplace “ S ” domain

$$\begin{aligned} \dot{x}(t) &\rightarrow X(s) \\ \ddot{x}(t) &\rightarrow s X(s) \end{aligned}$$

$$\begin{aligned} \dot{x}(t) &= Ax(t) + Bu(t) \\ y(t) &= Cx(t) + Du(t) \end{aligned}$$

$$\begin{aligned} s X(s) &= AX(s) + Bu(s) \\ Y(s) &= CX(s) + Du(s) \end{aligned}$$

$$\frac{Y(s)}{u(s)} = H(s) = C(sI - A)^{-1}B + D$$

$$H(s) = \frac{C [\text{cof}(sI - A)]^T B + D Q(s)}{Q(s)}$$

with $Q(s) = \det(sI - A)$ Characteristic Polynomial of matrix A

This implies that the **zeros** of this **Characteristic Polynomial** of matrix **A**

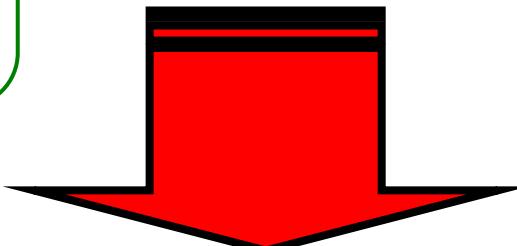
$$Q(s) = \det(sI - A) \quad \text{are also the } \textbf{POLES} \text{ of } H(s)$$

and so . . . the **POLES** of **H(s)** are from the expression of **Q(s)**

the **Eigen values** of matrix **A** ("valeurs propres" in french)

$$A = \begin{pmatrix} R1 & 1 \\ \frac{L1}{C1} & \frac{1}{L1} \\ \frac{1}{C1} & 0 \end{pmatrix}$$

Eigen values



$$-\frac{R1}{2L1} - \frac{\sqrt{C1^2 R1^2 - 4 C1 L1}}{2 C1 L1}$$

$$-\frac{R1}{2L1} + \frac{\sqrt{C1^2 R1^2 - 4 C1 L1}}{2 C1 L1}$$

POLES of the **Transfer Function** **H(s)**

$$A = \begin{pmatrix} R1 & 1 \\ L1 & L1 \\ 1 & 0 \\ \hline C1 & \end{pmatrix} \Rightarrow Q(s) = \det(sI - A)$$

please note that such **theoretical denominator expression**

is **literally** $(s - p_1) * (s - p_2)$

and thus has an **unit s^2** coefficient

$$Q(s) = s^2 + \frac{R1s}{L1} + \frac{1}{C1L1}$$

$$= \frac{C1L1s^2 + C1R1s + 1}{C1L1}$$

If you remember, this is effectively coherent with what we have found for the "**denominators**" when we used the **Kirchhoff laws** to compute **Laplace Transfer Functions**

$$\frac{vr1}{v} = \frac{C1R1s}{C1L1s^2 + C1R1s + 1}$$

$$\frac{i11}{v} = \frac{C1s}{C1L1s^2 + C1R1s + 1}$$

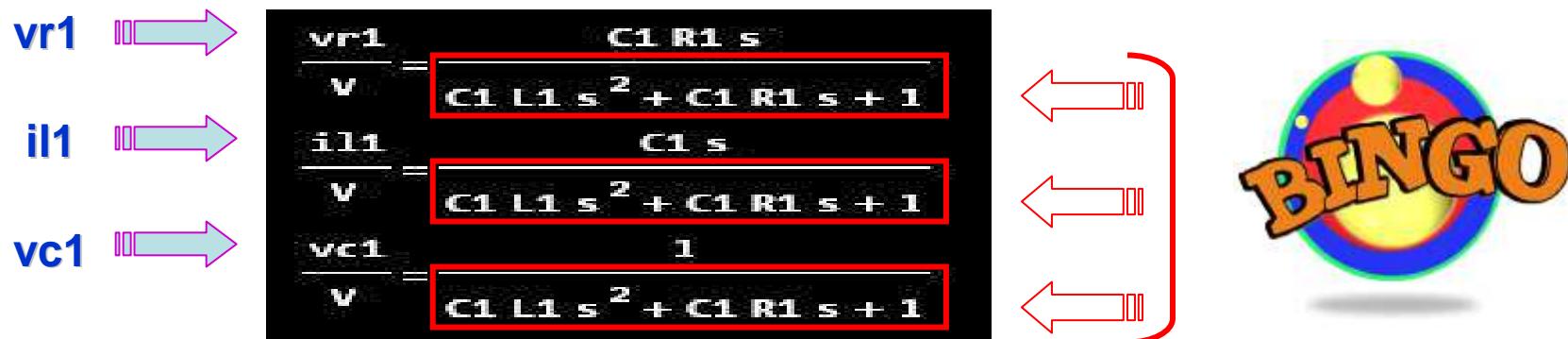
$$\frac{vc1}{v} = \frac{1}{C1L1s^2 + C1R1s + 1}$$

So, what have we learned ???

The **polynomial** that, *for a given circuit*, sounds “strangely” reappearing on the **denominator** of its various **output Transfer Functions** is getting *a justification* based on the

Characteristic Polynomial of the STATE matrix A

which is **independent** of the **observed OUTPUT**



So The expression “**Characteristic Polynomial**” is therefore taking its full sense since it truly **characterizes** the **INTRINSIC STATE** of the circuit

For the fun, let's verify that the **State Model** $vR1(s)$ **Transfer Function**

is effectively
the same than
the one found
with **Kirchhoff**

$$A = \begin{pmatrix} -\frac{R1}{L1} & -\frac{1}{L1} \\ \frac{1}{C1} & 0 \end{pmatrix} \quad B = \begin{pmatrix} \frac{1}{L1} \\ 0 \end{pmatrix} \quad C = \begin{pmatrix} R1 \\ 0 \end{pmatrix} \quad D = \begin{pmatrix} 0 \end{pmatrix}$$

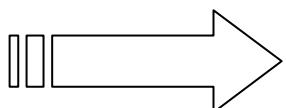
$$\frac{vR1(s)}{u(s)} = \frac{C [cof(sI - A)]^T B + D Q(s)}{Q(s)}$$

$$Q(s) = s^2 + \frac{R1 s}{L1} + \frac{1}{C1 L1}$$

$$[cof(sI - A)]^T =$$

transpose of the MATRIX of cofactors

$$\begin{bmatrix} s & -\frac{1}{L1} \\ \frac{1}{C1} & \frac{L1 s + R1}{L1} \end{bmatrix}$$

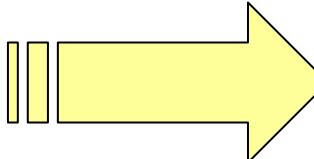


$$\frac{vR1(s)}{u(s)} = \frac{\frac{R1 s}{L1}}{\frac{C1 L1 s^2 + C1 R1 s + 1}{C1 L1}} = \frac{C1 R1 s}{C1 L1 s^2 + C1 R1 s + 1}$$

QED !!!

STATE MODEL



expression(t) ||  ***x(t)***

$$\dot{x} = ax + bu$$

$$\dot{x}(t) - ax(t) = bu(t)$$

$$e^{-at}(\dot{x}(t) - ax(t)) = e^{-at}(bu(t))$$

$$e^{-at}\dot{x} - e^{-at}ax = \frac{d}{dt}(e^{-at}x(t))$$

$$\frac{d}{dt}(e^{-at}x(t)) = e^{-at}bu$$

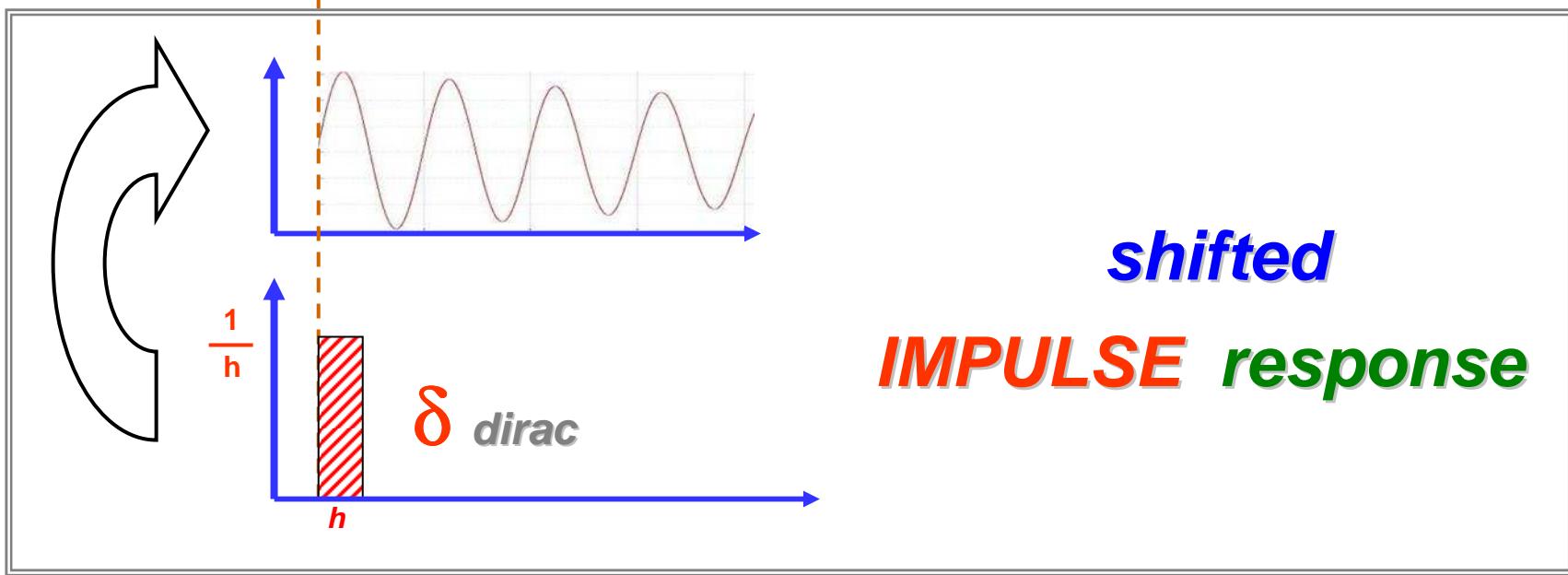
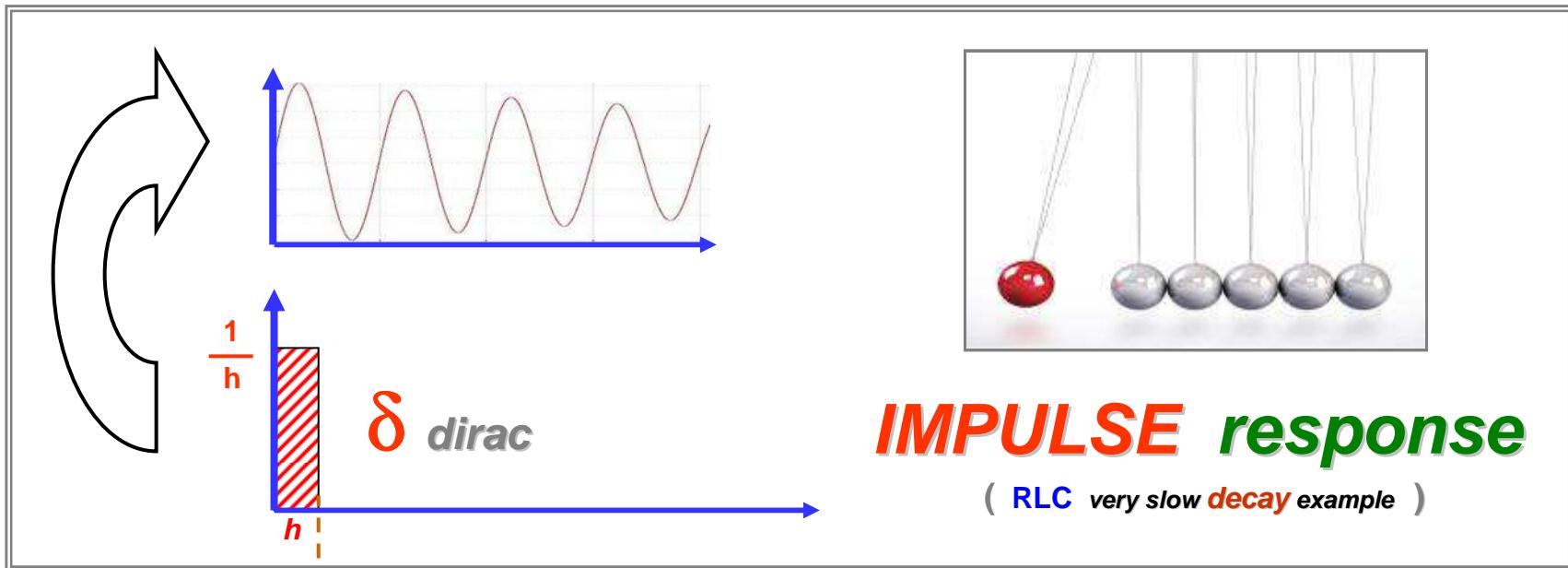
$$\int_0^t \frac{d}{d\tau}(e^{-a\tau}x(\tau)) d\tau = \int_0^t e^{-a\tau}bu(\tau) d\tau$$

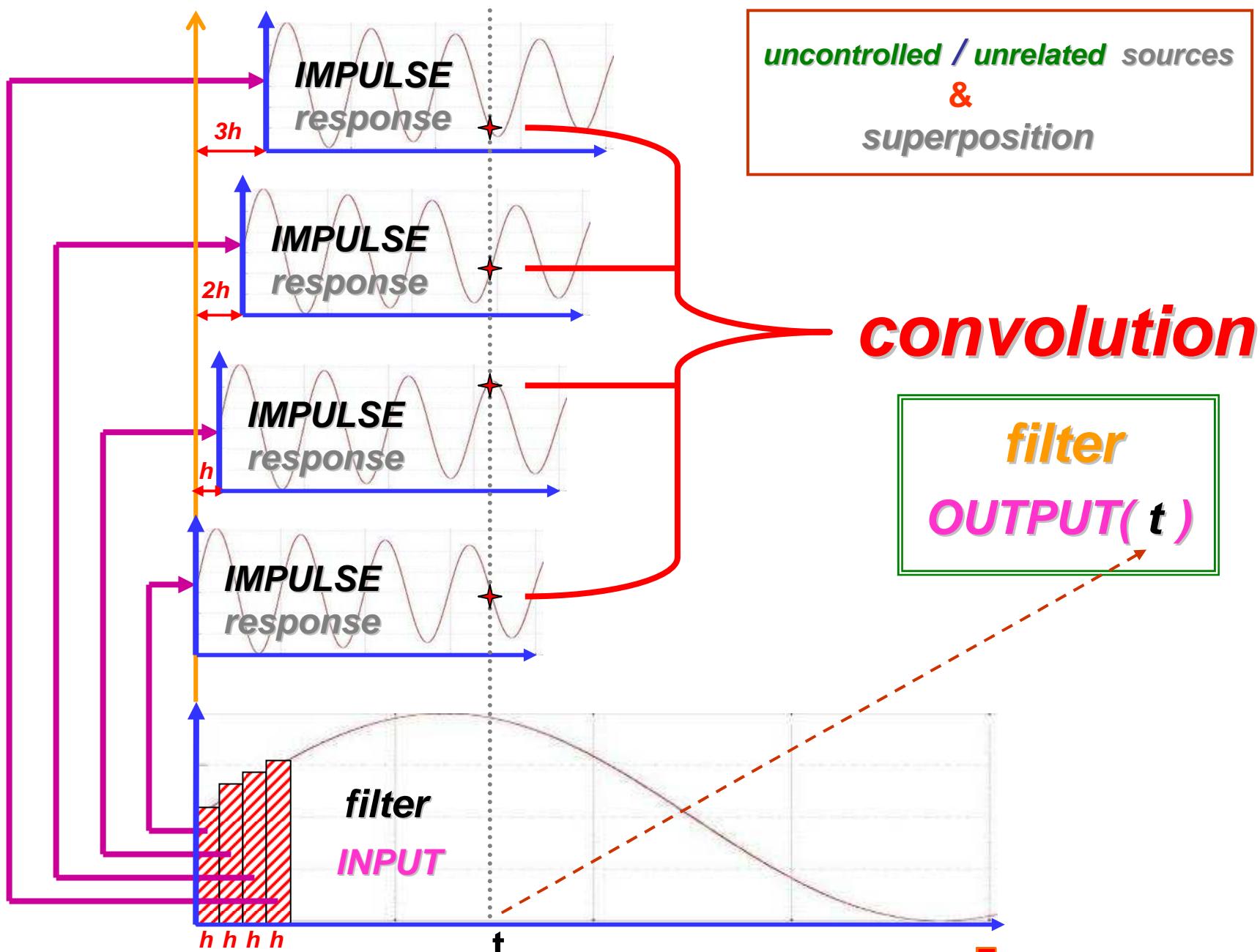
$$e^{-at}x(t) - x(0) = \int_0^t e^{-a\tau}bu(\tau) d\tau$$

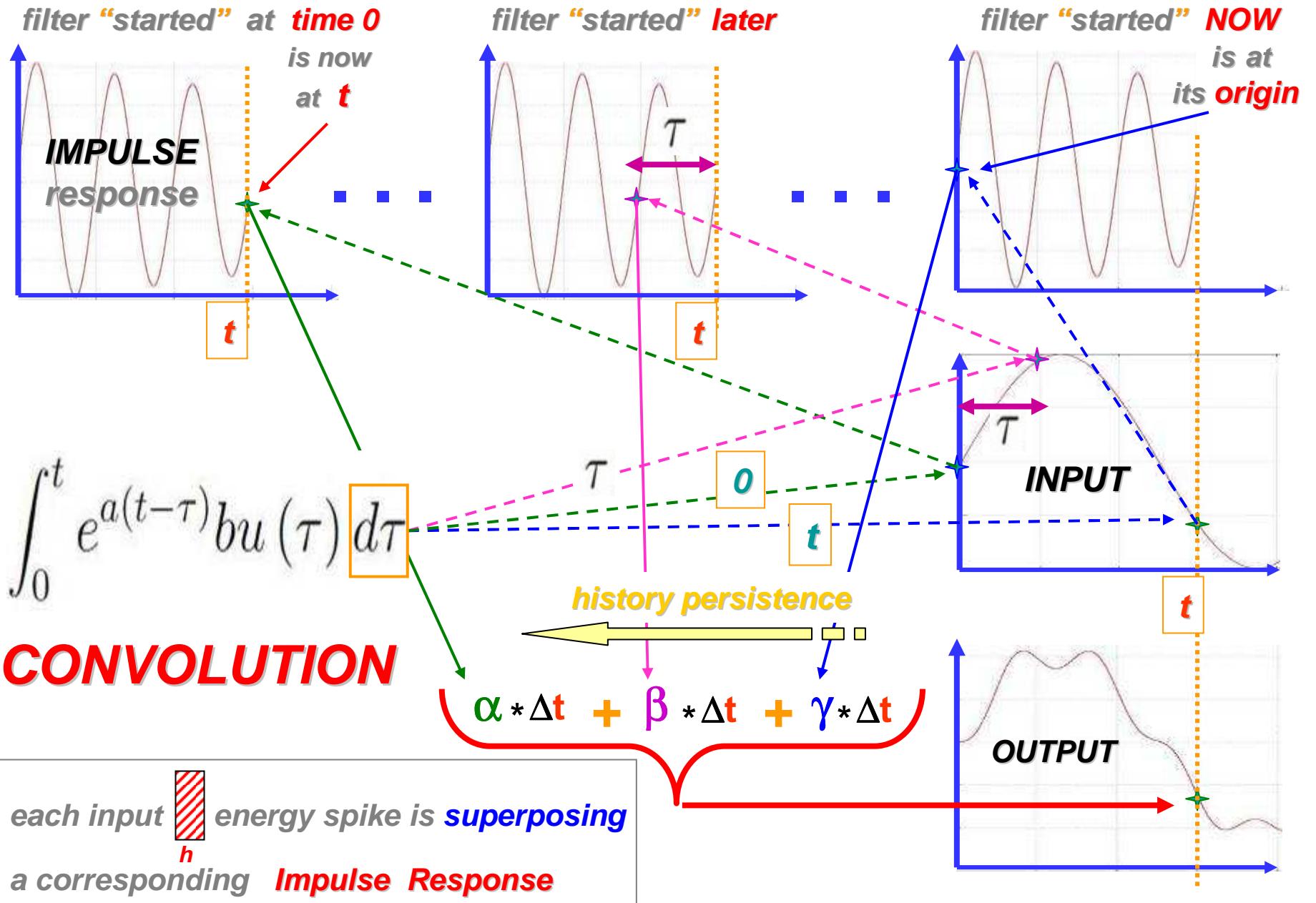
$$x(t) = e^{at}x(0) + \int_0^t e^{a(t-\tau)}bu(\tau) d\tau$$

convolution

integral





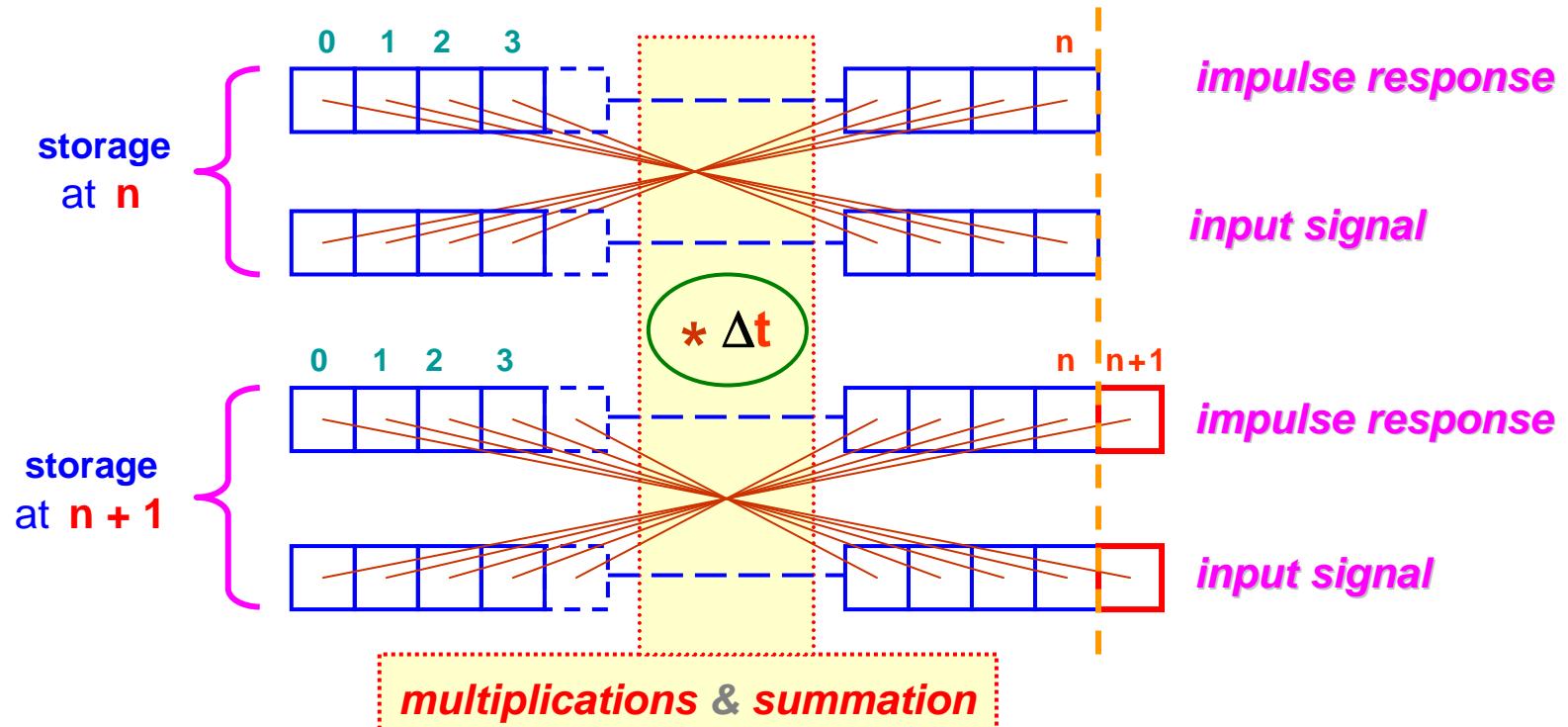


Unfortunately,

brut force convolution numerical usage is unlikely

💣 *unknown **input signal function $u(t)$** prevents **direct integration***

💣 *sample based discrete numerical evaluation is very **resource consuming***



STATE MODEL



expression(t) $x(t)$



$x(t)$ $x(t + h)$

$$x(t) = e^{at} x(0) + \int_0^t e^{a(t-\tau)} bu(\tau) d\tau$$

t_0

$$x(t_0) = e^{at_0} x(0) + \int_0^{t_0} e^{a(t_0-\tau)} bu(\tau) d\tau$$

t

$$x(t) = e^{at} x(0) + \int_0^{t_0} e^{a(t-\tau)} bu(\tau) d\tau + \int_{t_0}^t e^{a(t-\tau)} bu(\tau) d\tau$$

$$x(t) = e^{a(t-t_0)} \left(e^{at_0} x(0) + \int_0^{t_0} e^{a(t_0-\tau)} bu(\tau) d\tau \right) + \int_{t_0}^t e^{a(t-\tau)} bu(\tau) d\tau$$

$$x(t) = e^{a(t-t_0)} x(t_0) + \int_{t_0}^t e^{a(t-\tau)} bu(\tau) d\tau$$

Recursive Convolution Formula (RCF)

STATE MODEL

“time domain” evolution

STATE matrix A

$$x(t) = e^{A(t-t_0)} x(t_0) + \int_{t_0}^t e^{A(t-\tau)} B u(\tau) d\tau$$

new
STATE
at t

$u(t) = 0$
“free” mode
since STATE t_0

effective contribution from $u(t)$
“forced” mode
since STATE t_0

$$e^x = \sum_{k=0}^{\infty} \frac{1}{k!} x^k = 1 + x + \frac{1}{2}x^2 + \frac{1}{6}x^3 + \dots$$



$$e^{At} = \sum_{k=0}^{\infty} \frac{t^k A^k}{k!} = I + At + \frac{t^2 A^2}{2!} + \dots + \frac{t^k A^k}{k!} + \dots$$

The Transition Matrix

$$e^{At}$$

can be calculated

by searching *individual* matrix coefficient Inverse Laplace Transform

OR globally through *analytic* or *approximate* matricial methods

STATE MODEL



expression(t) \Rightarrow **$x(t)$**



$x(t)$ \Rightarrow **$x(t + h)$**



expression(t) \Rightarrow **e^{At}**



example: calculation based on **Inverse Laplace Transform** method

$$e^{At} = \text{Inverse Laplace Transform}\{K^{-1}\} \text{ with } K = sI - A$$

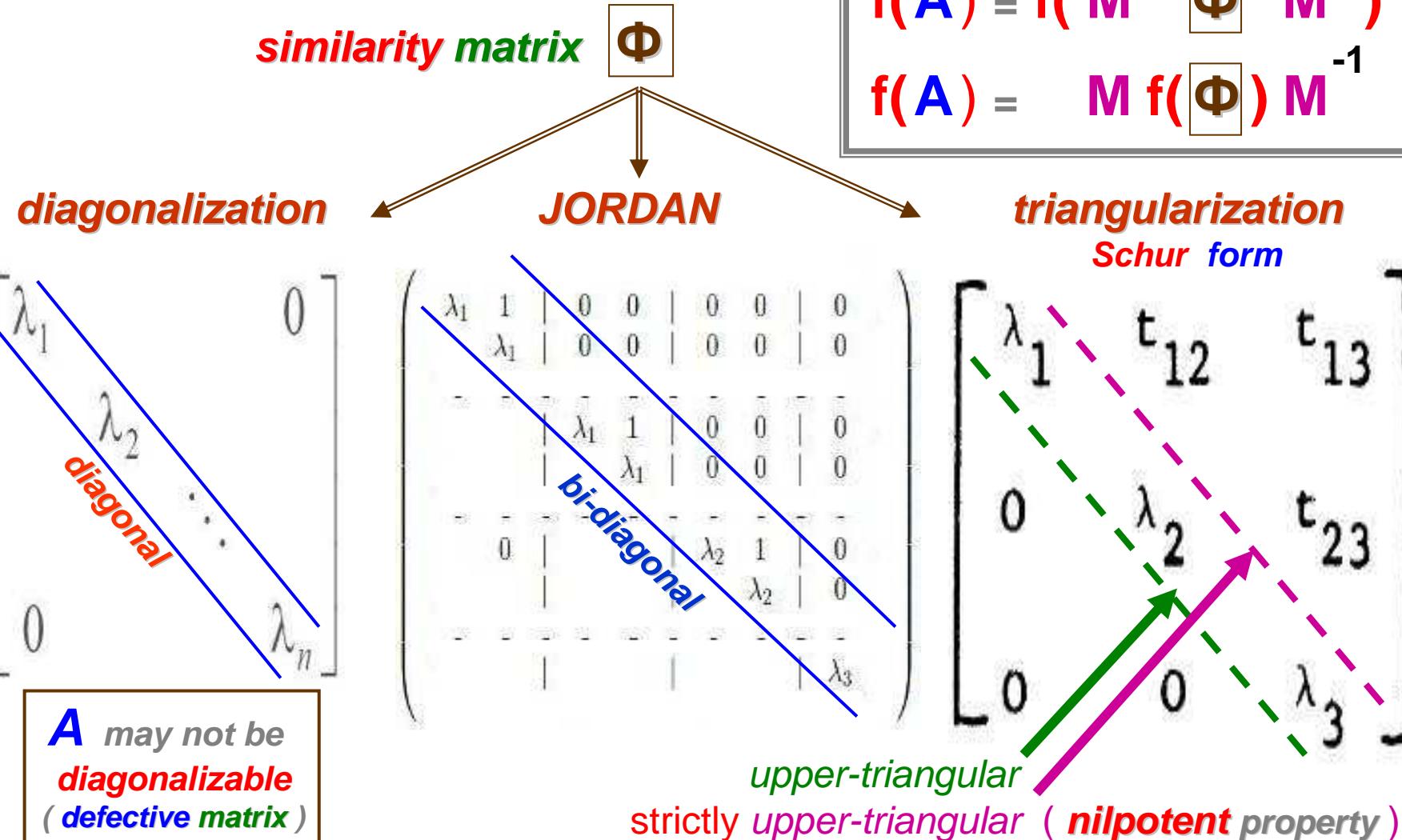
Characteristic Matrix
of matrix **A**

$$A = \begin{bmatrix} -\frac{R1}{L1} & -\frac{1}{L1} \\ \frac{1}{L1} & \frac{C1}{L1} \end{bmatrix} \Rightarrow K = \begin{bmatrix} \frac{L1s + R1}{L1} & \frac{1}{L1} \\ -\frac{1}{C1} & s \end{bmatrix} \Rightarrow K^{-1} = \begin{bmatrix} \frac{C1L1s}{C1L1s^2 + C1R1s + 1} & \frac{C1}{C1L1s^2 + C1R1s + 1} \\ \frac{L1}{C1L1s^2 + C1R1s + 1} & \frac{C1L1s + C1R1}{C1L1s^2 + C1R1s + 1} \end{bmatrix}$$

$$\Rightarrow e^{At} = \begin{bmatrix} \text{ILT}\left(\frac{sC1L1}{sC1R1+s^2C1L1+1}\right) & \text{ILT}\left(-\frac{C1}{sC1R1+s^2C1L1+1}\right) \\ \text{ILT}\left(\frac{L1}{sC1R1+s^2C1L1+1}\right) & \text{ILT}\left(\frac{C1R1+sC1L1}{sC1R1+s^2C1L1+1}\right) \end{bmatrix} = \begin{bmatrix} \text{ILT}\left(\frac{s}{(s-p1)(s-p2)}\right) & \text{ILT}\left(-\frac{1}{(s-p1)(s-p2)L1}\right) \\ \text{ILT}\left(\frac{1}{(s-p1)(s-p2)C1}\right) & \text{ILT}\left(\frac{R1+sL1}{(s-p1)(s-p2)L1}\right) \end{bmatrix}$$

transformed matrix is a symbolic “time domain” expression of e^{At}

MATRIX decomposition



example: e^{At} calculation based on **Jordan**

(even if this **A** is directly diagonalizable)

$$A = \begin{bmatrix} -\frac{\mathbf{R1}}{\mathbf{L1}} & -\frac{1}{\mathbf{L1}} \\ \frac{1}{\mathbf{C1}} & 0 \end{bmatrix} \quad \xrightarrow{\text{Jordan}} \quad \text{Jordan} = \begin{bmatrix} \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} + \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & 0 \\ 0 & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} - \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} \end{bmatrix} = \begin{pmatrix} p1 & 0 \\ 0 & p2 \end{pmatrix}$$

$$\boxed{A = MJM^{-1} \quad \text{similarity transform}}$$

$$A^k = (MJM^{-1})^k = M J^k M^{-1}$$

$$M = \begin{bmatrix} 1 & 1 \\ 1 & 1 \\ \frac{1}{p1 \mathbf{C1}} & \frac{1}{p2 \mathbf{C1}} \end{bmatrix} \quad e^{At} = \sum_{k=0}^{\infty} \frac{t^k A^k}{k!} = M e^{Jt} M^{-1}$$

transition matrix

$$= \begin{bmatrix} \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} + \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} - \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} + \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} - \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} \\ \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} - \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} + \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} - \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} + \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} \\ \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} + \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} - \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} + \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} - \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} \\ \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} - \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} + \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} - \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} & \frac{\sqrt{\mathbf{C1}^2 \mathbf{R1}^2 - 4 \mathbf{C1} \mathbf{L1} + \mathbf{C1} \mathbf{R1}}}{2 \mathbf{C1} \mathbf{L1}} \end{bmatrix}$$

that could be rewritten
based on the poles **p1** & **p2**
(that are conjugates)

$$\begin{bmatrix} \frac{p2 \cdot e^{p2 t} - p1 \cdot e^{p1 t}}{p2 - p1} & -\frac{e^{p2 t} - e^{p1 t}}{(p2 - p1)L1} \\ \frac{e^{p2 t} - e^{p1 t}}{(p2 - p1)C1} & \frac{p1 \cdot e^{p2 t} - p2 \cdot e^{p1 t}}{p2 - p1} \end{bmatrix}$$

even with **complex poles**, **A** real \rightarrow transition matrix real

Either way ...

$$e^{At} = \begin{bmatrix} \frac{p_2 e^{p_2 t} - p_1 e^{p_1 t}}{p_2 - p_1} & \frac{e^{p_2 t} - e^{p_1 t}}{p_2 - p_1} \\ \frac{e^{p_2 t} - e^{p_1 t}}{(p_2 - p_1)C_1} & \frac{p_1 e^{p_2 t} - p_2 e^{p_1 t}}{p_2 - p_1} \end{bmatrix}$$

poles
"distance"

let

$$p = r \pm v = r \pm z w$$



| | |
|----------|---------------|
| <i>r</i> | real |
| <i>z</i> | 1 or <i>i</i> |
| <i>w</i> | real > 0 |

r must be
negative
for stability

$$r = \frac{p_1 + p_2}{2} = -\frac{R_1}{Z L_1}$$

$$v = z w = z \left| f \left(z, \frac{p_1 - p_2}{2} \right) \right| = \frac{z \sqrt{z^2 (C_1^2 R_1^2 - 4 C_1 L_1)}}{2 C_1 L_1}$$

$$e^{At} = \begin{bmatrix} \frac{(v e^{2vt} + r e^{2vt} + v - r) e^{rt - vt}}{2v} & \frac{(-e^{vt} - 1)(e^{vt} + 1) e^{rt - vt}}{2 L_1 v} \\ \frac{(-e^{vt} - 1)(e^{vt} + 1) e^{rt - vt}}{2 C_1 v} & \frac{(v e^{2vt} - r e^{2vt} + v + r) e^{rt - vt}}{2v} \end{bmatrix}$$

$\overbrace{> 0 \rightarrow z = 1}$

$$V = \frac{1}{2 C_1 L_1} \sqrt{\frac{z^2 (C_1^2 R_1^2 - 4 C_1 L_1)}{z^2 (C_1^2 R_1^2 + 4 C_1 L_1)}}$$

$$e^{At} = \begin{bmatrix} \frac{(w e^{2wt} + r e^{2wt} + w - r) e^{rt-wt}}{2w} & \frac{-(e^{wt}-1)(e^{wt}+1) e^{rt-wt}}{2L_1 w} \\ \frac{(e^{wt}-1)(e^{wt}+1) e^{rt-wt}}{2C_1 w} & \frac{(w e^{2wt} - r e^{2wt} + w + r) e^{rt-wt}}{2w} \end{bmatrix}$$

$\overbrace{= 0 \rightarrow z = \text{don't care}}$

$$V = \frac{z}{2 C_1 L_1} \sqrt{z^2 (C_1^2 R_1^2 + 4 C_1 L_1)}$$

$$e^{At} = \infty$$

natural frequency

$$\frac{\sqrt{4 C_1 L_1 - C_1^2 R_1^2}}{4 \pi C_1 L_1}$$

Euler $e^{ix} = \cos x + i \sin x$

$\cos x = \frac{e^{ix} + e^{-ix}}{2}$
 $\sin x = \frac{e^{ix} - e^{-ix}}{2i}$

$\overbrace{< 0 \rightarrow z = \%i}$

$$V = \frac{\pm i}{2 C_1 L_1} \sqrt{\pm i^2 (C_1^2 R_1^2 - 4 C_1 L_1)}$$

decay
pulsation

$$e^{At} = \begin{bmatrix} \frac{\pm e^{rt} (r \sin(wt) + w \cos(wt))}{w} & \frac{\pm e^{rt} \sin(wt)}{L_1 w} \\ \frac{\pm e^{rt} \sin(wt)}{C_1 w} & \frac{\pm e^{rt} (r \sin(wt) - w \cos(wt))}{w} \end{bmatrix}$$

Those expressions are clearly highlighting the ***intrinsic influence*** of the ***Transition Matrix*** on the overall ***state*** behavior ***regardless of the system input activation***

Nineteen Dubious Ways to Compute the Exponential of a Matrix, Twenty-Five Years Later

Cleve Moler , Charles Van Loan

SIAM review 1978 and then 2003



not e^x specialized

- 1 *Taylor series*
- 2 *Padé approximation*
- 3 *Scaling and squaring*
- 4 *Chebyshev rational approximation*
- 5 *General purpose ODE solver*
- 6 *Single step ODE method*
- 7 *Multistep ODE solver*
- 8 *Cayley - Hamilton*
- 9 *Lagrange interpolation*
- 10 *Newton interpolation*
- 11 *Vandermonde matrix*
- 12 *Inverse Laplace transform*
- 13 *Companion matrix*
- 14 *Eigenvectors*
- 15 *Triangular systems of eigenvectors*
- 16 *Jordan canonical form*
- 17 *Schur decomposition*
- 18 *Block diagonal matrix*
- 19 *Splitting method*
- 20 *Krylov space methods*

(added in 2003)



Brook Taylor (1685 – 1731)



Henri Padé (1863 – 1953)

$$e^x \approx \text{“Tailor series”} \quad \text{order 3}$$

$$1 + x + \frac{x^2}{2} + \frac{x^3}{6} \quad \times$$

series development
approximation

“Padé approximants” order 3

$$e^z \approx \frac{1 + \frac{1}{2}z + \frac{1}{10}z^2 + \frac{1}{120}z^3}{1 - \frac{1}{2}z + \frac{1}{10}z^2 - \frac{1}{120}z^3}$$

“rational fraction”

$$e^z = \cfrac{1}{1 - \cfrac{z}{1 + \cfrac{\frac{1}{2}z}{1 - \cfrac{\frac{1}{6}z}{1 + \cfrac{\frac{1}{10}z}{1 - \cfrac{\frac{1}{120}z}{\ddots}}}}}}$$

“continued fraction” form

$$[L/M](z) = \frac{a_0 + a_1 z + \dots + a_L z^L}{b_0 + b_1 z + \dots + b_M z^M}$$

$$f(z) = \sum_{i=0}^{\infty} c_i z^i = [L/M](z) + O(z^{L+M+1})$$

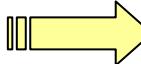
“approximant” error

Tailor Padé

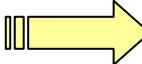


STATE MODEL

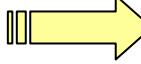


expression(t)  $x(t)$



$x(t)$  $x(t + h)$



expression(t)  e^{At}



unknown input $u(\tau)$  *approximation*



Semi Analytical Recursive Convolution algorithm (**SARC**)



$$\int_{t-h}^t e^{a(t-\tau)} u(\tau) d\tau$$



$$e^{at} \int_{t-h}^t e^{-a\tau} u(\tau) d\tau$$



$$e^{at} \int_{t-h}^t e^{-a\tau} \text{interpolation}(\tau) d\tau$$

Lagrange polynomial

$$u(\tau) = \sum_{i=0}^n L_i(\tau) u_i$$

with

$$L_i(\tau) = \prod_{\substack{j=0 \\ j \neq i}}^n \frac{(\tau - \tau_j)}{(\tau_i - \tau_j)}$$

$$\begin{aligned} L_i(\tau_i) &= 1 \\ L_i(\tau_j) &= 0 \quad (j \neq i) \end{aligned}$$

order $n = 0$ $u(\tau)$ is approximated on the interval $[t-h, t]$ by a constant



order $n = 1$ $u(\tau)$ is approximated on the interval $[t-h, t]$ by a slope



etc ...

interpolation error

$$O(h^{n+1}) \quad \frac{h}{2} \xrightarrow{\text{error}} \frac{h}{2^{n+1}}$$

$$e(\tau) = \frac{u^{(n+1)}(\xi)}{(n+1)!} \lambda(\tau)$$

with

$$\lambda(\tau) = (\tau - \tau_0)(\tau - \tau_1)\dots(\tau - \tau_n)$$

$$\int_{t-h}^t e^{a(t-\tau)} u(\tau) d\tau \rightarrow e^{at} \int_{t-h}^t e^{-a\tau} \text{interpolation}(\tau) d\tau$$

order 0

$$e^{at} \int_{t-h}^t e^{-a\tau} \left[u_t \right] d\tau$$

order 1

$$e^{at} \int_{t-h}^t e^{-a\tau} \left[\frac{t-\tau}{h} u_{t-h} + \frac{\tau-(t-h)}{h} u_t \right] d\tau$$

$$\Phi = e^{ah}$$

constant

$$\frac{\Phi - 1}{a} u_t$$

$$\frac{1}{a} \left(\Phi - \frac{\Phi - 1}{ah} \right) u_{t-h} + \frac{1}{a} \left(\frac{\Phi - 1}{ah} - 1 \right) u_t$$



does not depend on

$$e^{at}$$

!!! (Markov process)

The “shaking” due to [t-h, t] input is independent of the current oldness of the system

putting everything together

$$x(t) = e^{at} x(0) + \int_0^t e^{a(t-\tau)} b u(\tau) d\tau$$

$$x(t) = e^{a(t-t_0)} x(t_0) + \int_{t_0}^t e^{a(t-\tau)} b u(\tau) d\tau$$

$$x(t) = e^{a h} x(t-h) + \int_{t-h}^t e^{a(t-\tau)} b u(\tau) d\tau$$

$$x(t) = e^{a h} x(t-h) + b e^{at} \int_{t-h}^t e^{-a\tau} \text{interpolation}(\tau) d\tau$$

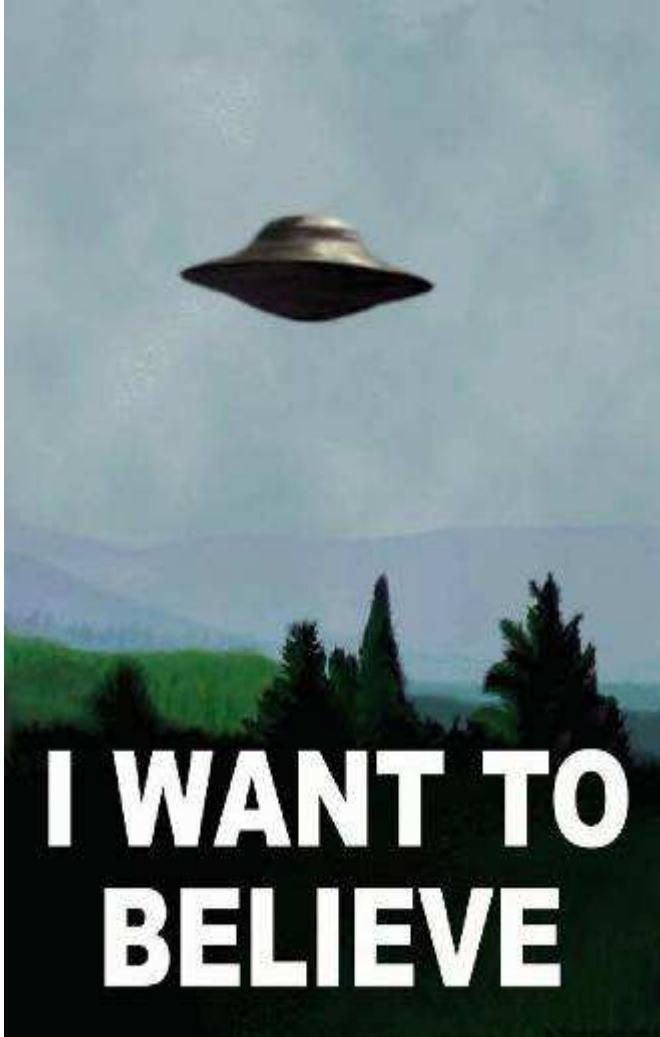
order 0

$$x_t = \Phi x_{t-h} + \frac{\Phi - 1}{a} b u_t$$

$$\Phi = e^{ah}$$

amazingly trivial
compared to Spice

constant *constant*



SARC equations are also applicable to STATE MODEL matrices



SARC *ERROR*



Can I help ???

input interpolation

order
step

$$\Phi = e^{Ah}$$

choose finer calculation method
(may be jeopardized by **poles** configuration)

Euler $e^{ix} = \cos x + i \sin x$

complex number

floating number

bits
digits
exponent

| IEEE STANDARD 754 | | | |
|-------------------|-------------------|--------------------|--------------------|
| float | double | long double | x87 standard |
| 32 | 64 | 80 96 | 128 |
| 7 | 16 | 19 | 34 |
| $\approx \pm 38$ | $\approx \pm 308$ | $\approx \pm 4932$ | $\approx \pm 4932$ |

contribution of *interpolation* to

SARC step *ERROR*

$$x(t) = e^{a(t-t_0)} x(t_0) + \int_{t_0}^t e^{a(t-\tau)} b u(\tau) d\tau$$

$$x(t) = e^{at} x(t-h) + b e^{at} \int_{t-h}^t e^{-a\tau} \text{interpolation}(\tau) d\tau$$

Lagrange interpolation step error

$$\mathcal{E}_{\text{step interpolation}} = b e^{at} \frac{\mathbf{U}^{(n+1)}(\xi)}{(n+1)!} \int_{t-h}^t e^{-a\tau} \lambda(\tau) d\tau$$

with

$$\lambda(\tau) = (\tau - t_0) \dots * (\tau - t_n)$$

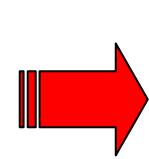
$\xi \in \text{interpolation involved segments}$

order n = 0

SOB

step constant = U_{t-h}

$$\varepsilon(k)_{\substack{\text{step} \\ \text{interpolation}}} = \frac{b e^{at}}{1!} \sup_{[t-h, t]} \left\{ \frac{d u(t)}{dt} \right\} \int_{t-h}^t e^{-a\tau} (\tau - (t-h)) d\tau$$



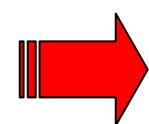
$$\varepsilon(k)_{\substack{\text{step} \\ \text{interpolation}}} = \frac{1}{a^2} (\Phi - ah - 1) \sup_{[t-h, t]} \left\{ b \frac{d u(t)}{dt} \right\}$$

order n = 0

SOE

step constant = U_t

$$\varepsilon(k)_{\substack{\text{step} \\ \text{interpolation}}} = \frac{b e^{at}}{1!} \sup_{[t-h, t]} \left\{ \frac{d u(t)}{dt} \right\} \int_{t-h}^t e^{-a\tau} (\tau - (t)) d\tau$$



$$\varepsilon(k)_{\substack{\text{step} \\ \text{interpolation}}} = \frac{1}{a^2} ((1 - ah) \Phi - 1) \sup_{[t-h, t]} \left\{ b \frac{d u(t)}{dt} \right\}$$

order $n = 1$

slope between

u_{t-h}

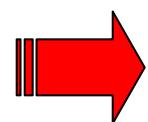
&

u_t

S1

$$\varepsilon(k) \underset{\substack{\text{step} \\ \text{interpolation}}}{=} \frac{b e^{at}}{2!} \sup_{[t-h, t]} \left\{ \frac{d^2 u(t)}{dt^2} \right\} \int_{t-h}^t e^{-a\tau} (\tau - (t-h)) (\tau - (t)) d\tau$$

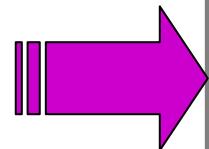
interpolation points



$$\varepsilon(k) \underset{\substack{\text{step} \\ \text{interpolation}}}{=} \frac{1}{2a^3} \left((2 - ah)\Phi - ah - 2 \right) \sup_{[t-h, t]} \left\{ b \frac{d^2 u(t)}{dt^2} \right\}$$

SARC *cumulated interpolation ERROR*

RECURSIVITY



$$\varepsilon(k) = \sum_{i=1}^k (\phi^{k-i} \varepsilon(i))_{\text{step interpolation}}$$

cumulated interpolation

constant "majorant"

MAJORANT

for large k



$$t = kh$$

$$\varepsilon(k) < \frac{1 - \phi^k}{1 - \phi}$$

cumulated max interpolation

$$\varepsilon_{\max}^{\text{step interpolation}}$$

$$\varepsilon(\text{only large } t) <$$

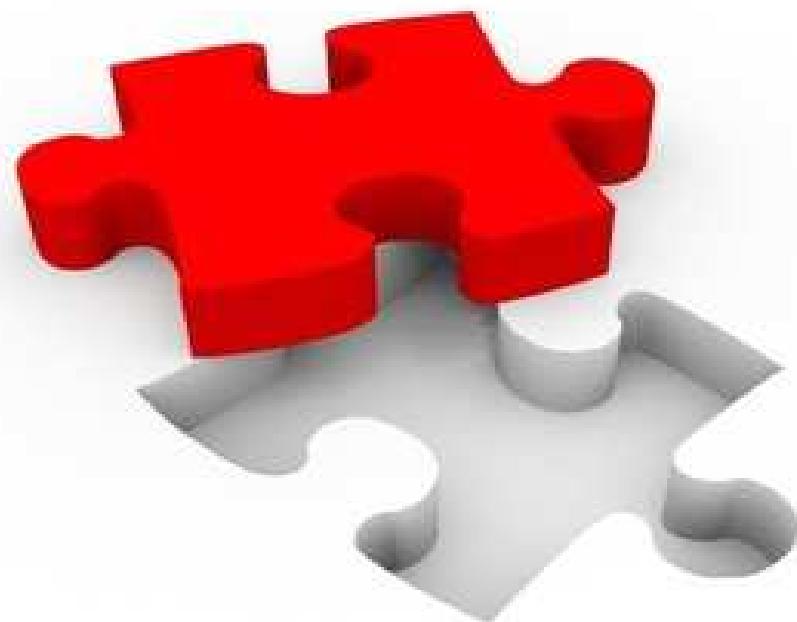
cumulated max interpolation

$$\frac{\varepsilon_{\max}^{\text{step interpolation}}}{1 - \phi}$$

asymptotic
error majorant

(proof of **absolute stability**)

*estimated **SARC error**, in practice*



cumulated SARC error can be **estimated** at **run time**
due to *interpolation*

example, for **order n = 0**

with **step constant** =

$$u_{t-h}$$

$$\frac{\Delta u}{\Delta t}$$



$$\varepsilon(k)_{\substack{\text{step} \\ \text{interpolation}}} = \frac{1}{a^2} \left[\Phi - ah - 1 \right] \left[b \frac{u(t) - u(t-h)}{h} \right]$$

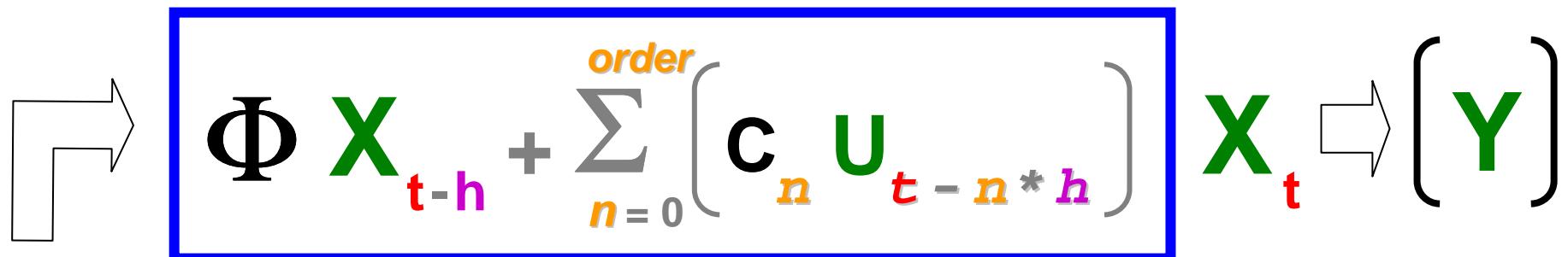
$$\varepsilon(k)_{\substack{\text{cumulated} \\ \text{interpolation}}} = \phi \varepsilon(k-1)_{\substack{\text{cumulated} \\ \text{interpolation}}} + \varepsilon(k)_{\substack{\text{step} \\ \text{interpolation}}}$$

$$\varepsilon(k)_{\substack{\text{cumulated max} \\ \text{interpolation}}} \approx \max \left(\left| \varepsilon(k-1)_{\substack{\text{cumulated} \\ \text{interpolation}}} \right|, \left| \varepsilon(k)_{\substack{\text{cumulated} \\ \text{interpolation}}} \right| \right)$$

SARC SUMMARY

$$\Phi = e^{Ah}$$

recursive convolution

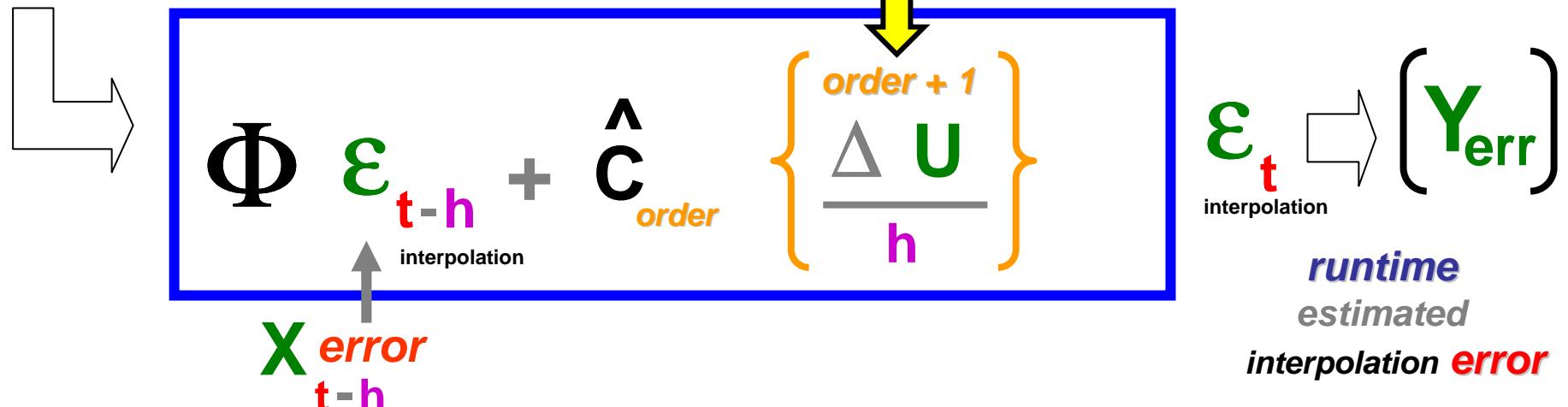


$[U]$ multi *inputs* vector

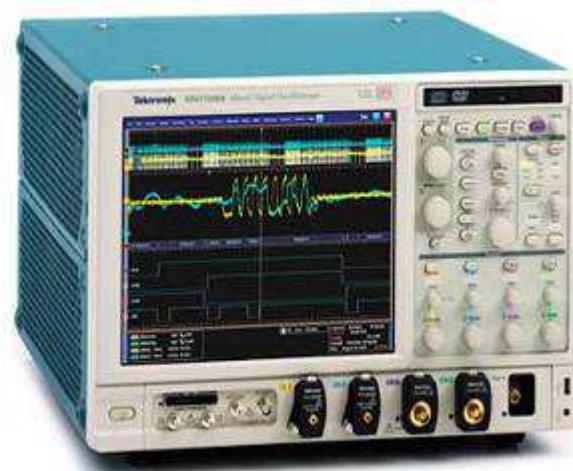
likely to engender low noise

multi *outputs* vectors

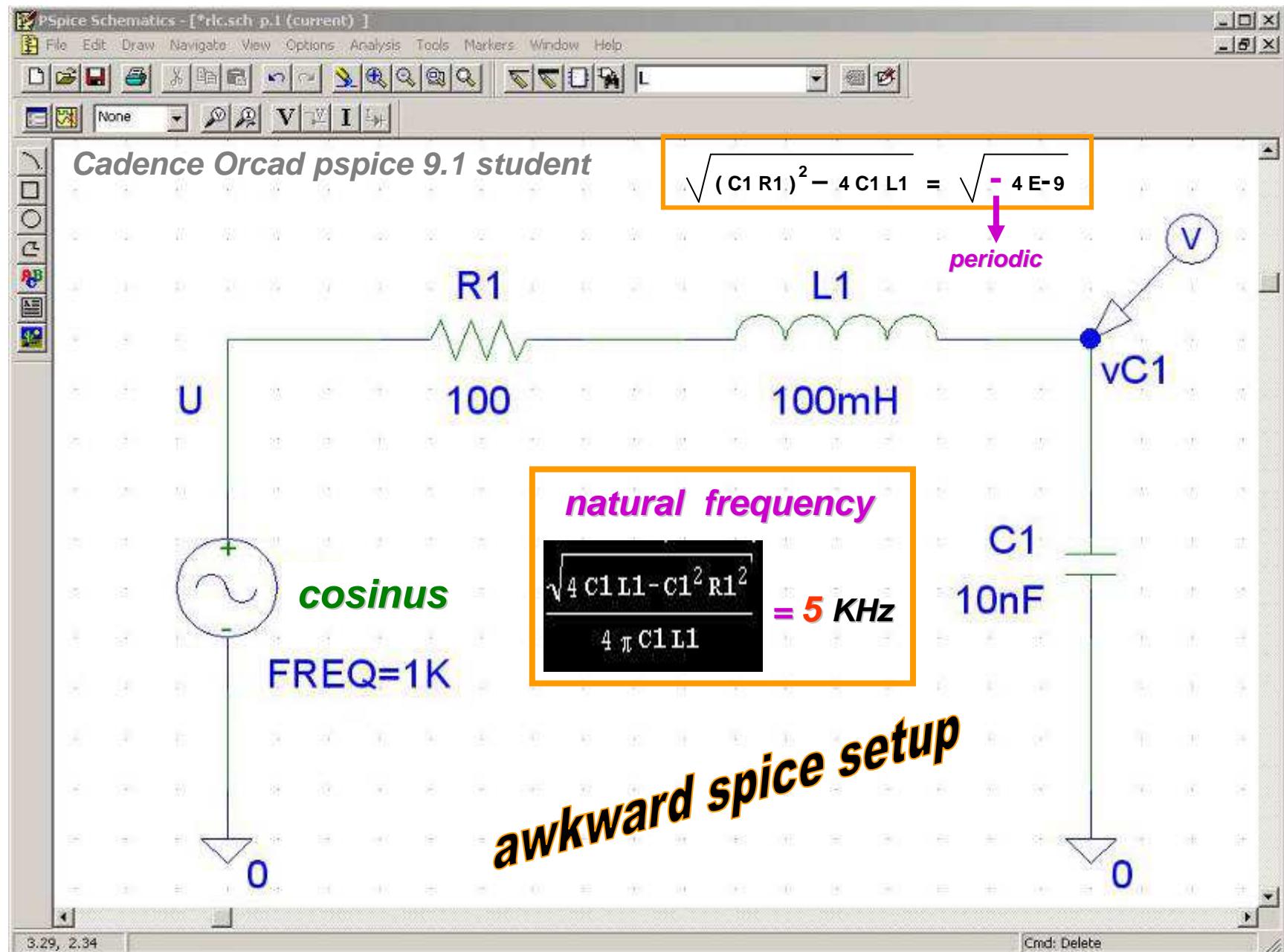
recursive convolution

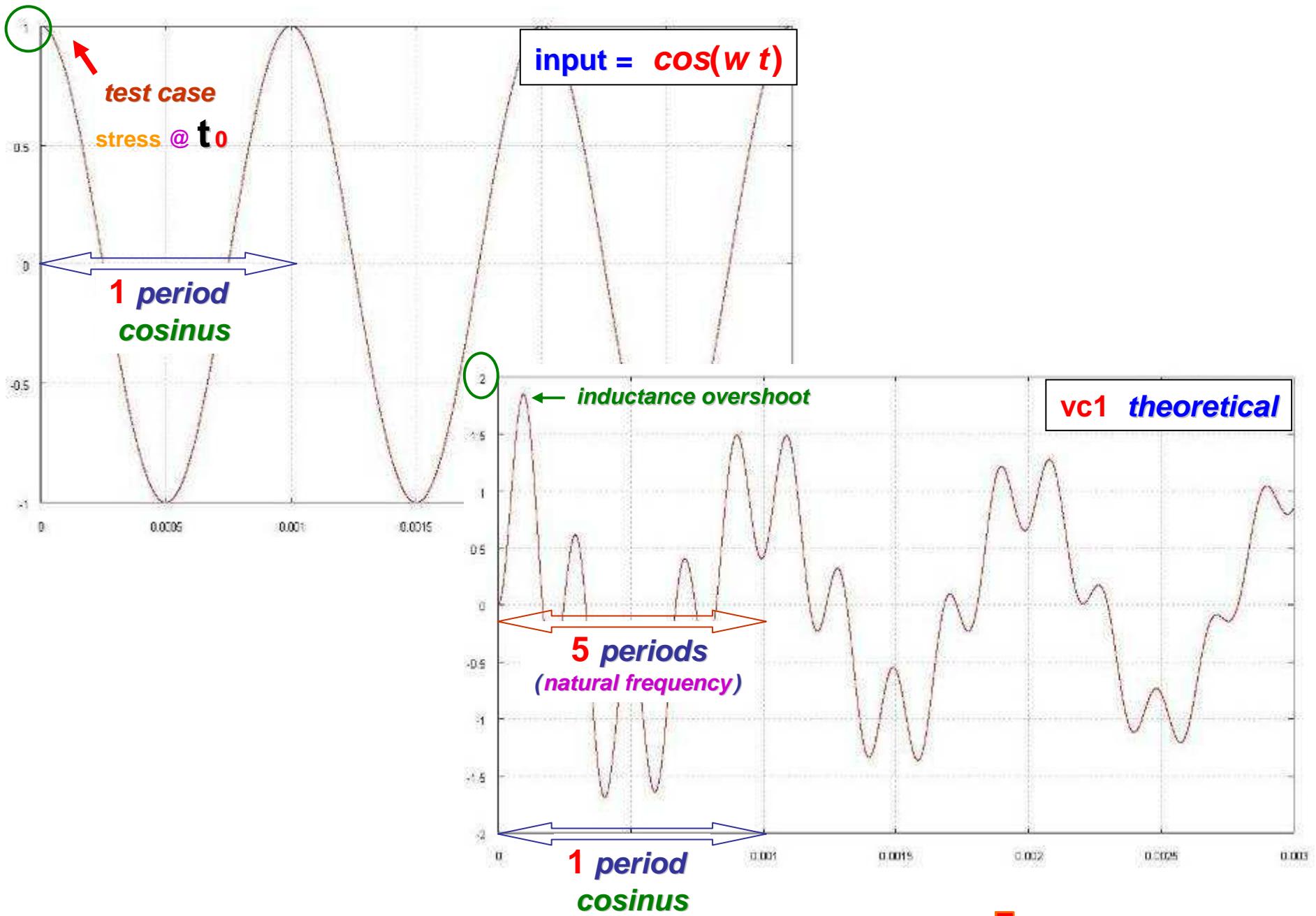


TEST

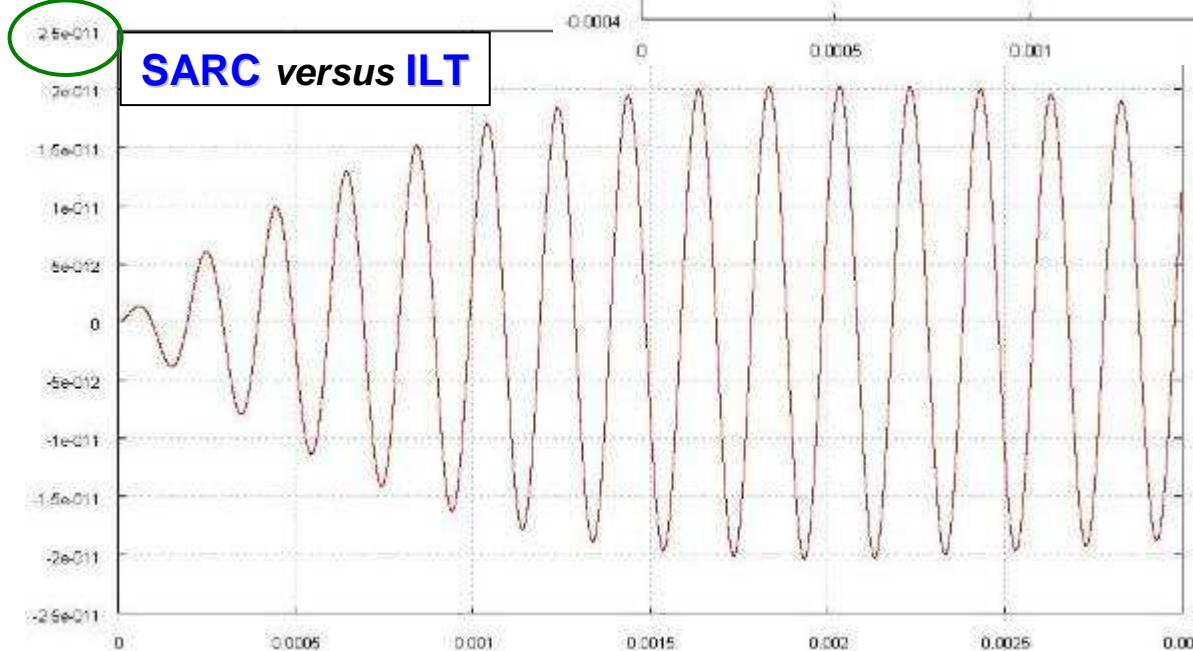
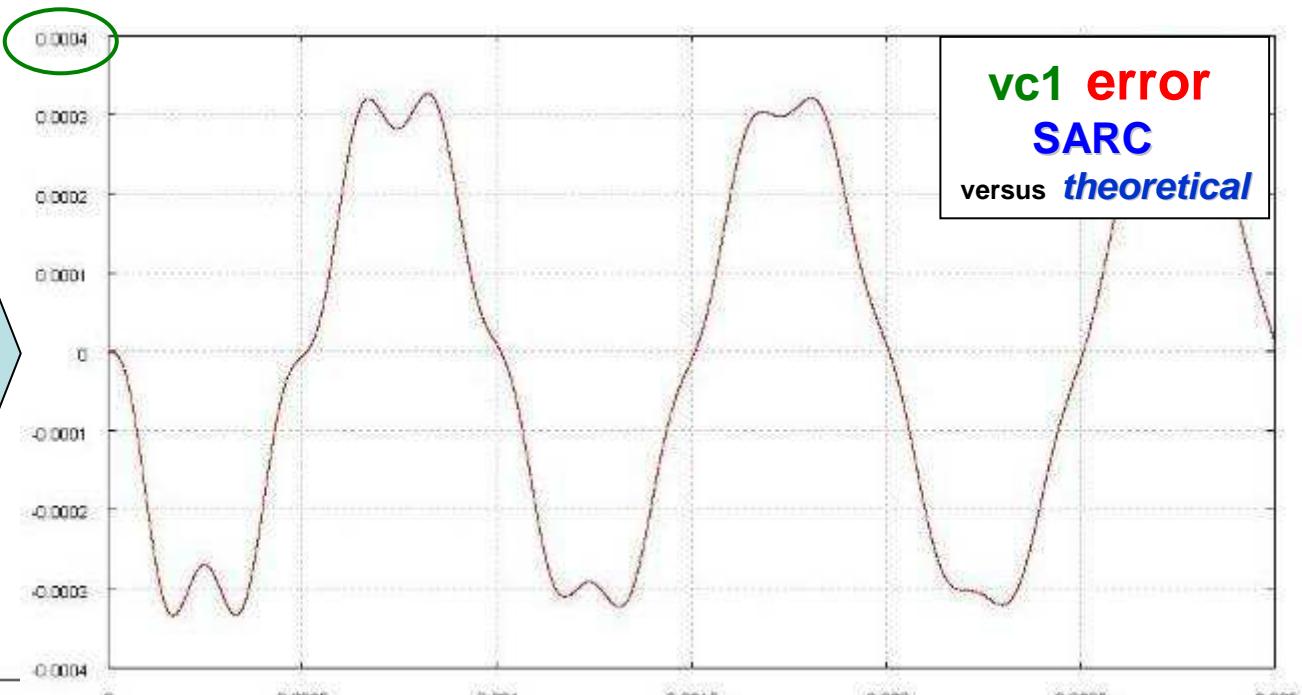
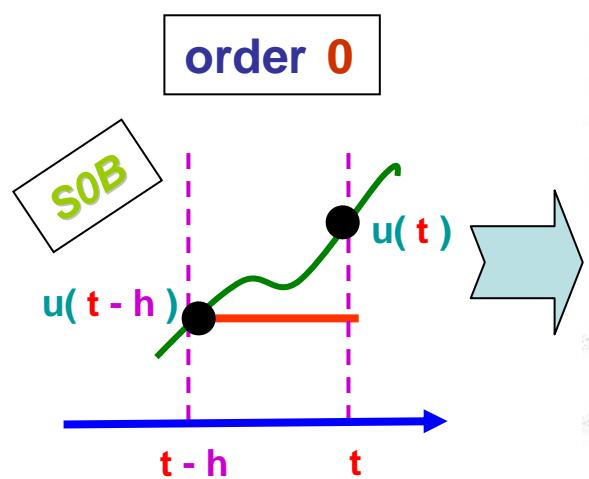


Jacques MEQUIN
NICE EWTBU / DTE
2012 j-mequin@ti.com



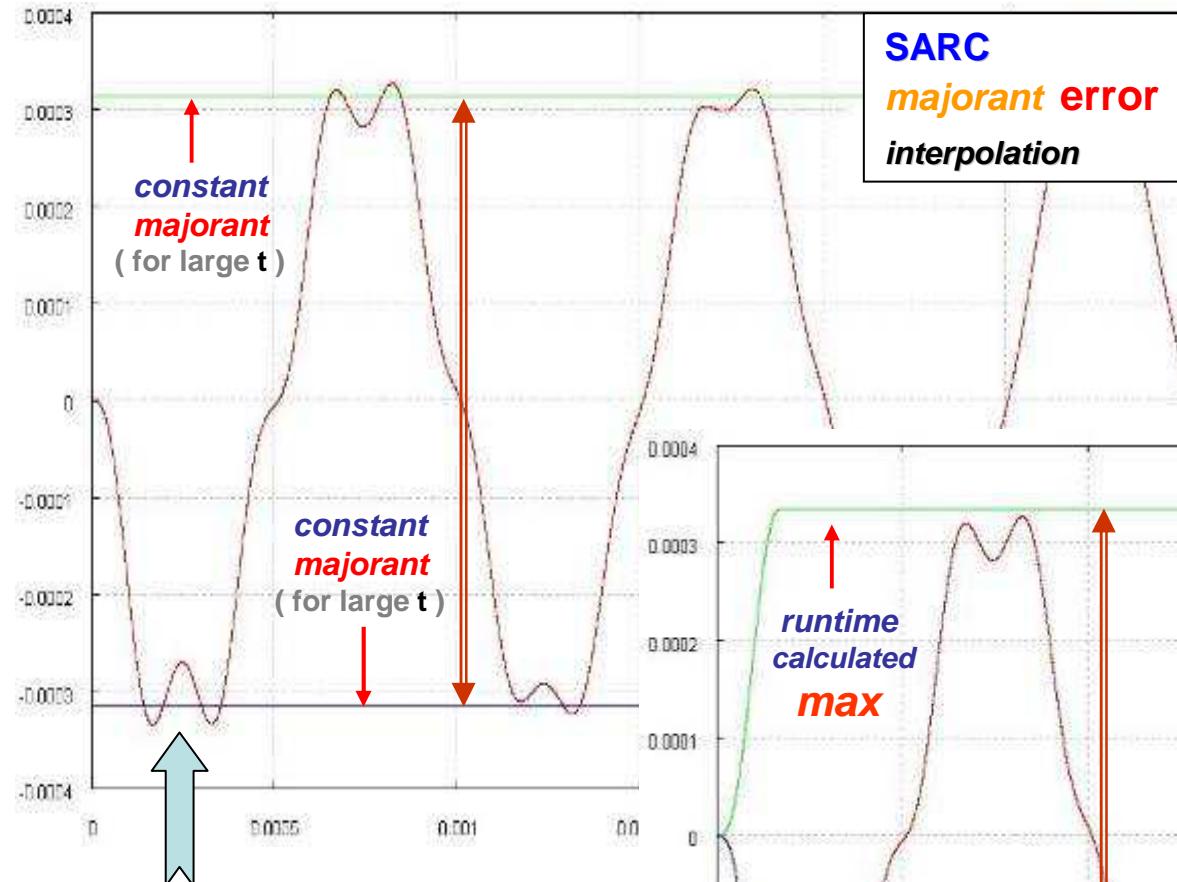


Lagrange interpolation



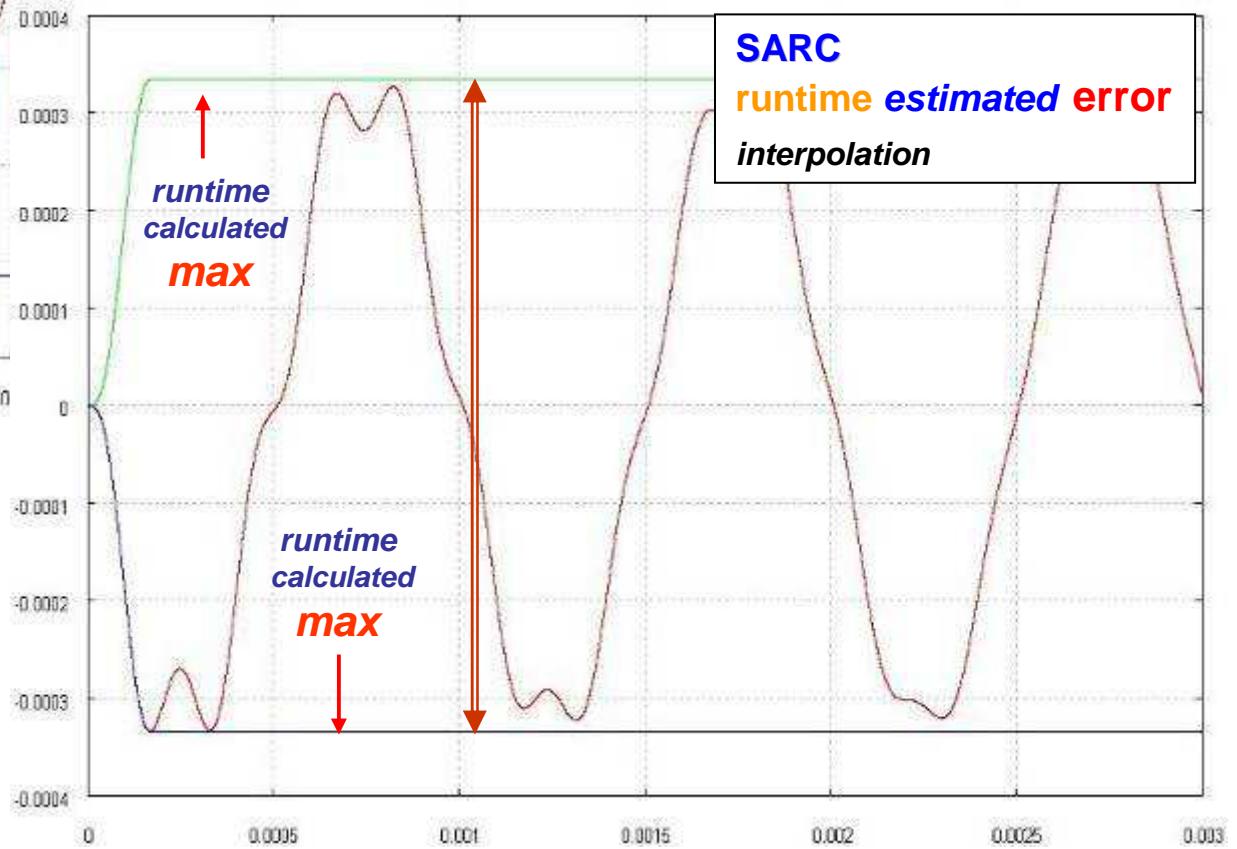
methods comparison

$\approx 10^{-11}$!!!



SARC
majorant error
interpolation

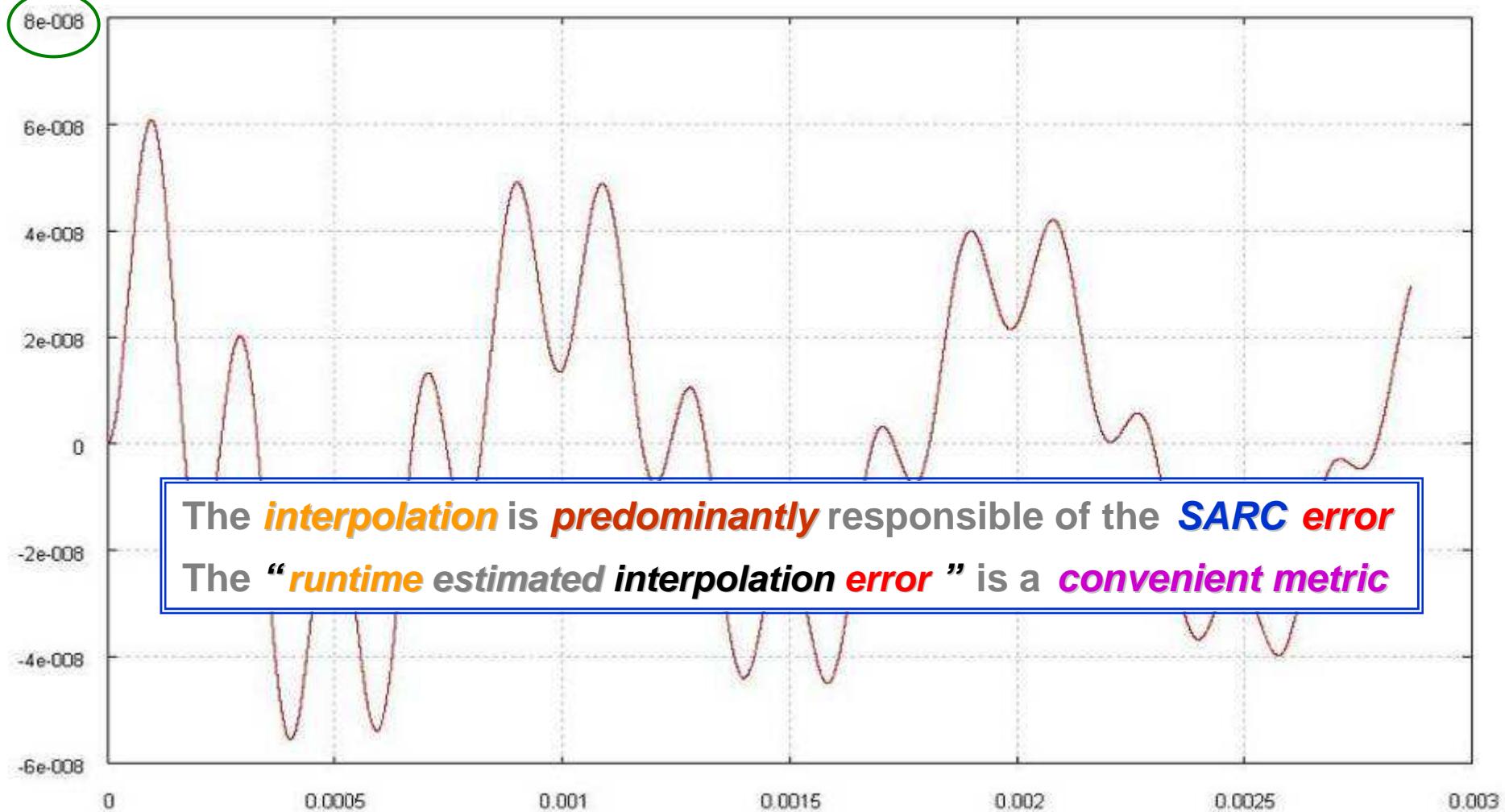
inaccurate
for small t

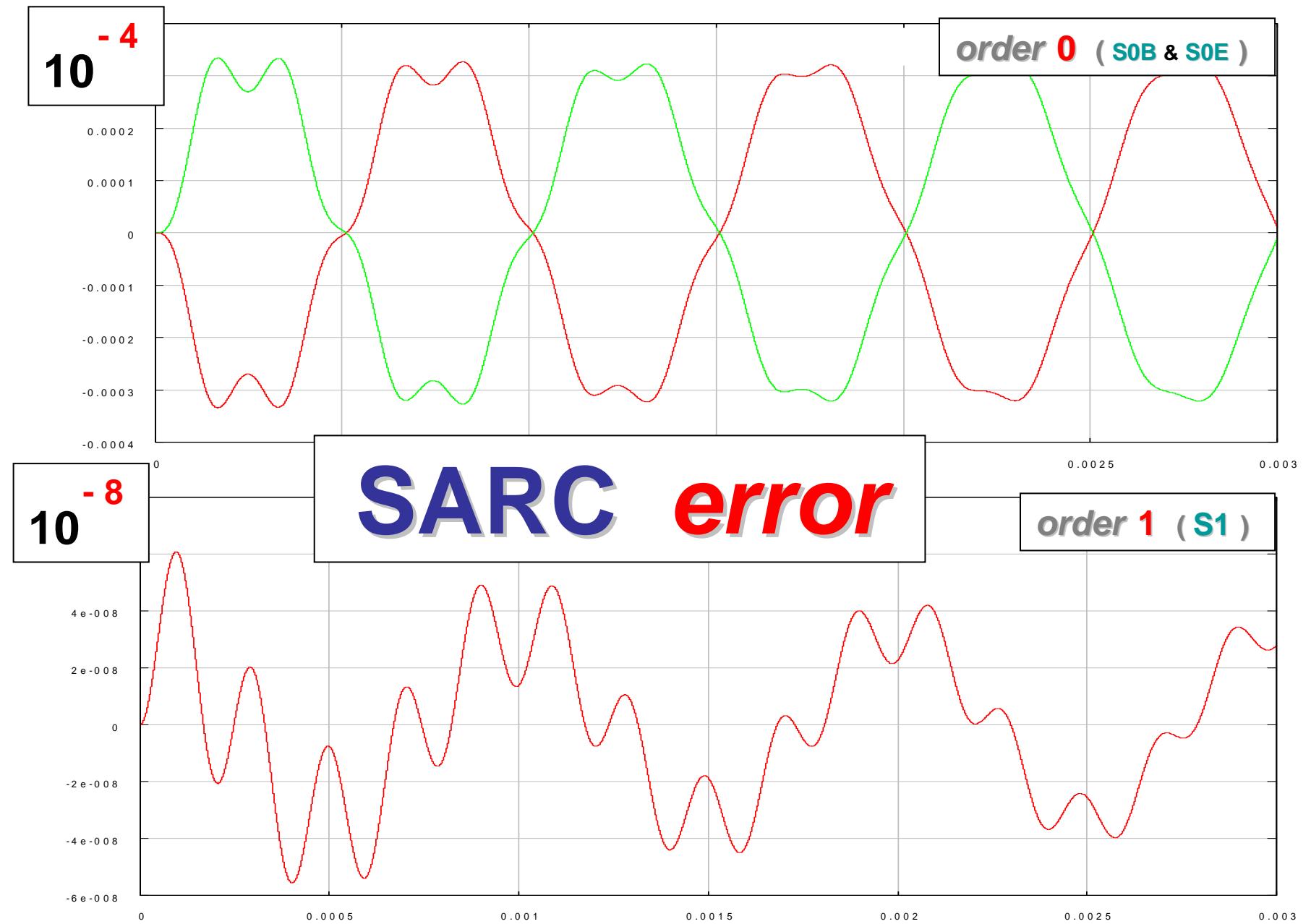


SARC
runtime estimated error
interpolation

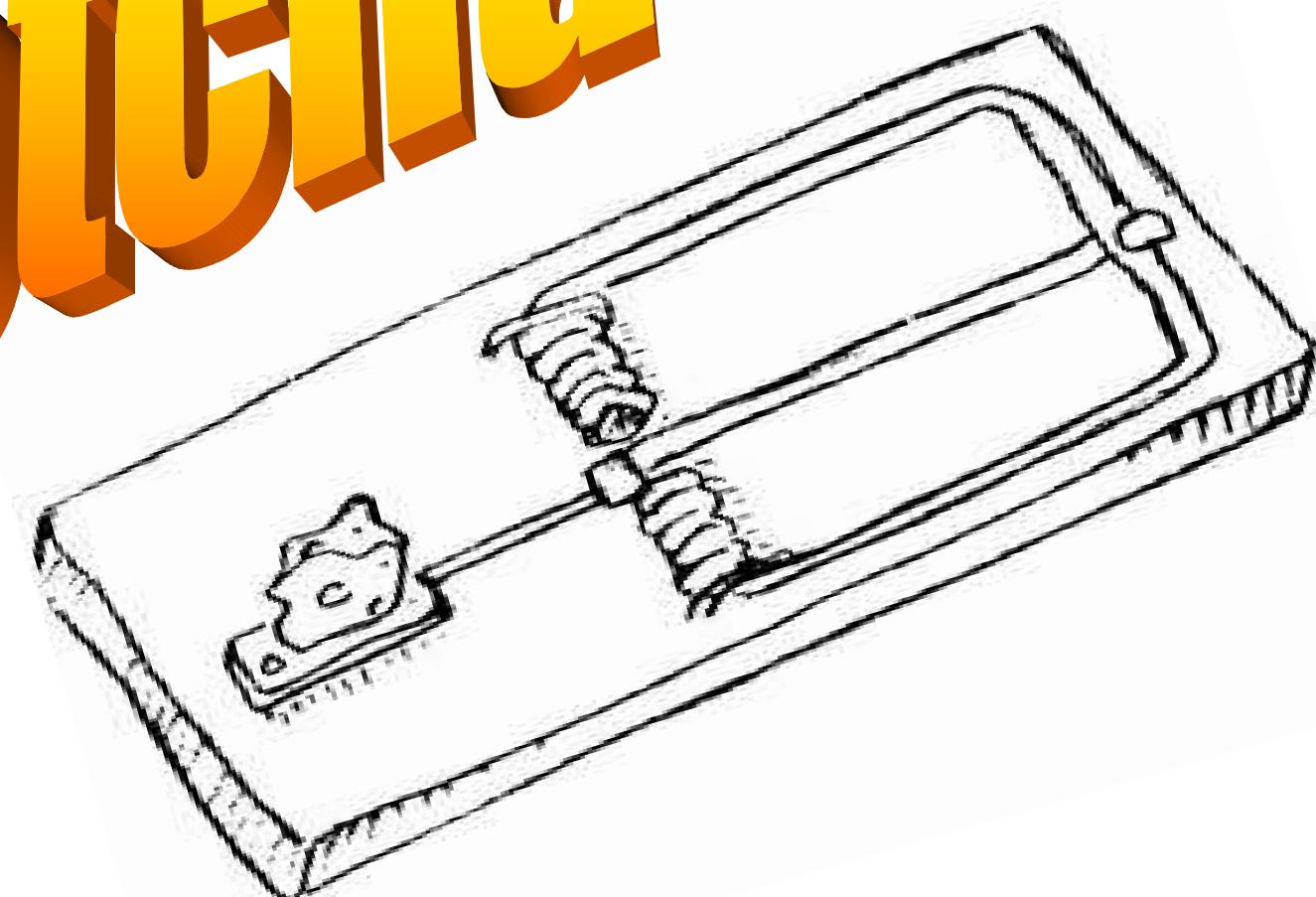
full SARC error

(*theoretical - SARC*) - *runtime estimated interpolation error*





Gotcha

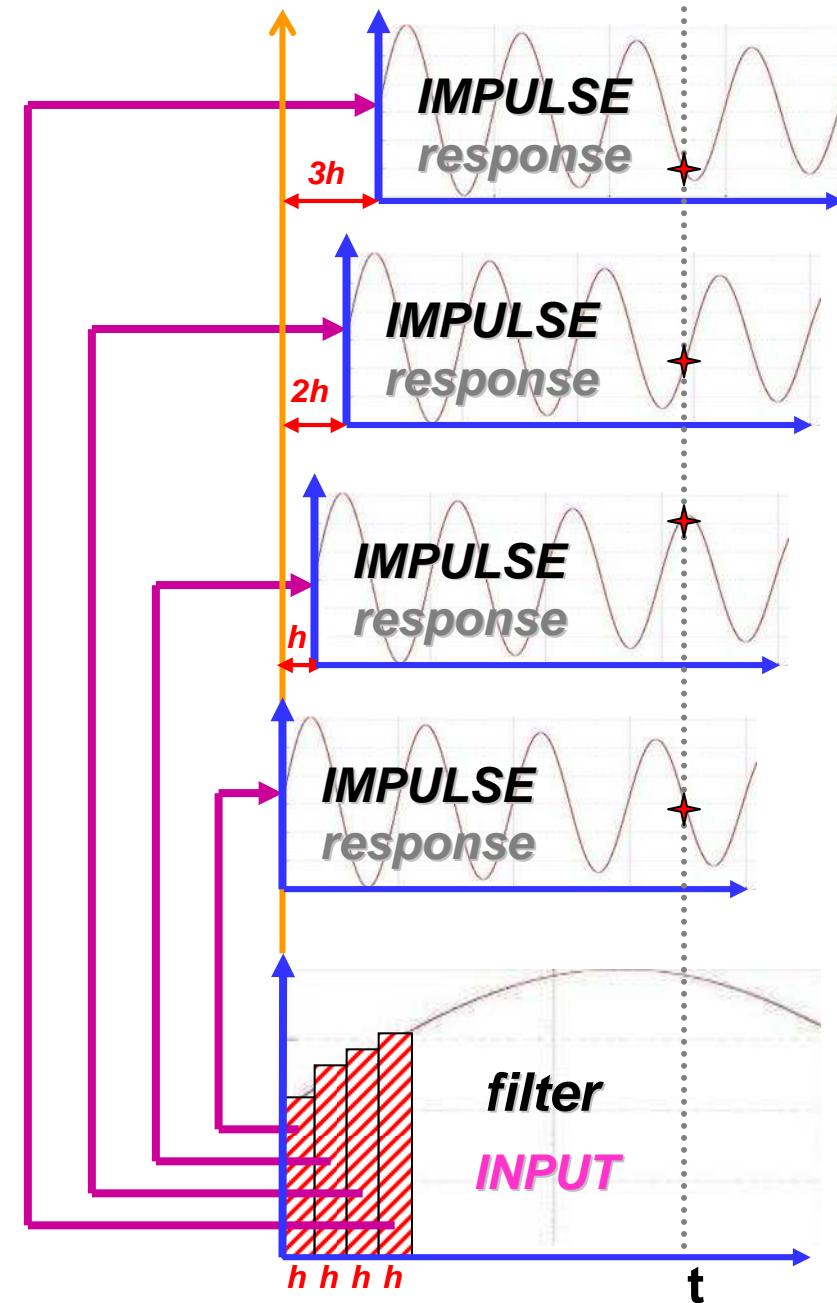


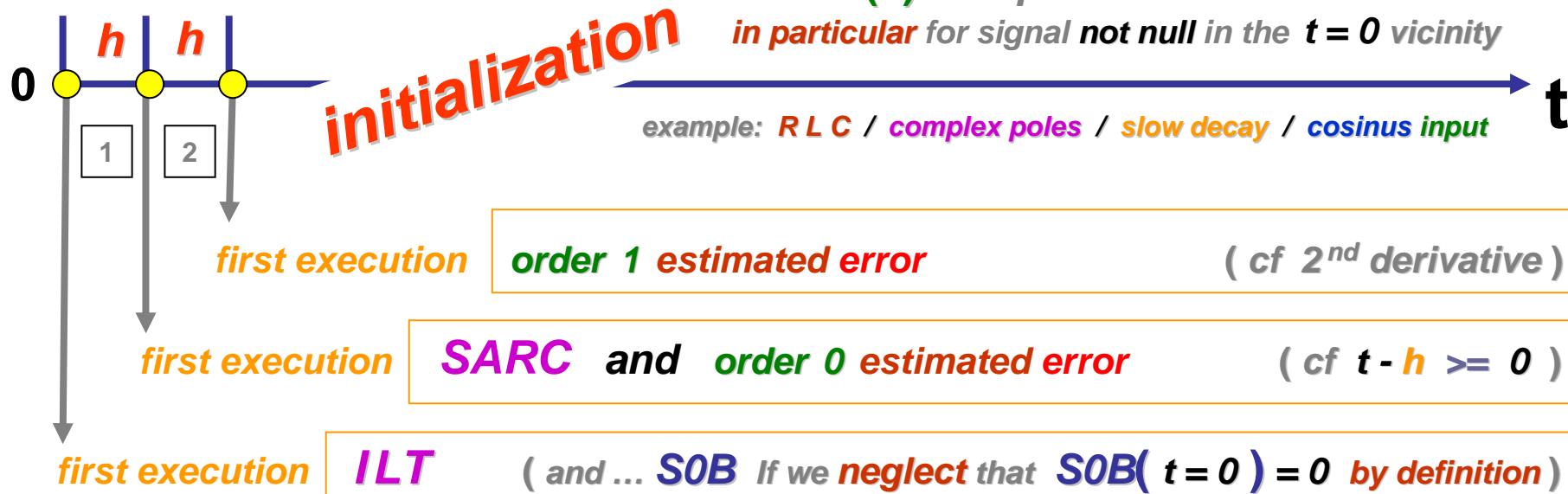
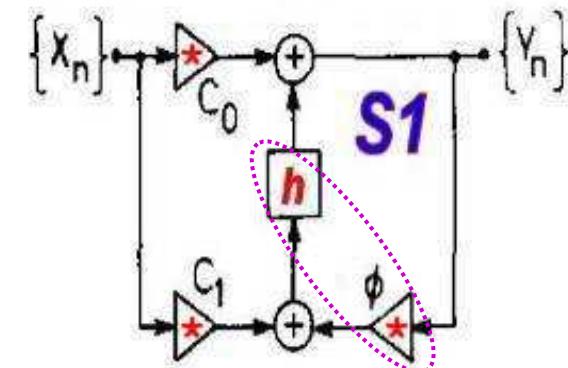
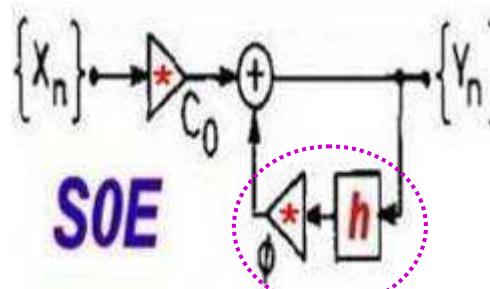
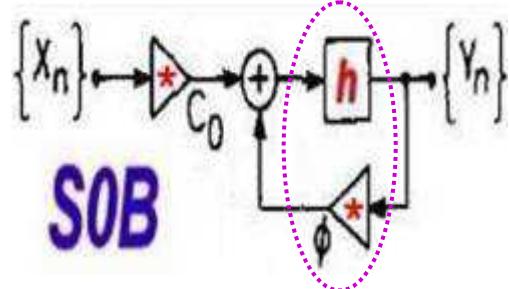
INITIALIZATION

a single **missing**
misplaced
 or **aberrant**

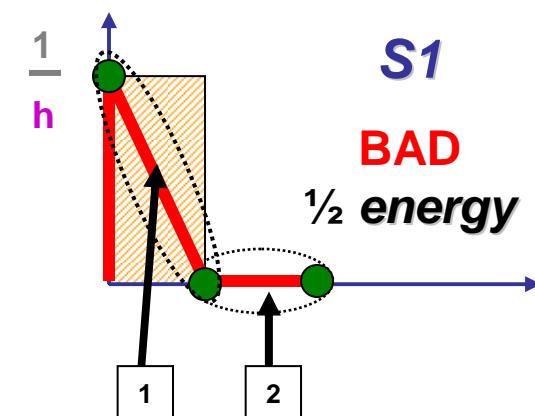
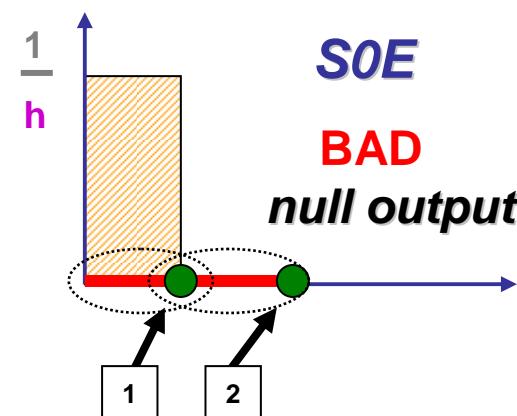
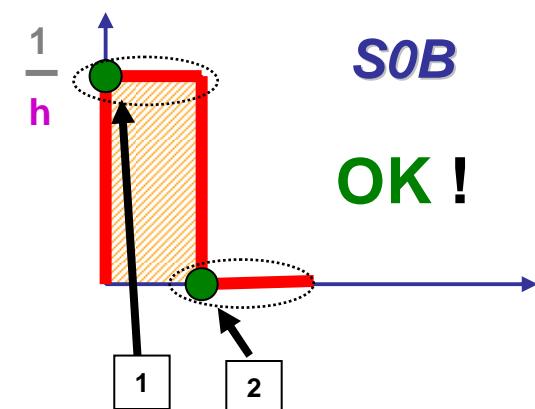
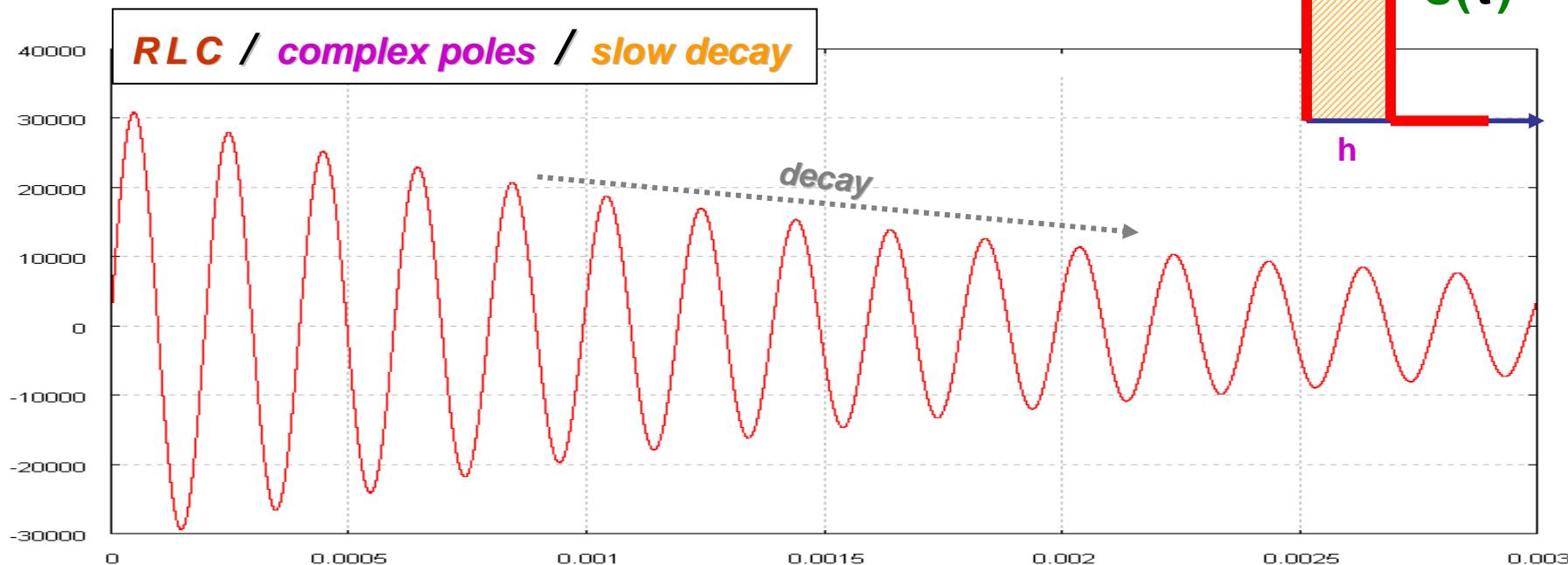
input sample response

may result in a significant **ERROR**
 over a long interval





impulsional response

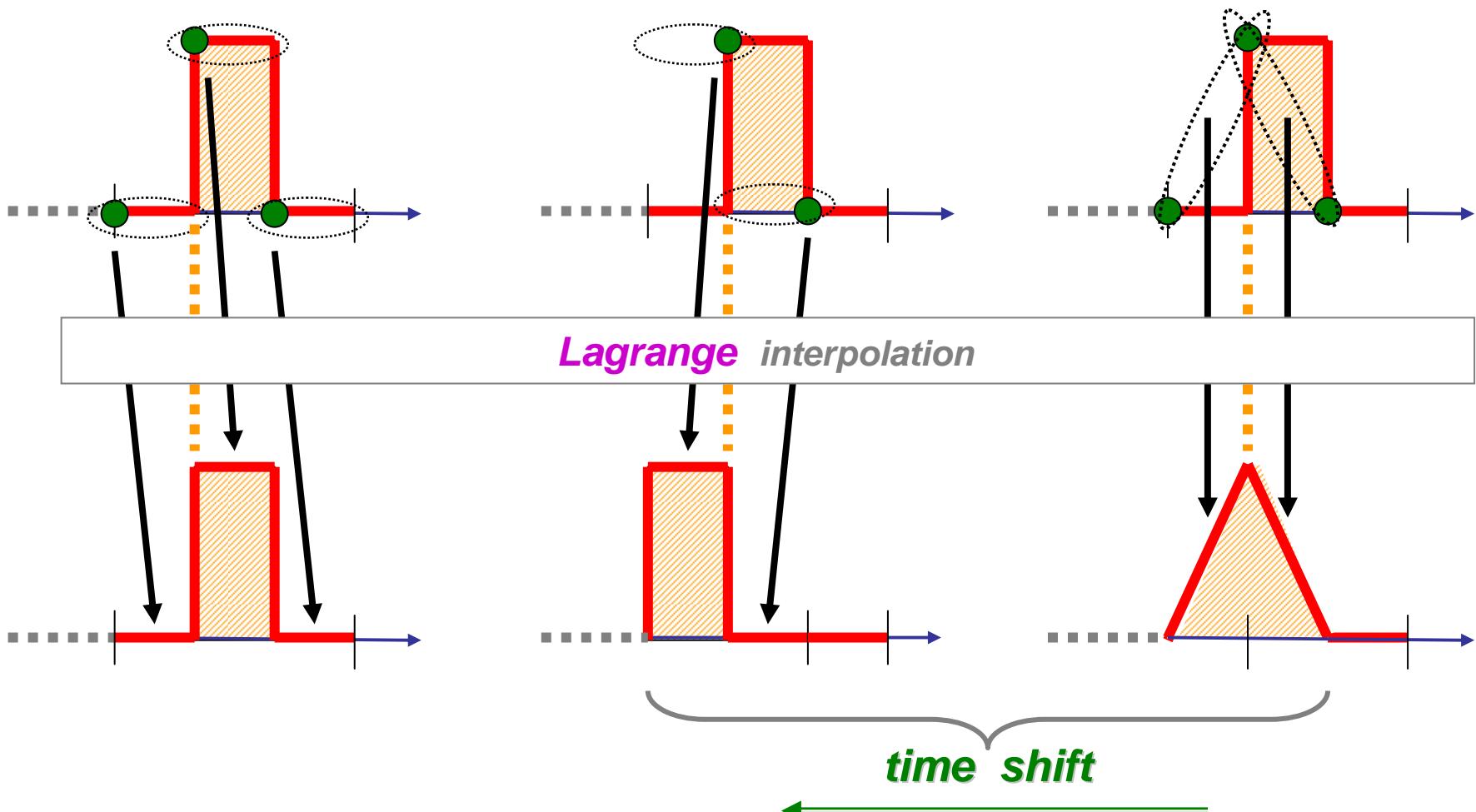


S0B

S0E

S1

input signal $U(t)$



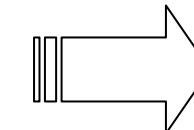
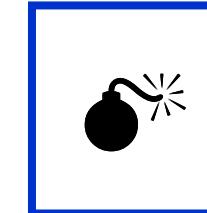
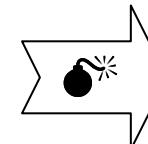
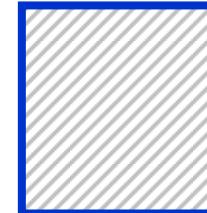
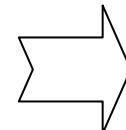
BUT

order 1 is better than **order 0**

general purpose simulators cascading sub-circuits

cannot tolerate initialization missing **output samples**

U(0)



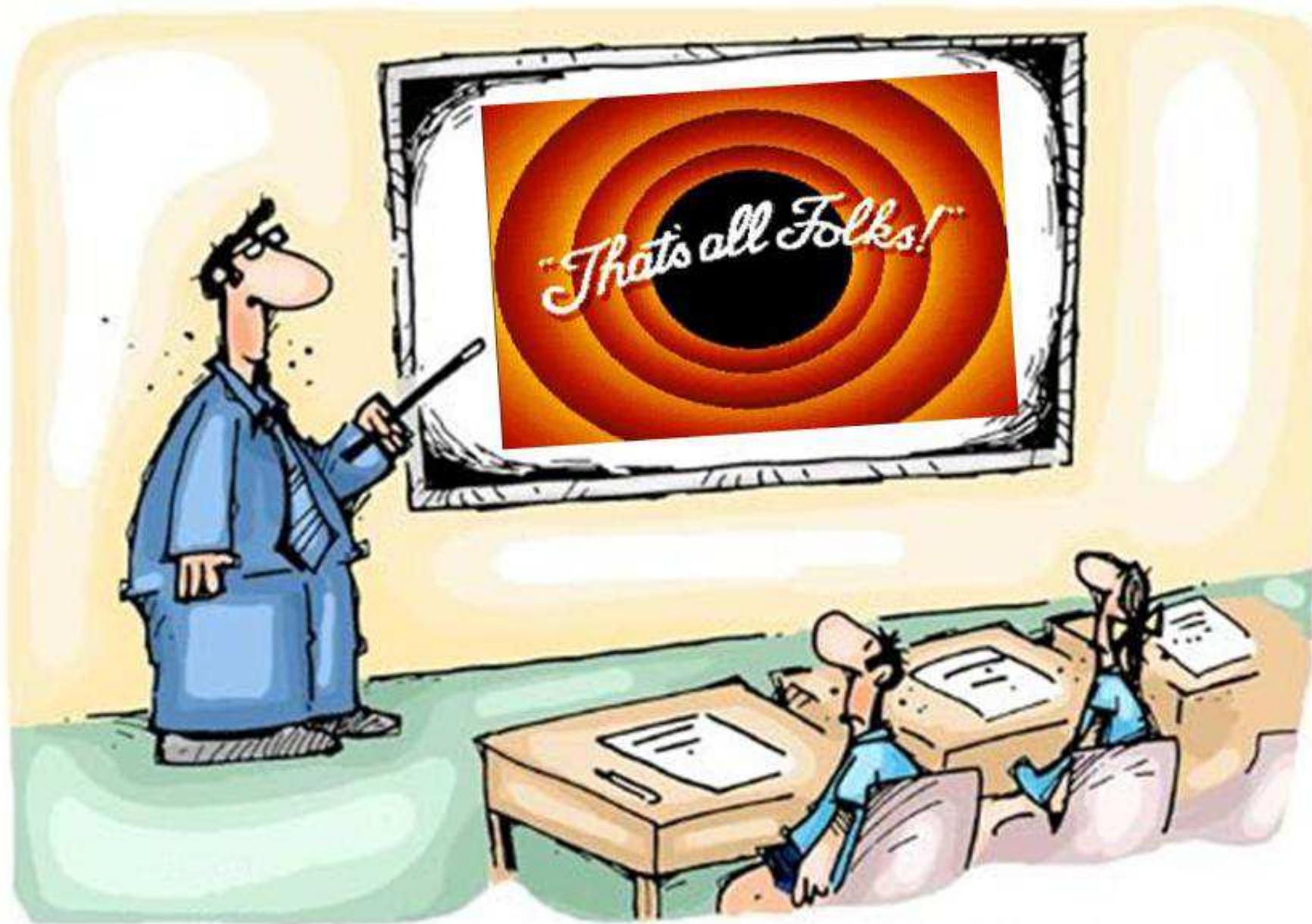
So, the



is more likely to be

S0B

similar to **ILT**





*Wait a minute ,
we have not talked
about
PERCUSSIONS*

CHARGE conservation

$$Q = C V$$

$$Q = C' V'$$

$$V' = \frac{C}{C'} V$$

FLUX conservation

$$\phi = L i$$

$$\phi = L' i'$$

$$i' = \frac{L}{L'} i$$

percussion C

$$C' = \frac{1}{2} C$$

$$V' = \frac{2}{1} V$$

L C components \Rightarrow "state" variables

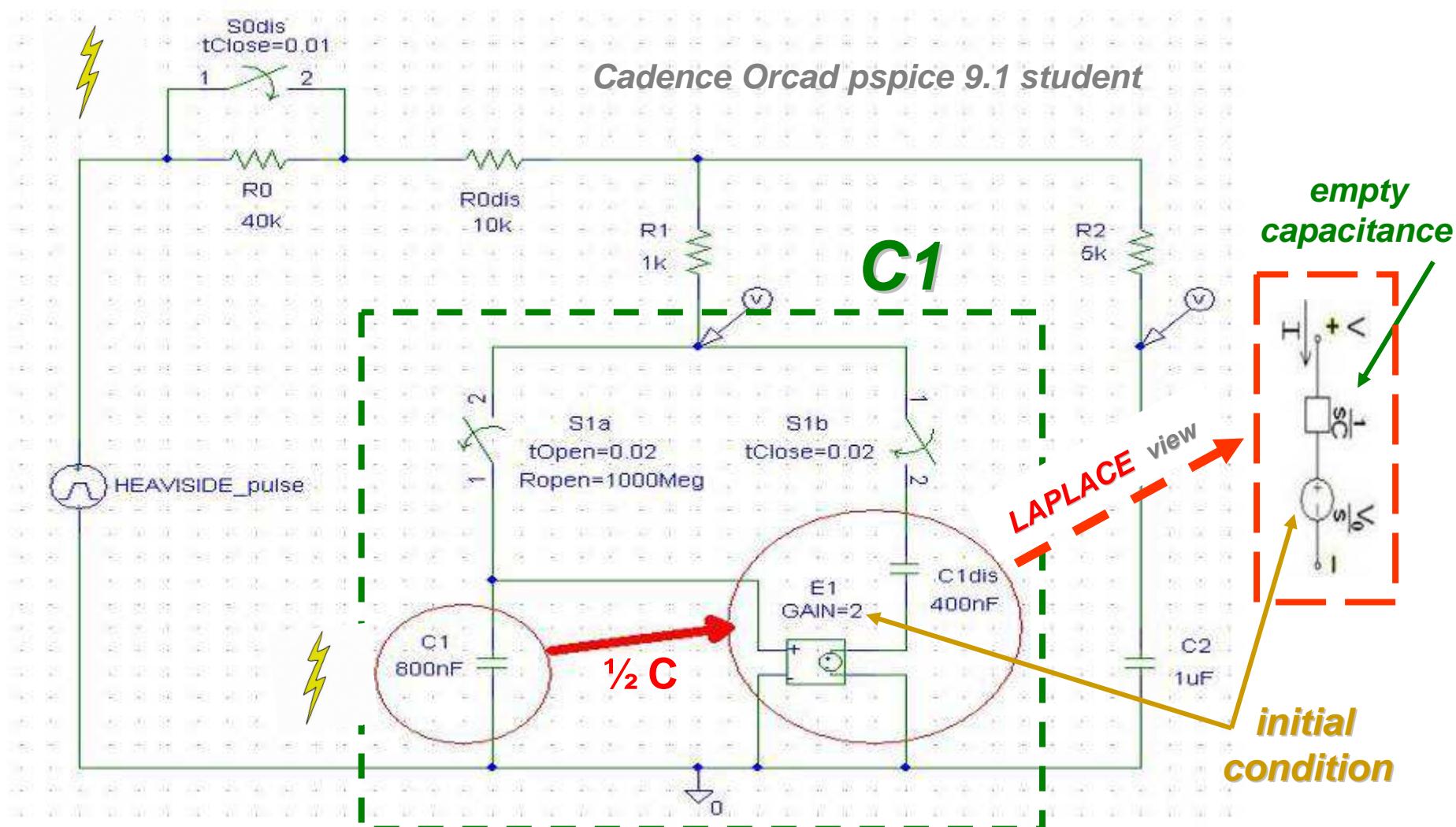
$$L' = \frac{1}{2} L$$

$$i' = \frac{2}{1} i$$

percussion L

PERCUSSION

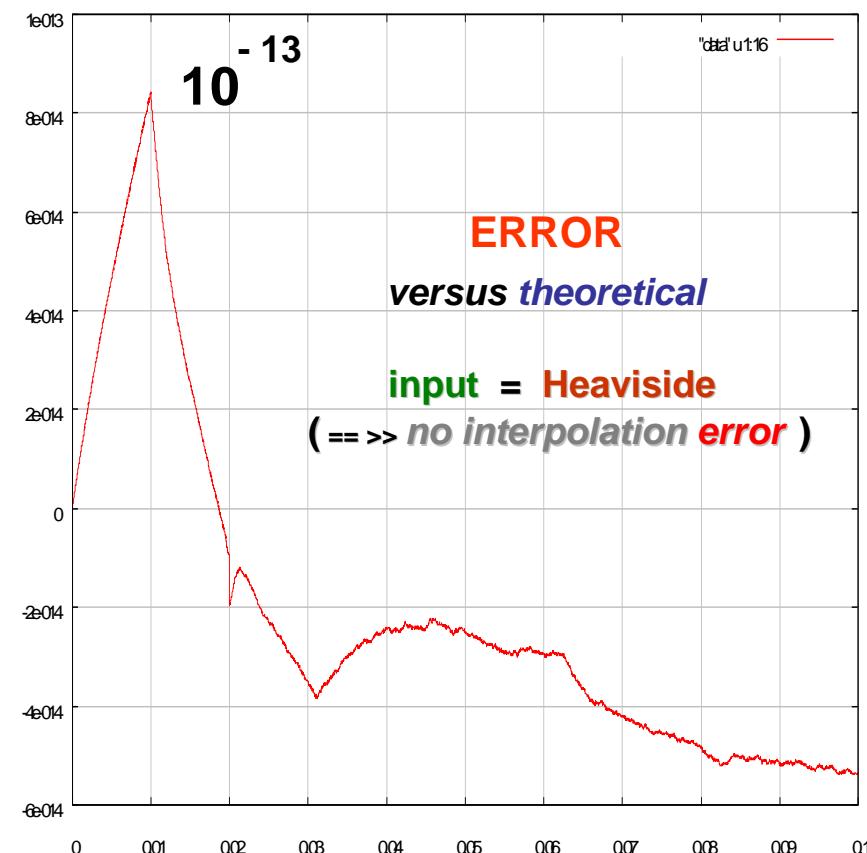
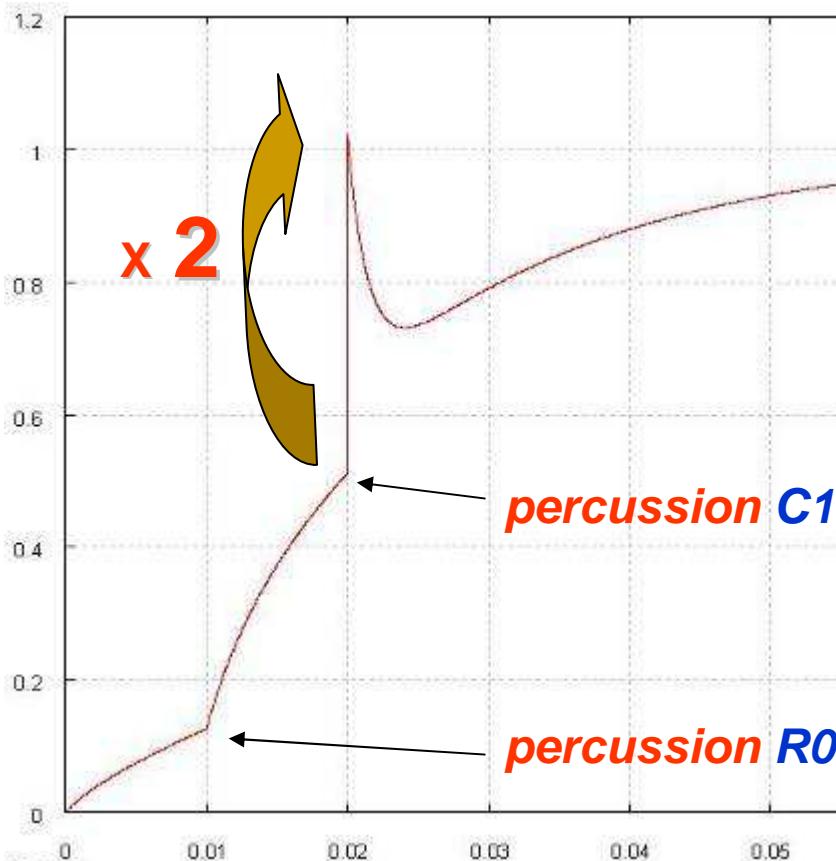
at $t = 0.01$ $R_0 = 50 \text{ Kohm} \Rightarrow 10 \text{ Kohm}$
 at $t = 0.02$ $C_1 = 800 \text{ nF} \Rightarrow 400 \text{ nF}$



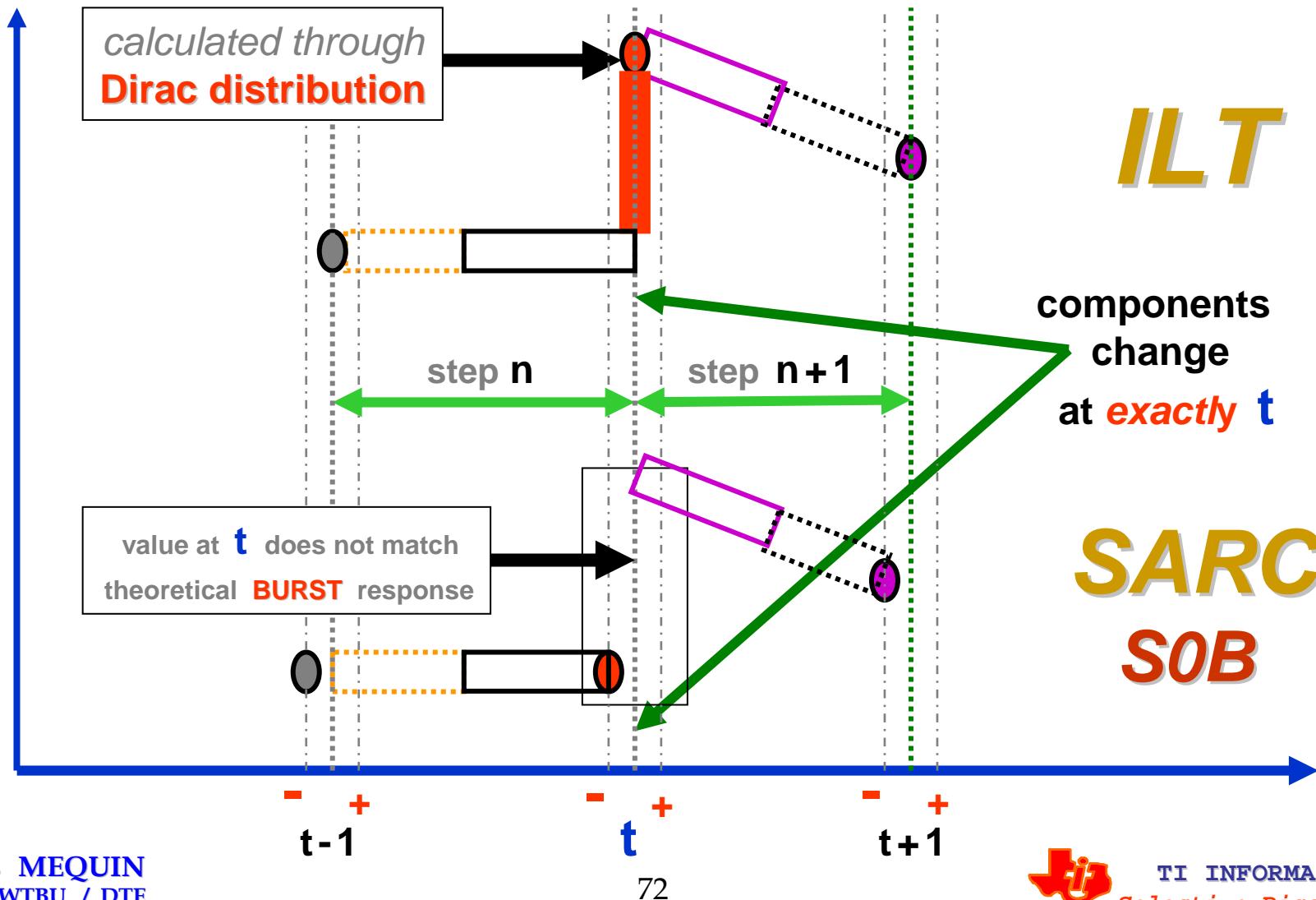
In fact, at least for network not “degenerate”, it is trivial !!!

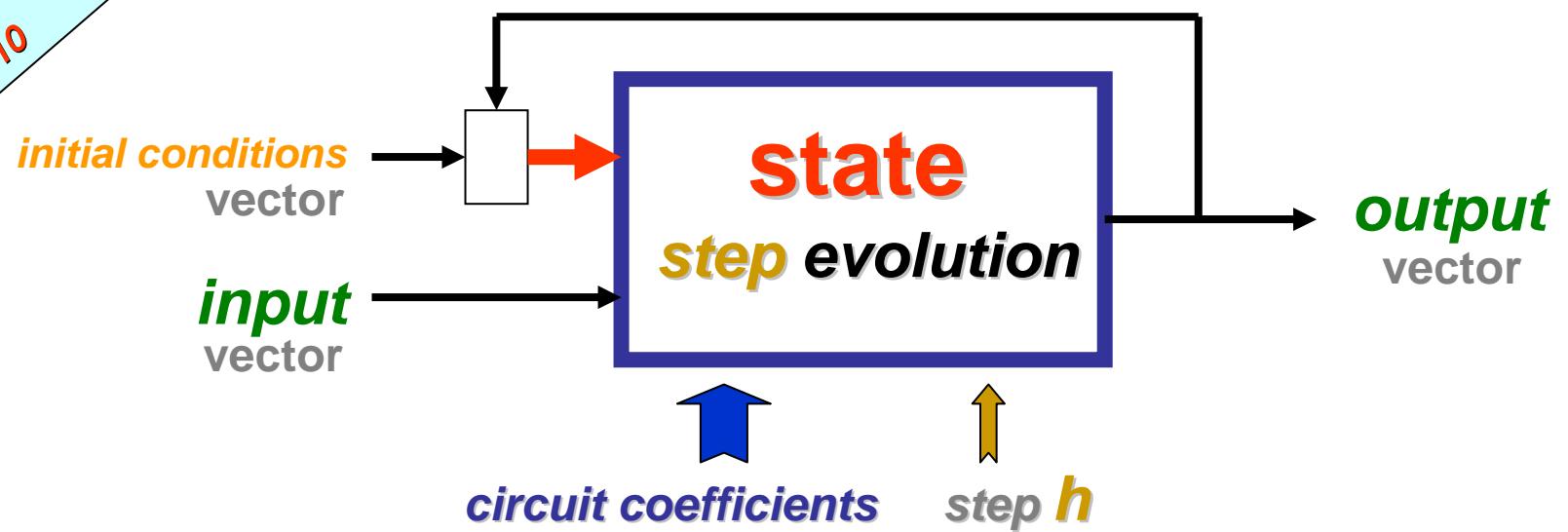
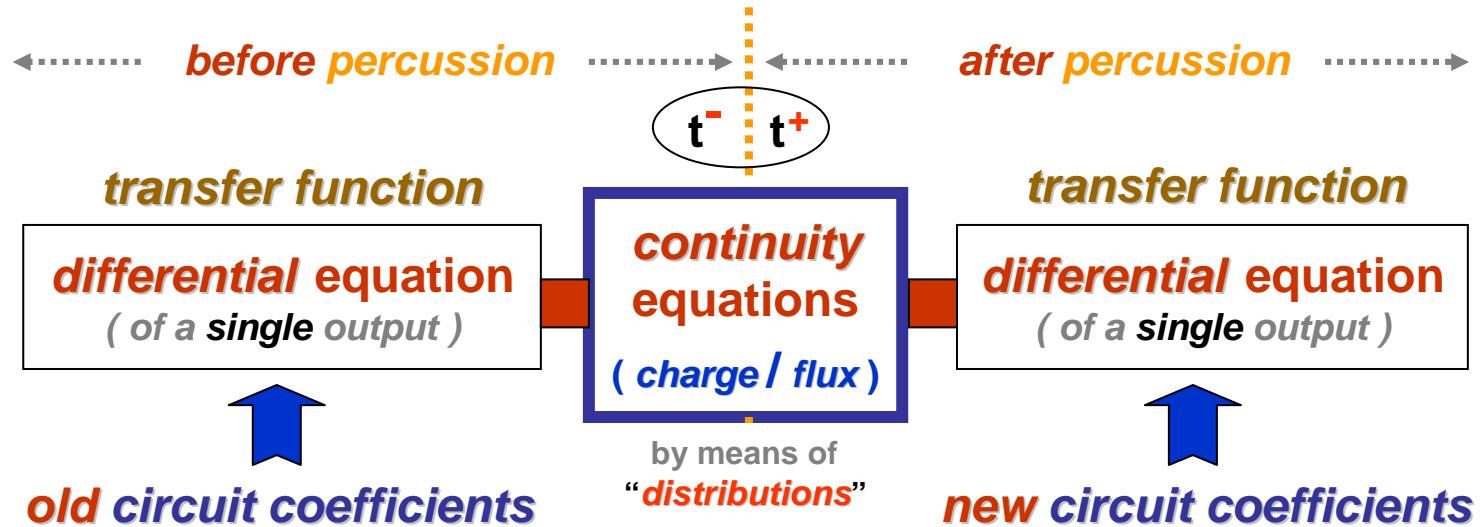
At the correct instant, update the percutred component values (R_0 or C_1)

and for the capacitance C_1 “double” its STATE VARIABLE voltage



SARC versus ILT percussion





QUESTIONS ?



Component PERCUSSION

of

***Linear Time Invariant
analog circuit***





Digital and Analog simulations are the 2 faces of a same coin

In the **digital world** architects require (through **ESL**) **high level functional** simulations to perform numerous “**what if**” and **analysis** in order to challenge the concept and to **increase confidence** about its validity

Unfortunately in the **analog world** architects do have much less tools and methodologies available to perform equivalent **high level functional** simulations

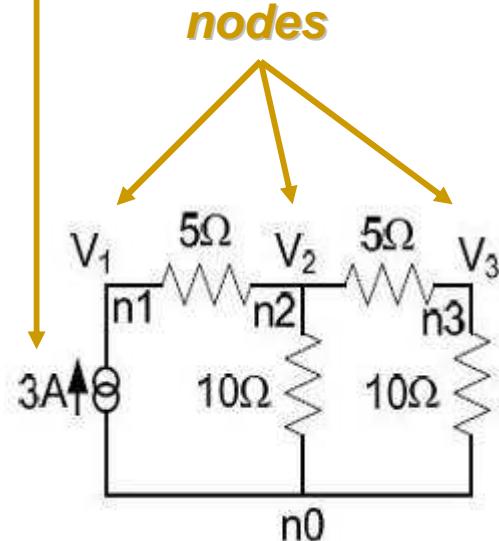
In many cases, **SPICE** turns to be **the only practical approach**

This is unfortunate because **SPICE** is so **low level** that its **runtime** quickly becomes a bridle to the architect ambition

Less experiments, Less tested alternatives, Less confidence ...

current source

(“**Norton**” equivalent for **Voltage**)



“NODAL” equations based on KCL

(Kirchhoff Current Law $\Sigma I = 0$)

$$[G] * [V] = [I]$$

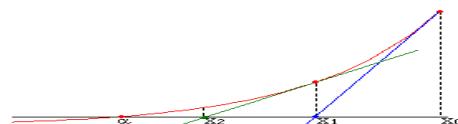
conductance

$$\begin{bmatrix} 0.2 & -0.2 & 0 \\ -0.2 & 0.5 & -0.2 \\ 0 & -0.2 & 0.3 \end{bmatrix} \begin{bmatrix} V_1 \\ V_2 \\ V_3 \end{bmatrix} = \begin{bmatrix} 3 \\ 0 \\ 0 \end{bmatrix}$$

“Gaussian” elimination
to make it “triangular”
then substitutions ...

$$\begin{bmatrix} 0.2 & -0.2 & 0 \\ 0 & 0.3 & -0.2 \\ 0 & 0 & 0.25 \end{bmatrix} \begin{bmatrix} V_1 \\ V_2 \\ V_3 \end{bmatrix} = \begin{bmatrix} 3 \\ 3 \\ 3 \end{bmatrix}$$

NON-LINEAR COMPONENT behavior is approximated by doing **iterations** (**convergence**)
based on **Newton-Raphson** method (improvement of an equation root “**guess**”)



$$f(x) \cong f(x_0) + f'(x_0)(x - x_0)$$

=> **recurrence**

$$x_{n+1} = x_n - \frac{f(x_n)}{f'(x_n)}$$

NUMERICAL INTEGRATION (for **linear** inductances and capacitances) is approximated according to **3 user-selectable methods** : **Trapezoidal** , **backward-Euler** , **Gear**

NUMERICAL INTEGRATION

(for **linear** inductances **and** capacitances)

Why several different methods ?

Depending on the “expected” waveform, one method may have an advantage over another

What makes them strong or weak ?

Accuracy How much error (**Local Truncation Error**) each method introduces ?

Stability Does the error accumulate (**that's bad**) or decrease to zero (**that's good**) over time ?

Each method has its own strengths and weaknesses

Trapezoidal

$$x_{n+1} = x_n + h \cdot \frac{(\dot{x}_{n+1} + \dot{x}_n)}{2}$$

Better accuracy than stability
may oscillate in some cases

Backward Euler

$$x_{n+1} = x_n + h \cdot \dot{x}_{n+1}$$

Average accuracy and stability

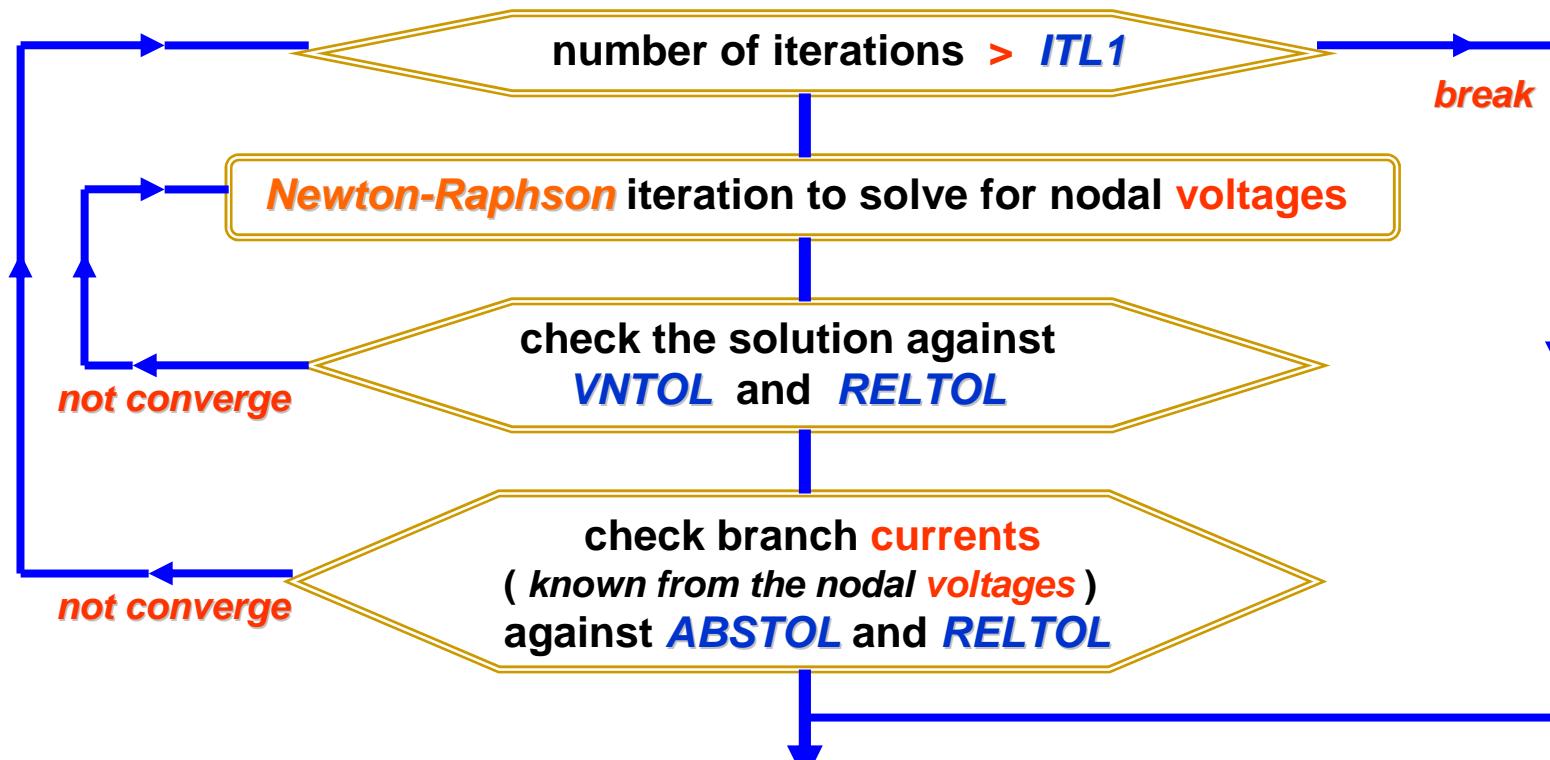
Gear-2

$$x_{n+1} = \frac{4}{3} \cdot x_n - \frac{1}{3} \cdot x_{n-1} + \frac{2}{3} \cdot h \cdot \dot{x}_{n+1}$$

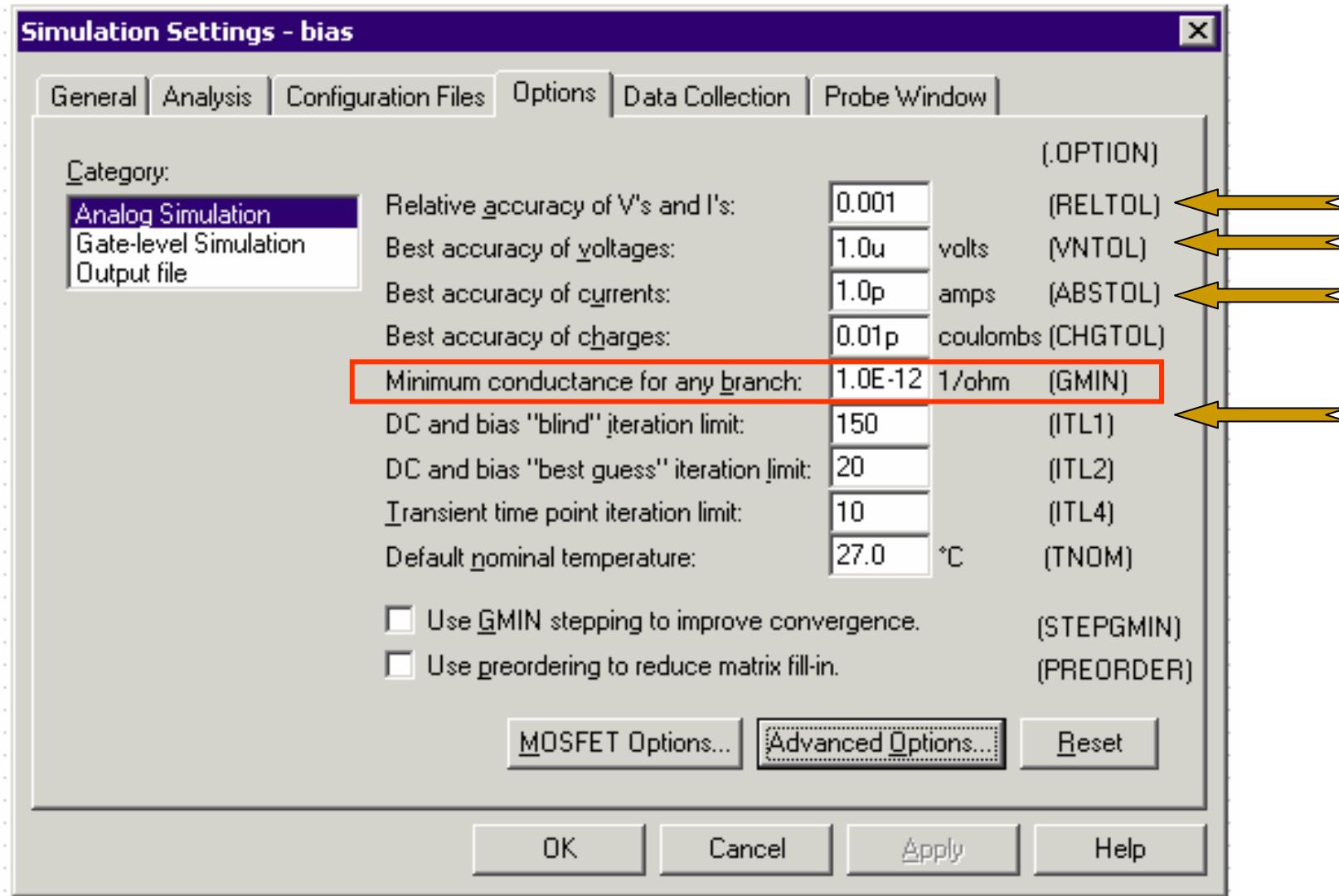
Better stability than accuracy
Most stable method

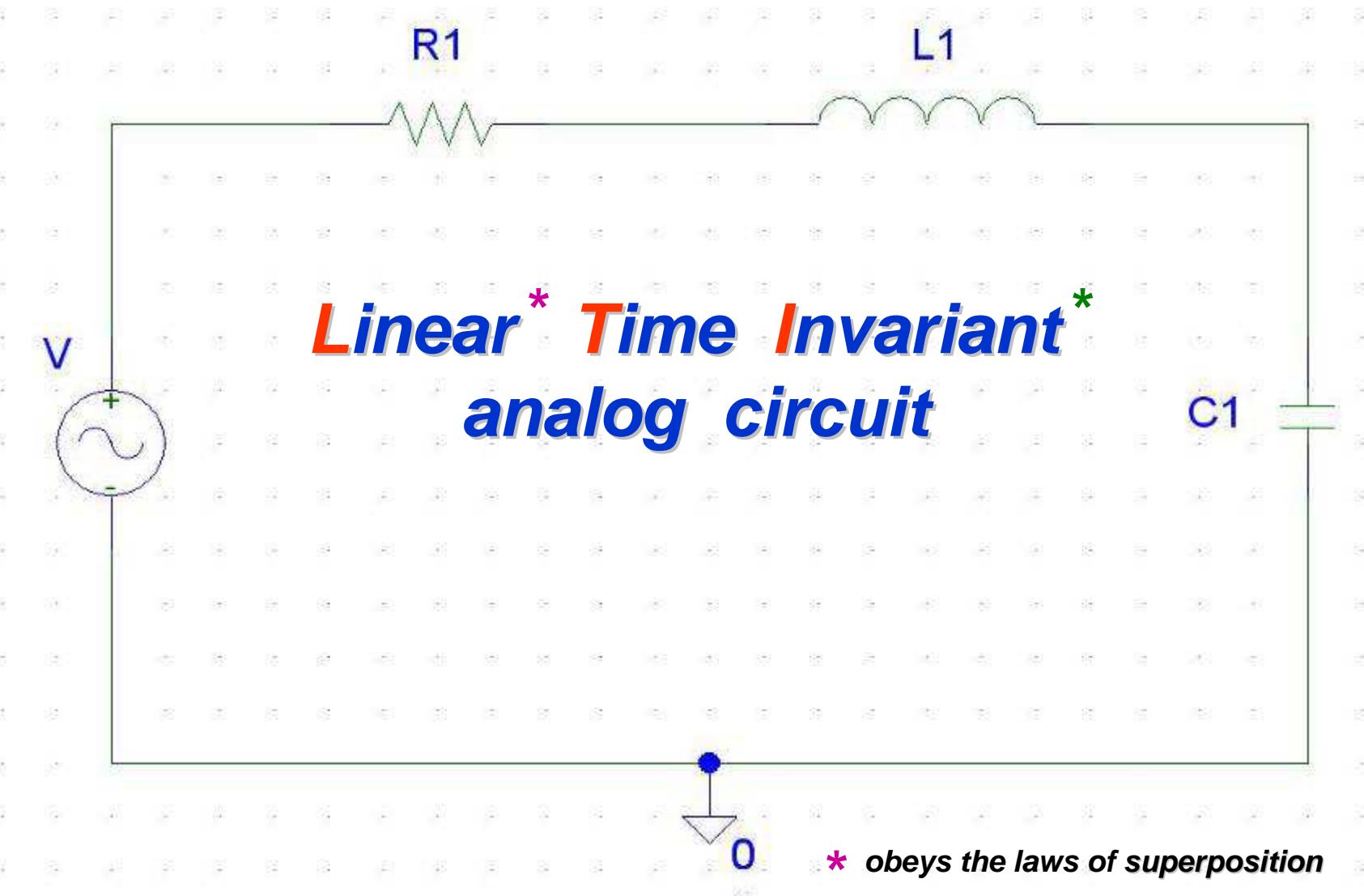
SPICE “convergence” control variables

| | |
|---------------|--|
| ITL1 | number of iterations allowed for a bias point calculation |
| VNTOL | node voltage tolerance |
| RELTOL | fractional tolerance of voltages and currents |
| ABSTOL | current branch tolerance |



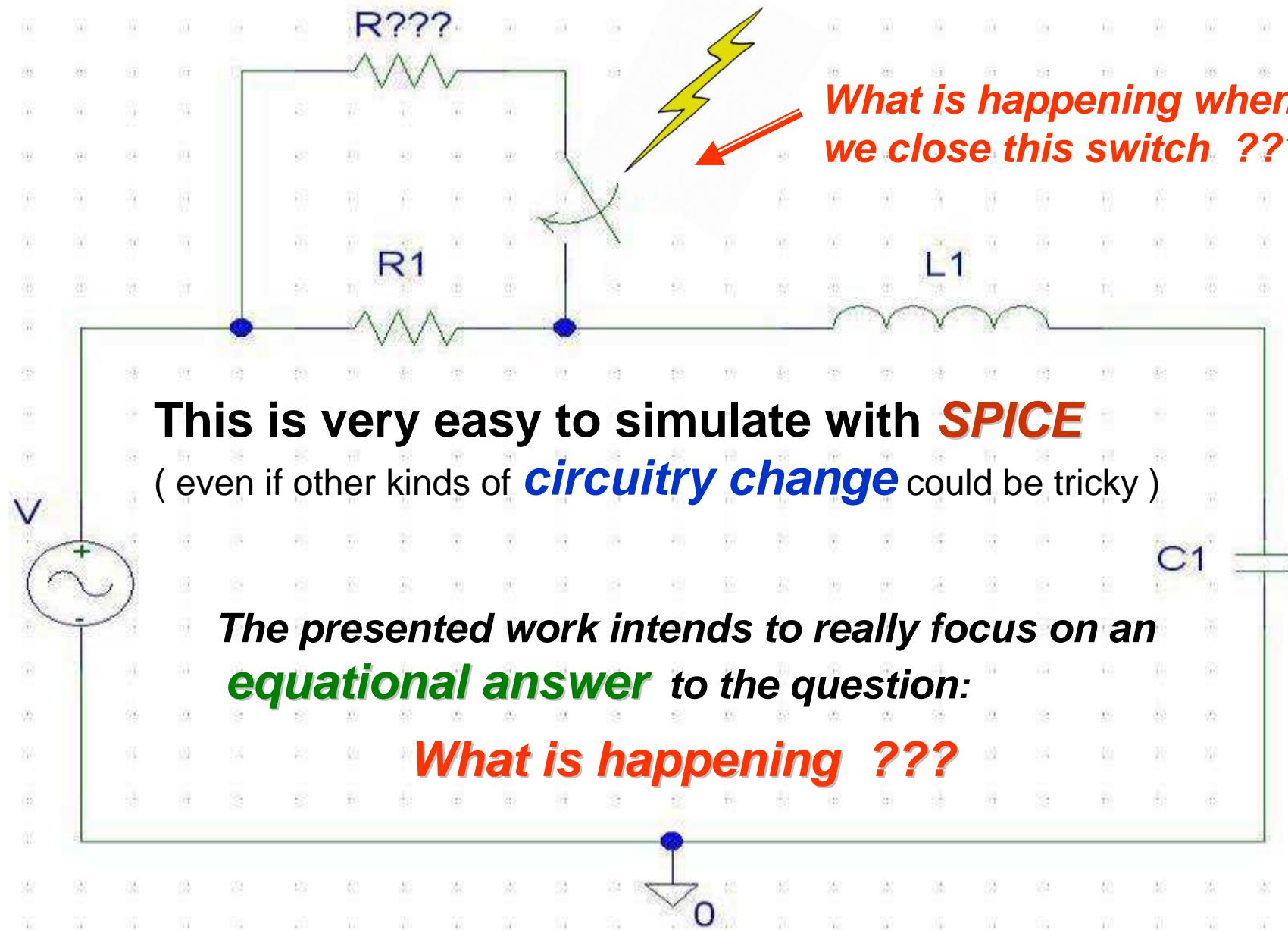
SPICE





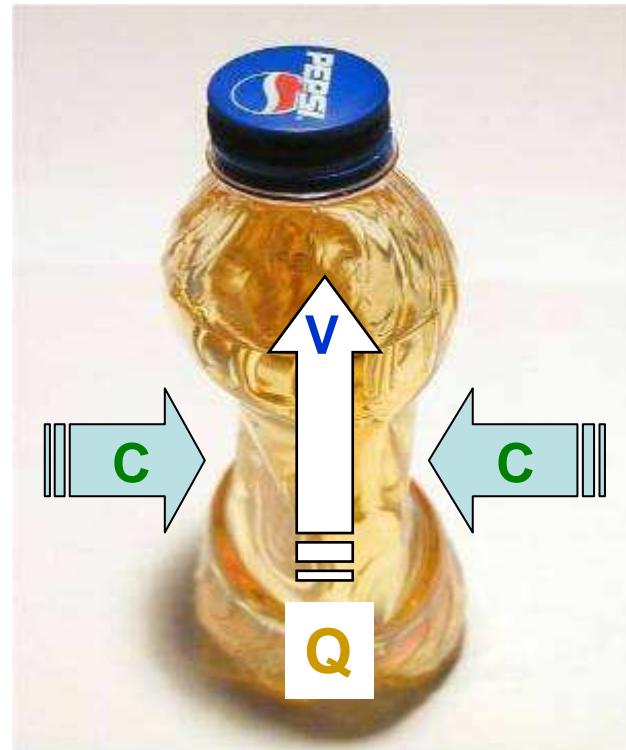
* obeys the laws of superposition

* does not change its behavior over time
(i.e. constant intrinsic parameters)



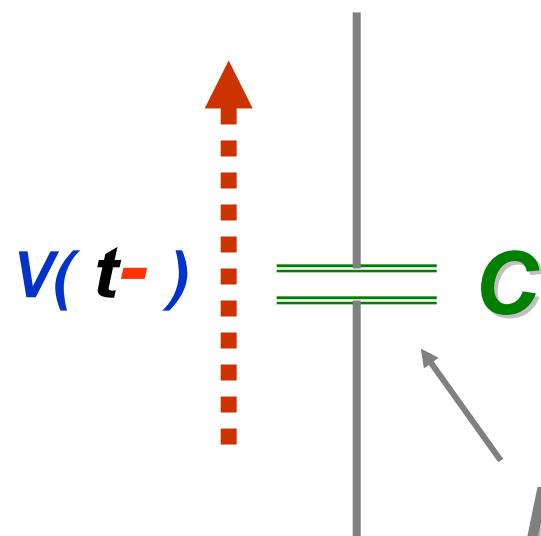
other example:

**capacitance
change ?**



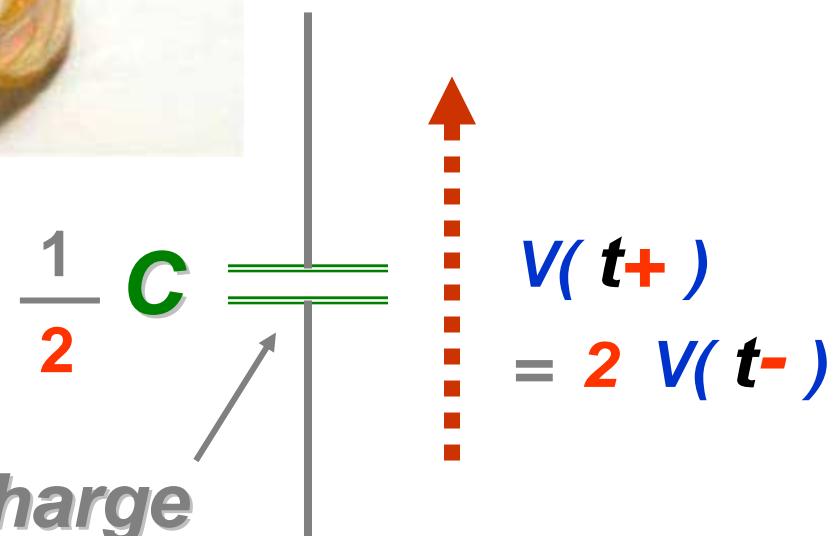
$$Q = C V$$

$$V = \frac{Q}{C}$$

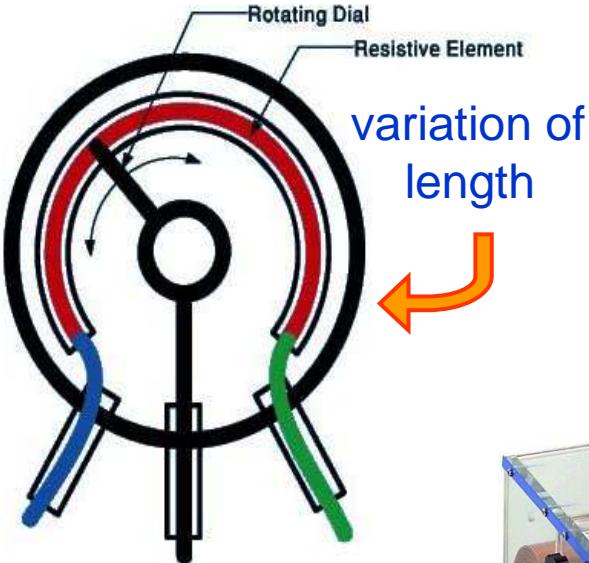


preserved charge

$$Q(t-)$$

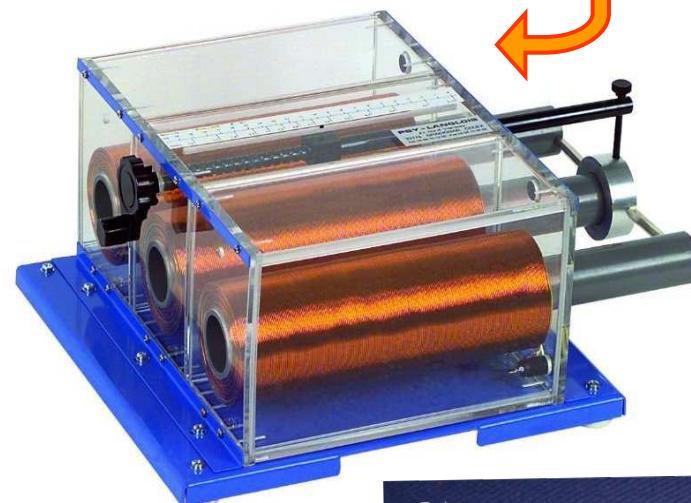


$$V(t+) = 2 V(t-)$$

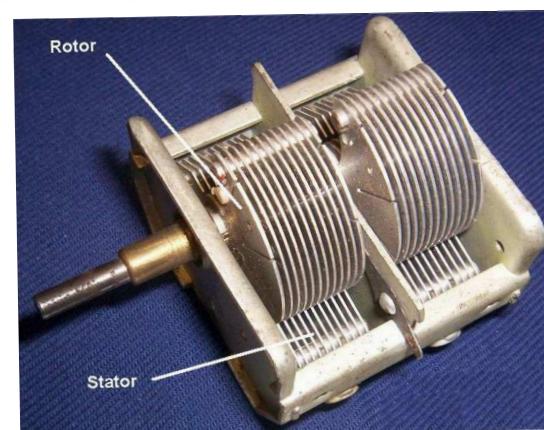


R

variation of
permeability (μ)



L



C

percussion

Definition:



- 1) **musical instruments that are hit:** the group of musical instruments that produce sound by being struck, including drums and cymbals, or the section of the orchestra playing such instruments
- 2) **act of detonating percussion cap:** the striking or detonating of a percussion cap in a firearm
- 3) **impact:** the impact of one object striking another, or the noise or
shock created when two objects hit each other
- 4) **medicine tapping of body:** examination of part of a patient's body by tapping with the fingers to assess the presence of fluid, the enlargement of organs, or the solidification of normally hollow parts

Wait a minute, is there anything **wrong** with the title of this presentation ???

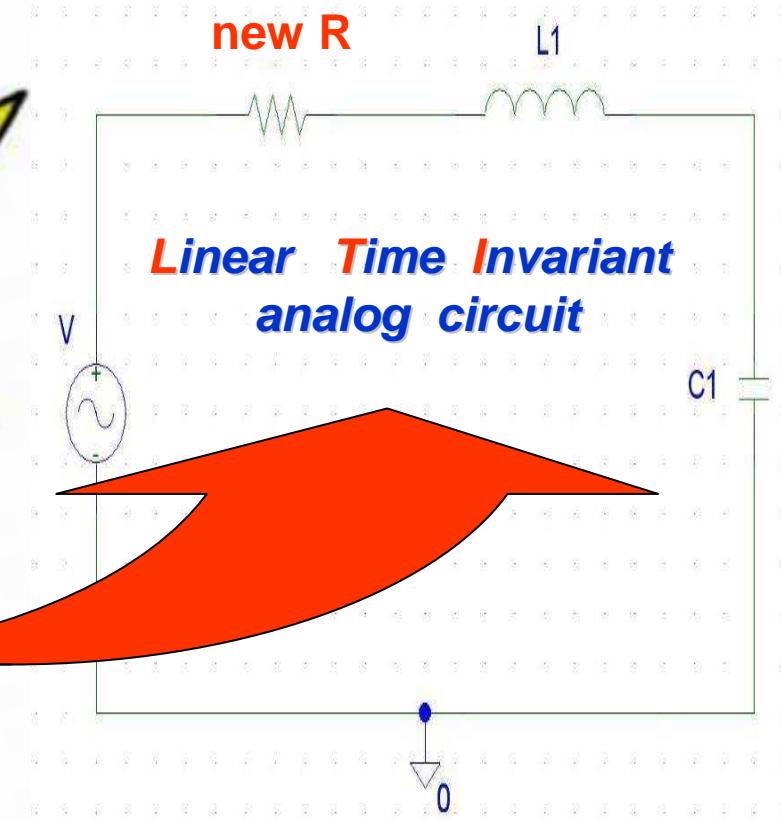
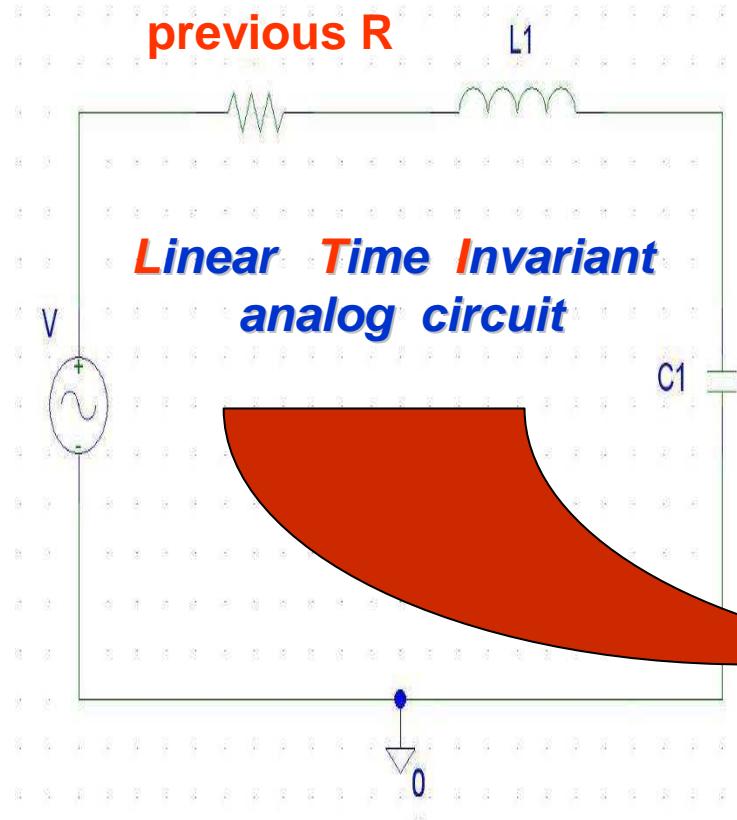
Component **PERCUSSION** of **Linear^{*} Time Invariant^{*}** **analog circuit**

- * obeys *the laws of superposition*
- * does not change its behavior over time
(i.e. **constant intrinsic parameters**)

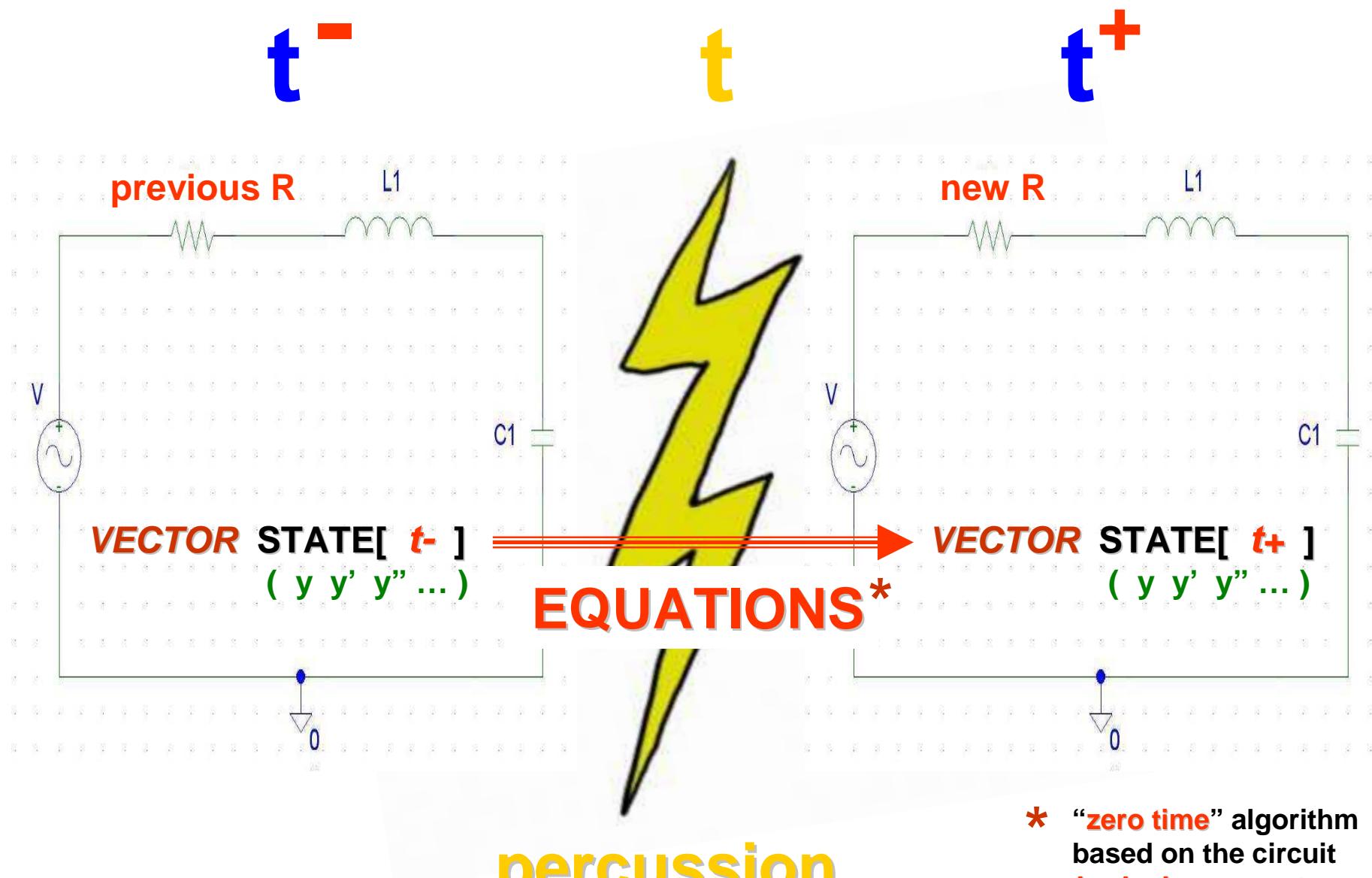
t^-

t

t^+



percussion



* “zero time” algorithm
 based on the circuit
intrinsic parameters
 not a convergence !

Moreover,

the circuit "*initialization*"
is just a *particular case* of such
instantaneous state change

circuit
“*at rest*”

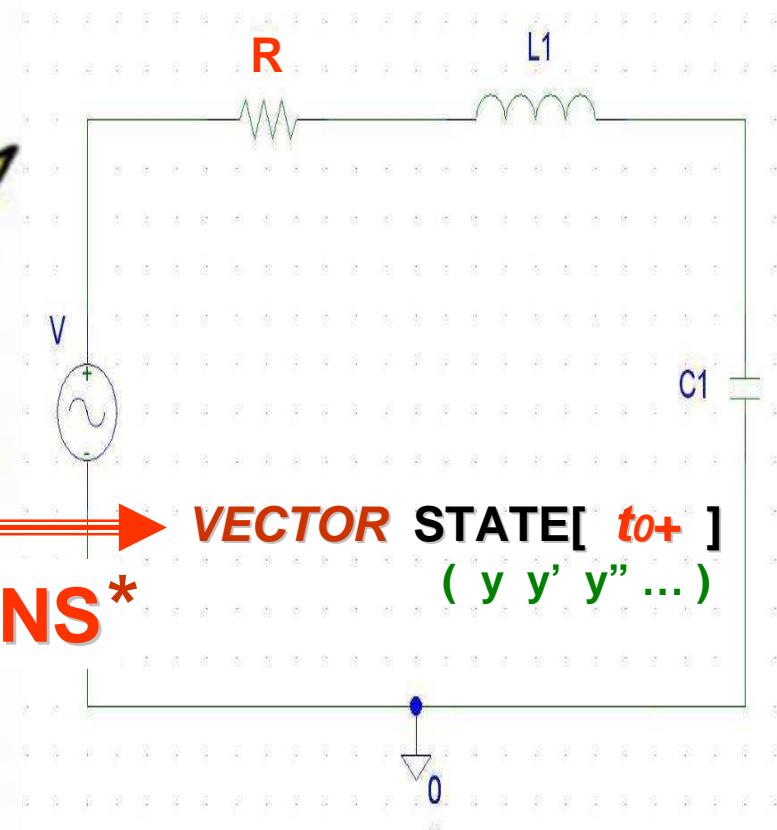
percussion

t_0

t_0^+



EQUATIONS*



* “zero time” algorithm
based on the circuit
intrinsic parameters
not a convergence !

In summary, a **PERCUSSION** can be



- a “**sharp**” change of the **INPUT** (including **INITIALIZATION**)
- an “**INSTANTANEOUS**” change of **R** , **L** or **C**

However,

the processing of **INPUT** and **COMPONENT PERCUSSIONS** being “*similar*”
but one modeling an **external shock**
and the other one an **internal structural change**

in this **Inverse Laplace Transform software package**

a “**sharp**” change of an **INPUT** is referenced as a **DISCONTINUITY**
but a modification of **R**, **L** or **C** is effectively called a **PERCUSSION**

In fact, we are playing with words



a **discontinuity** on the INPUT results in a **percussion** of the circuitry

a **percussion** of a COMPONENT results in a **discontinuity** on the OUTPUT

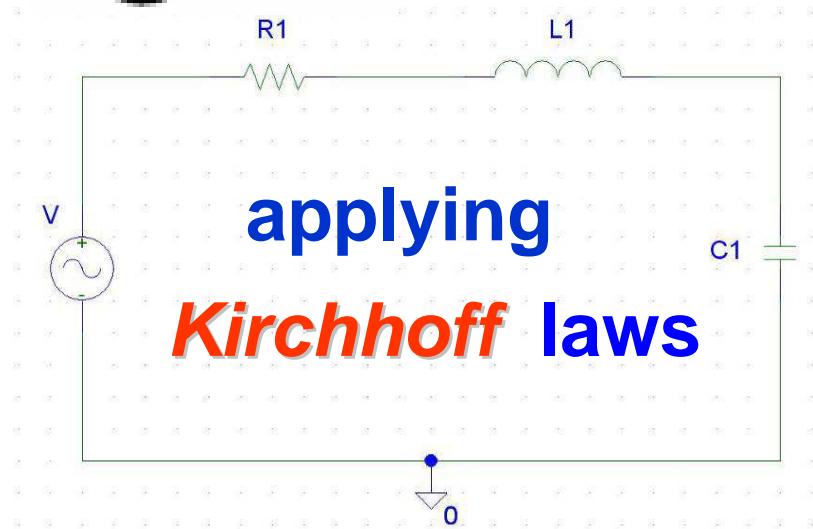


92



MAXIMA

*Symbolic
(Computer Algebra System)*



```
EQ : [
    vr1 = R1 * ir1 ,
    v11 = ZL1 * ill ,
    vc1 = ZC1 * ic1 ,
    ir1 = ill ,
    ir1 = ic1 ,
    v = vr1 + v11 + vc1
] ;
```

```
solve( eliminate( EQ , [ v11, vc1, ir1, ill, ic1 ] ) , vr1 ) ;
solve( eliminate( EQ , [ vr1, v11, vc1, ir1, ic1 ] ) , ill ) ;
solve( eliminate( EQ , [ vr1, v11, ir1, ill, ic1 ] ) , vc1 ) ;
```

$$[\mathbf{vr1} = \frac{\mathbf{R1} \mathbf{v}}{\mathbf{ZL1} + \mathbf{ZC1} + \mathbf{R1}}]$$

$$[\mathbf{ill} = \frac{\mathbf{v}}{\mathbf{ZL1} + \mathbf{ZC1} + \mathbf{R1}}]$$

$$[\mathbf{vc1} = \frac{\mathbf{v} \mathbf{ZC1}}{\mathbf{ZL1} + \mathbf{ZC1} + \mathbf{R1}}]$$

UNFORTUNATELY,



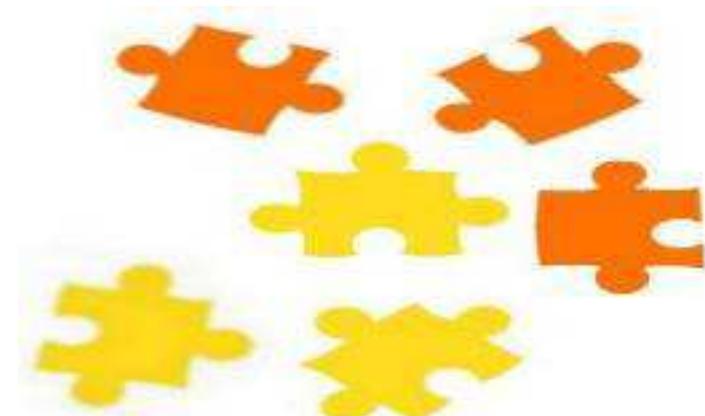
*Multiple approaches exist **to model** such kind of filters :*

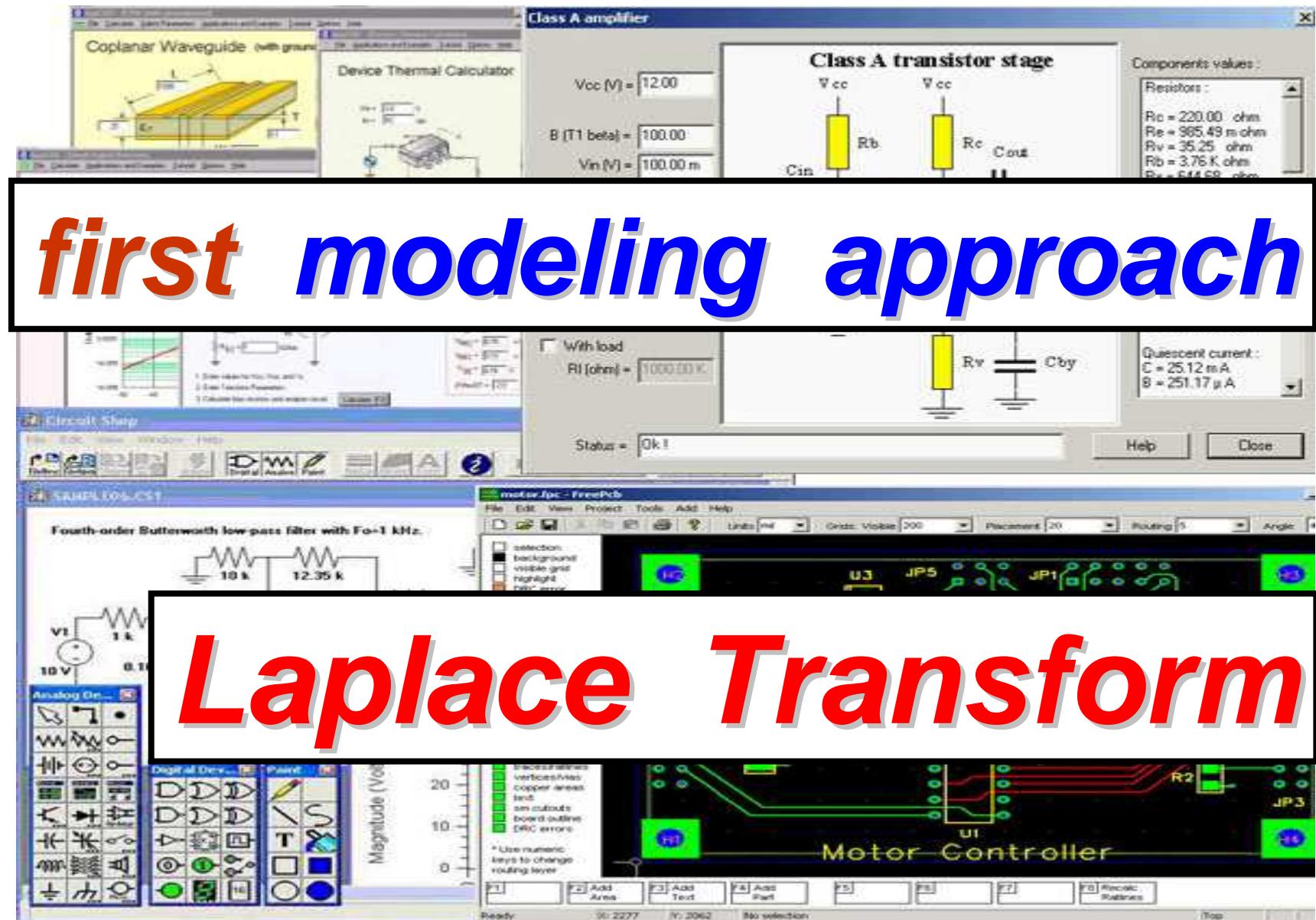
★ **Laplace Transform**

★ **State Model**

★ **Differential equations**

⊕ **Differential equations** by means of **Distributions**

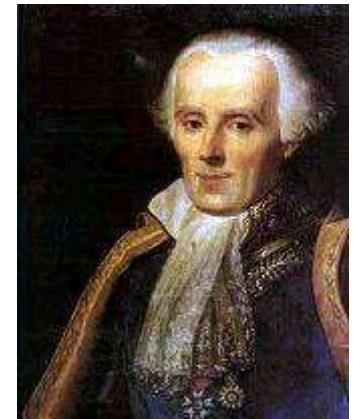




LAPLACE transform

$$L\{f(t)\} = F(s) = \int_0^{\infty} f(t) e^{-st} dt$$

$$s = \alpha + j\omega$$

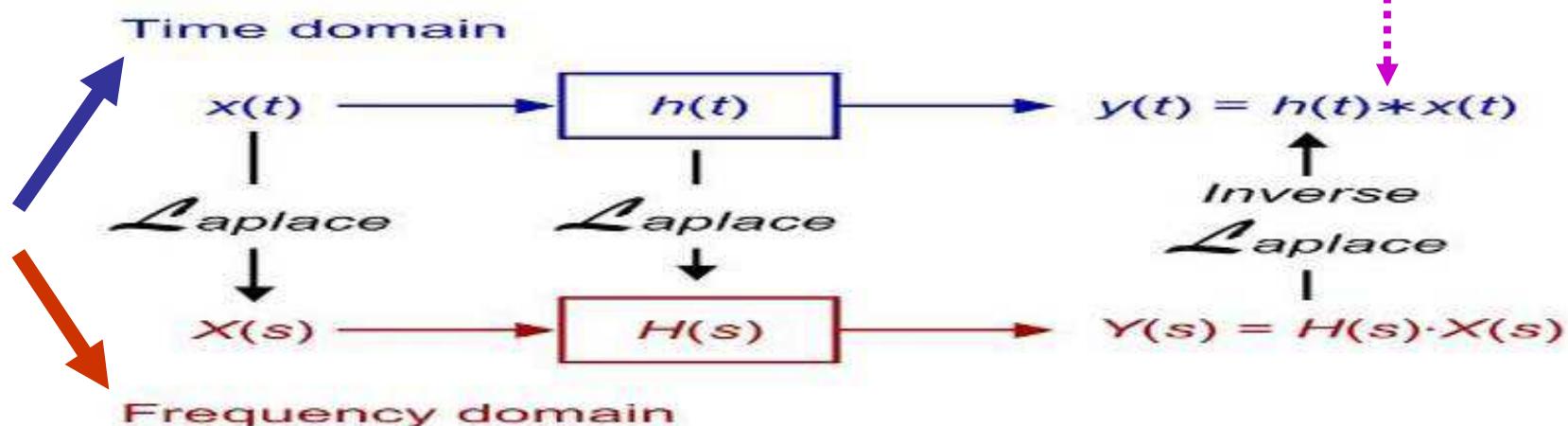


1749 - 1827

for causal energy bounded signals

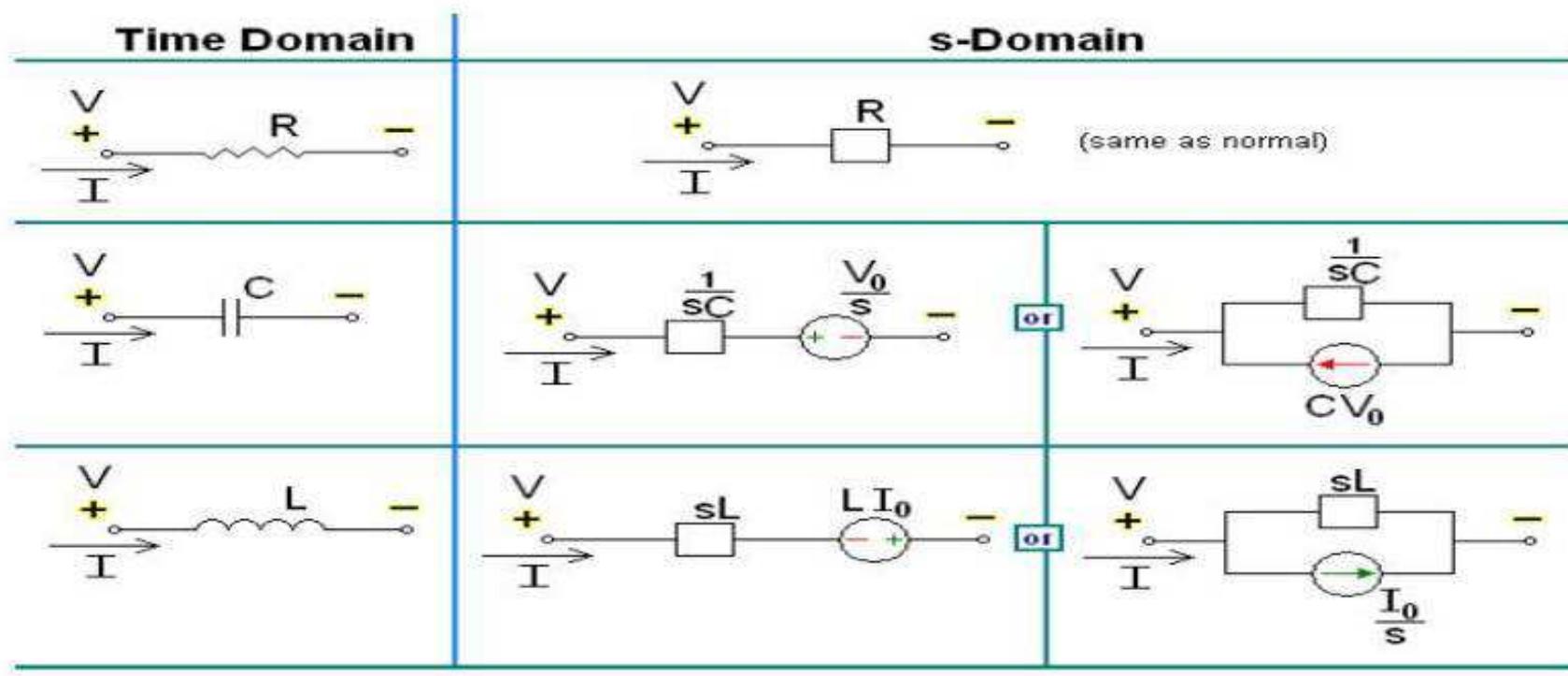
| | | |
|-------------------|---|-------------------------|
| Fourier | = | Laplace |
| $\mathcal{F}x(f)$ | = | $\mathcal{L}x(j2\pi f)$ |

convolution product



Laplace

“*s*” (or “*p*”) transform



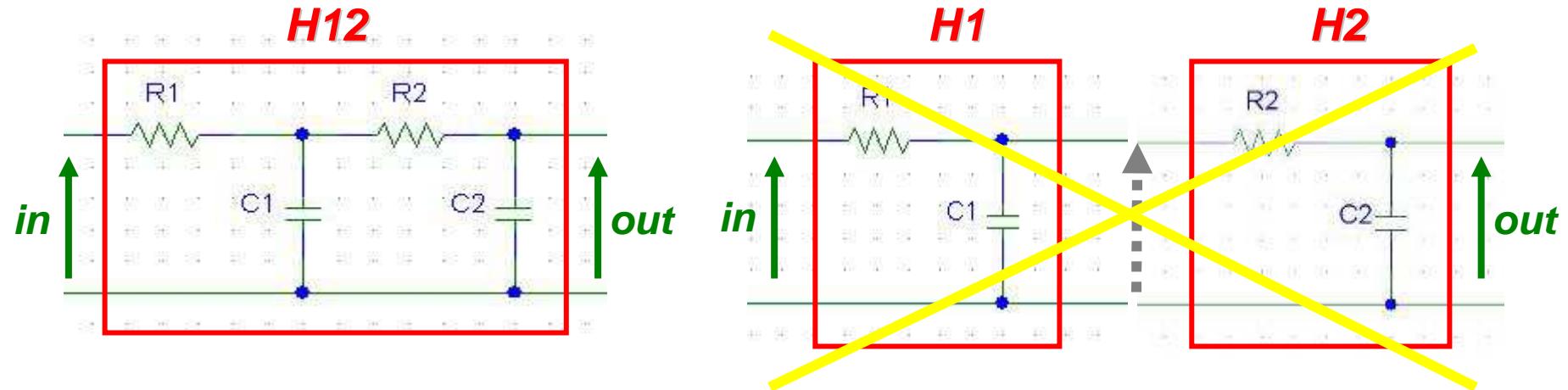
terminology

resistance R
conductance G

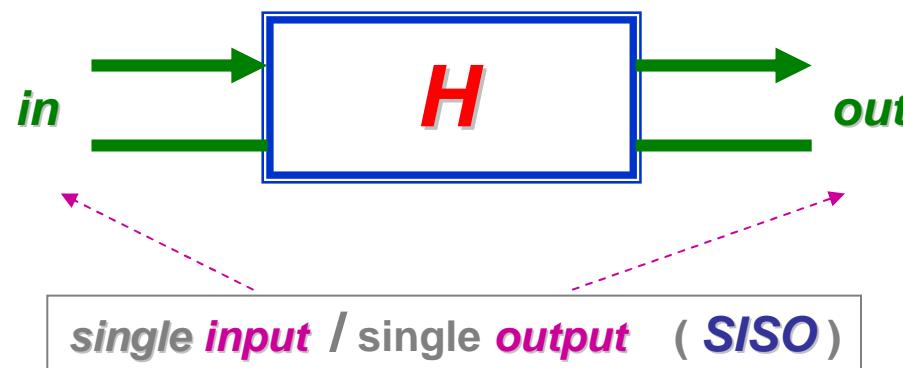
impedance $Z(s)$
admittance $Y(s)$

} immittance

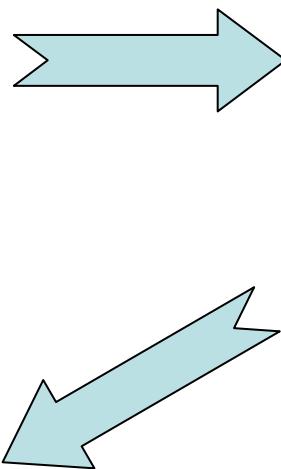
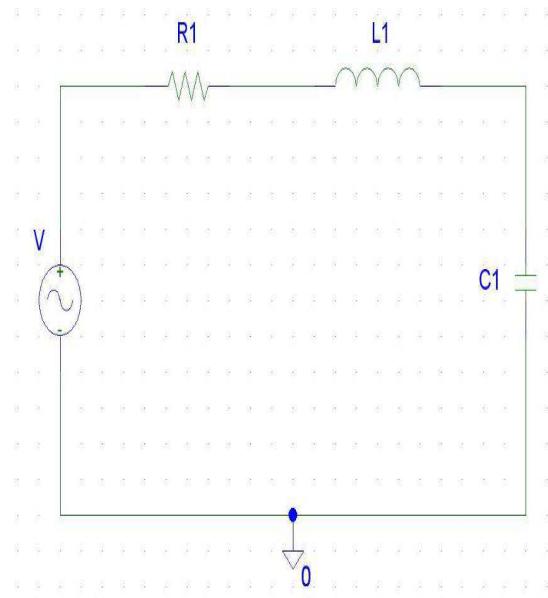
important NOTE



The concept of ***Transfer Function*** assumes that
the ***input*** and ***output*** are **ideal** (= no *interfering impedance*)



UNIT gain ***OPAMP*** may be used to safely cascade ***FILTERS***



$$Z_{L1} = L1 * S$$

$$Z_{C1} = \frac{1}{C1 * S}$$



$$[v_{r1} = \frac{R1 * V}{Z_{L1} + Z_{C1} + R1}]$$

$$[i_{l1} = \frac{V}{Z_{L1} + Z_{C1} + R1}]$$

$$[v_{c1} = \frac{V * Z_{C1}}{Z_{L1} + Z_{C1} + R1}]$$

$$\frac{v_{r1}}{V} = \frac{C1 * R1 * s}{C1 * L1 * s^2 + C1 * R1 * s + 1}$$

$$\frac{i_{l1}}{V} = \frac{C1 * s}{C1 * L1 * s^2 + C1 * R1 * s + 1}$$

$$\frac{v_{c1}}{V} = \frac{1}{C1 * L1 * s^2 + C1 * R1 * s + 1}$$

Those **3** “**transfer functions**” have the general form
(transmittance)

$$\frac{Y(s)}{X(s)} = \frac{s^2 N_2 + s N_1 + N_0}{s^2 D_2 + s D_1 + D_0}$$

An expression of the **Y(t)** “**time response**” can be determined

when

$$\frac{s^2 N_2 + s N_1 + N_0}{s^2 D_2 + s D_1 + D_0} X(s)$$

(or its **decomposition**) can be
retrieved from an
Inverse Laplace Transform Table

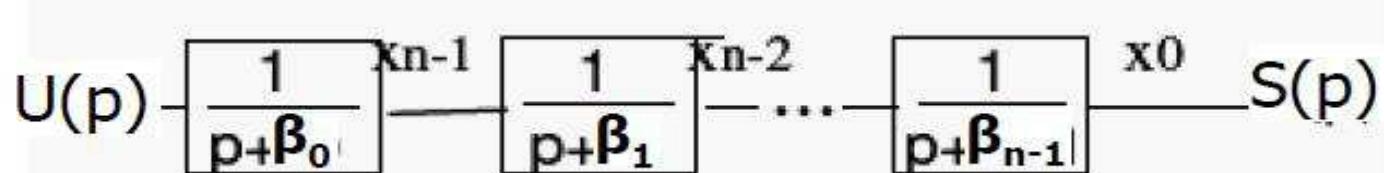


Obviously, this supposes that a “**S**” **transform** of the **INPUT X(t)** is known !!!

DECOMPOSITION

CASCADE

$$H(p) = \frac{1}{(p+\beta_0)(p+\beta_1)\dots(p+\beta_{n-1})}$$



MODAL

$$H(p) = \frac{\alpha_0}{p+\beta_0} + \frac{\alpha_1}{p+\beta_1} + \dots + \frac{\alpha_{n-1}}{p+\beta_{n-1}}$$

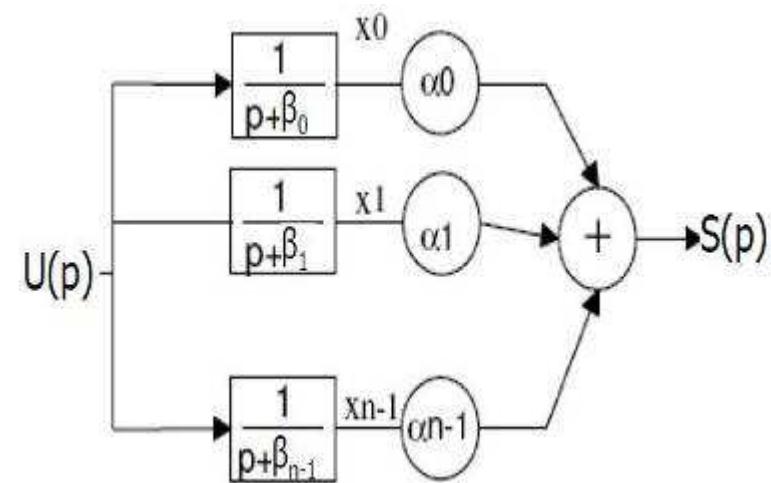
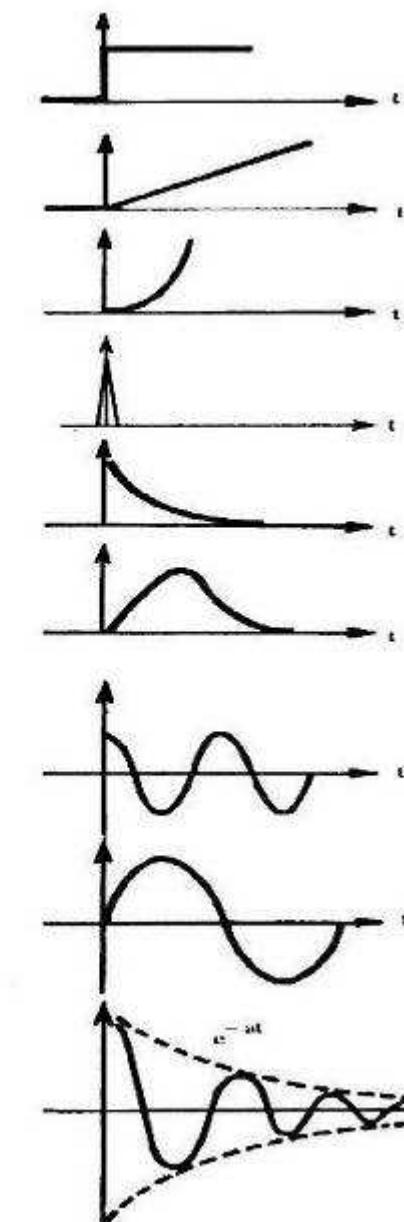
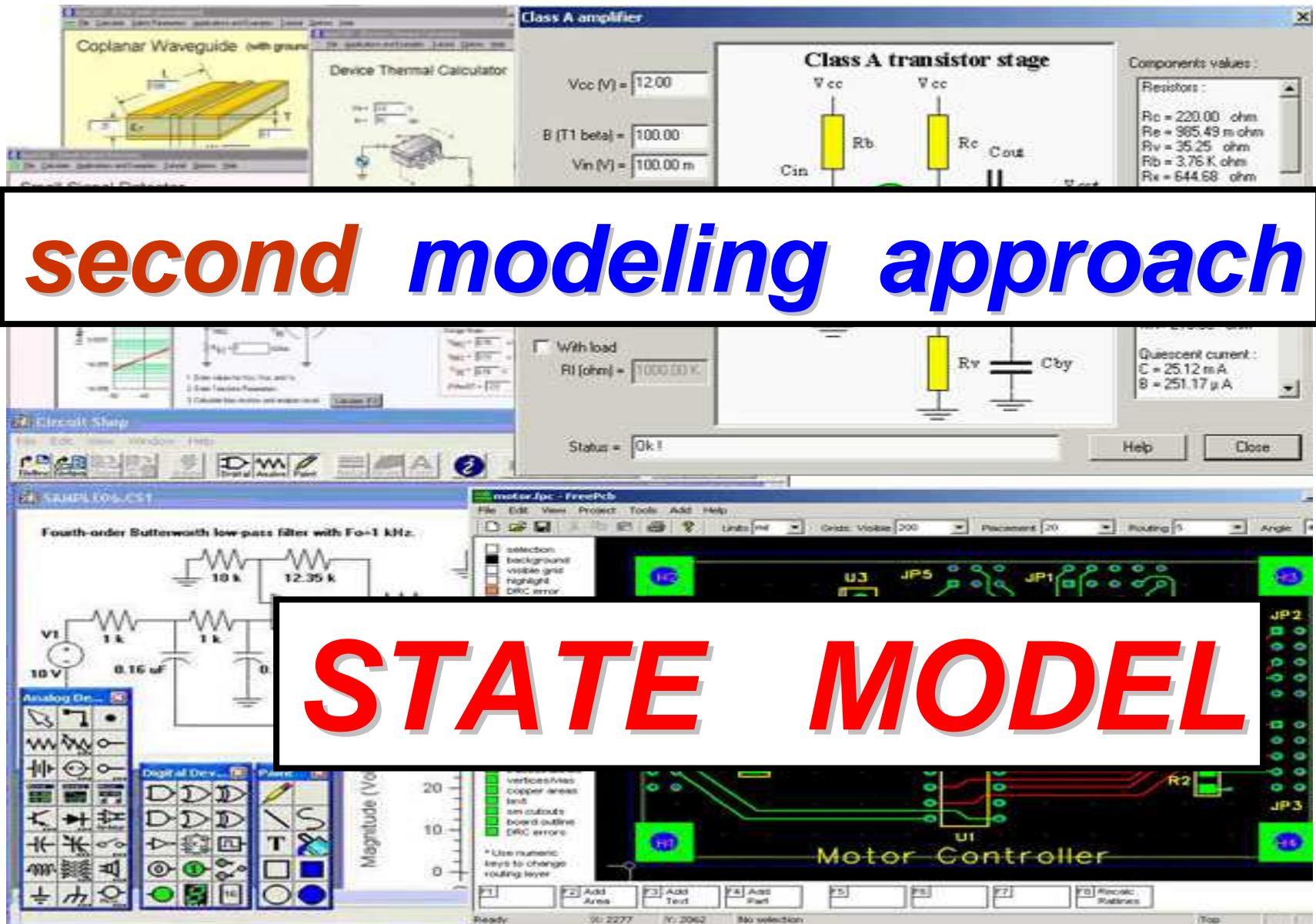


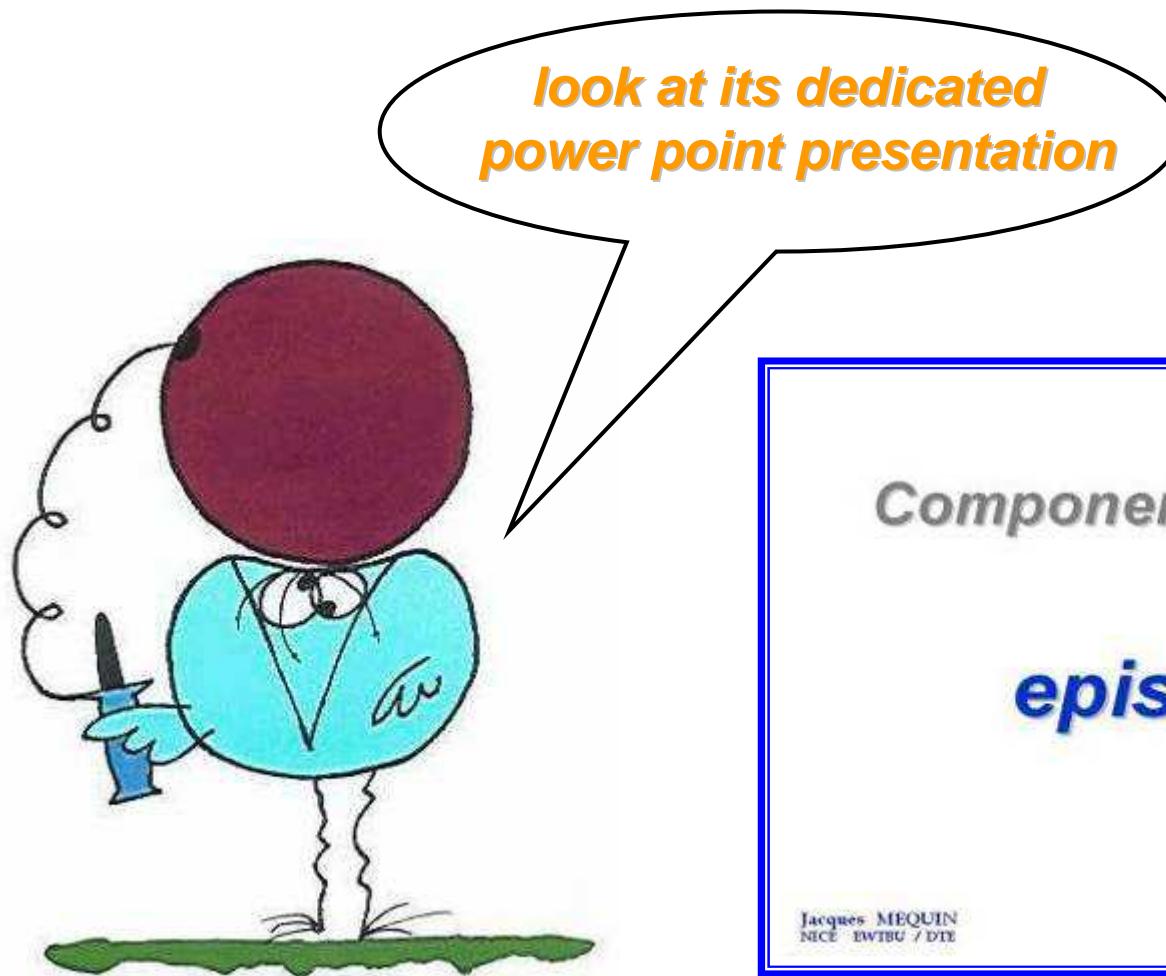
TABLE EXTRACT

| | | | |
|---------------------|---------------------------------------|------------------|--|
| Unit | $\frac{1}{p}$ | HEAVISIDE | 1 |
| Speed | $\frac{1}{p^2}$ | | t |
| Acceleration | $\frac{1}{p^3}$ | | $\frac{t^{n-1}}{(n-1)!}$ |
| Dirac | 1 | DIRAC | at |
| 1 order | $\frac{1}{p+a}$ | | e^{-at} |
| 2 order | $\frac{1}{(p+a)^2}$ | | te^{-at} |
| (n+1) order | $\frac{1}{(p+a)^{n+1}}$ | | $\frac{t^n e^{-at}}{n!}$ |
| Cosinus ... | $\frac{p}{p^2 + \omega^2}$ | | $\cos \omega t$ |
| | $\frac{p}{p^2 - \omega^2}$ | | $\operatorname{ch} \omega t$ |
| Sinus ... | $\frac{\omega}{p^2 + \omega^2}$ | | $\sin \omega t$ |
| | $\frac{\omega}{p^2 - \omega^2}$ | | $\operatorname{sh} \omega t$ |
| | $\frac{p+a}{(p+a)^2 + \omega^2}$ | | $e^{-at} \cos \omega t$ |
| | $\frac{ap+\beta}{(p+a)^2 + \omega^2}$ | | $Ae^{-at} \cos(\omega t + \varphi)$ |
| | | | $A = \frac{1}{\omega} \sqrt{a^2 \omega^2 + (\beta - a\omega)^2}$ |
| | | | $\varphi = -\arctan(\beta - a\omega)$ |





STATE MODEL



Component PERCUSSION

episode II

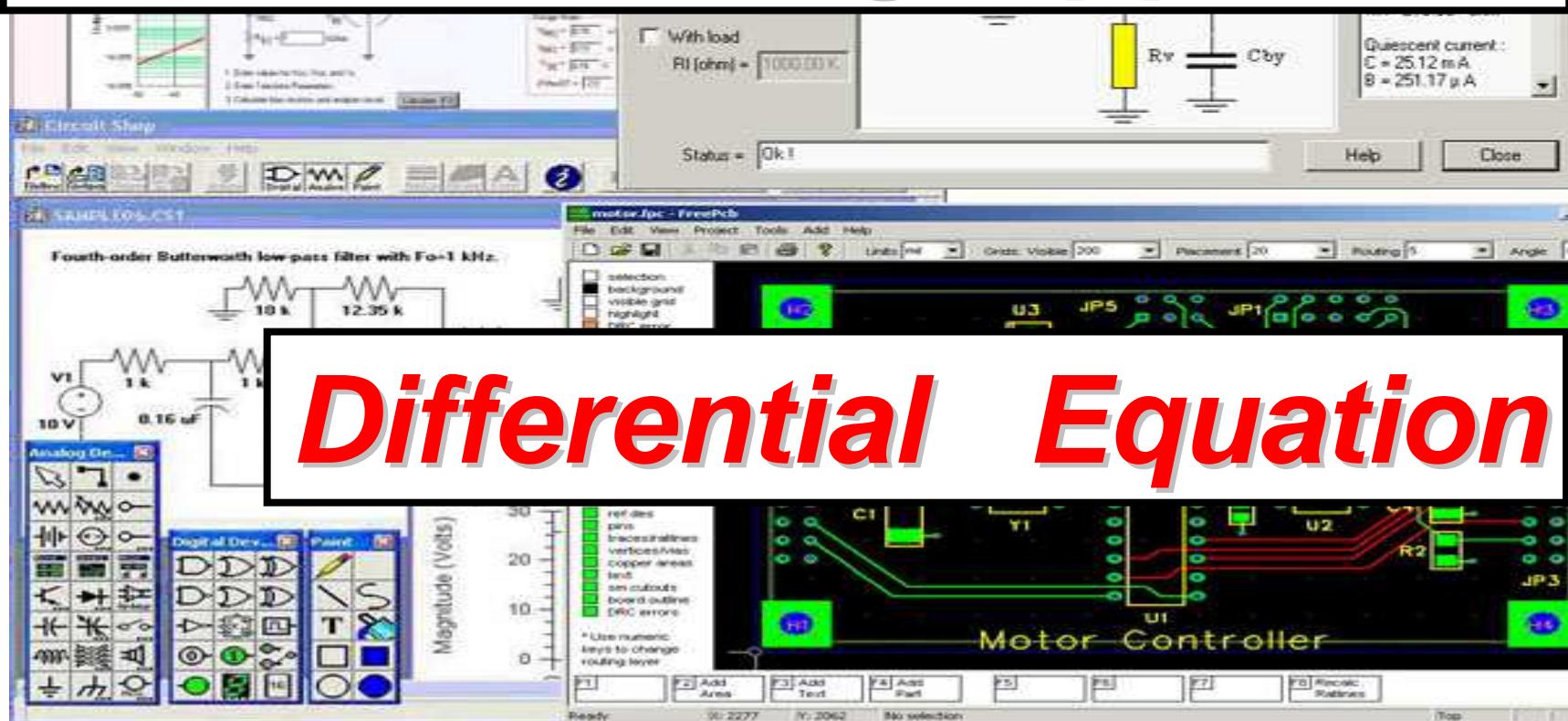
Jacques MEQUIN
NICE EWTBU / DTE

TEXAS INSTRUMENTS
PROPRIETARY INFORMATION
INTERNAL DATA

release



third modeling approach



Differential Equation

A **generic method** to **simulate** the “**time domain**” response of the **transfer function**

$$\frac{Y(s)}{X(s)} = \frac{s^2 N_2 + s N_1 + N_0}{s^2 D_2 + s D_1 + D_0}$$

is to use its equivalent **time domain** “**differential equation**” *

$$\left(\frac{d^2}{dt^2} Y(t) \right) D_2 + \left(\frac{d}{dt} Y(t) \right) D_1 + Y(t) D_0 = \left(\frac{d^2}{dt^2} X(t) \right) N_2 + \left(\frac{d}{dt} X(t) \right) N_1 + X(t) N_0$$

This only requires that **samples** of the **INPUT X(t)** can be evaluated

* **LCCODE** : linear constant coefficient ordinary differential equation

They are numerous methods available to **integrate Ordinary Differential Equation (ODE)**

Among them, the **RK4** method (“*fourth order*” **Runge-Kutta**) is both **accurate** (error $O(h^5)$) and **easy** to implement

Expecting for $y(t+h)$ an expression such as

$$y_{i+1} = y_i + (a_1 k_1 + a_2 k_2 + a_3 k_3 + a_4 k_4) (x_{i+1} - x_i)$$

assuming

$$\frac{dy}{dx} = f(x, y)$$

it is based on the determination of the value of the coefficients that “**match**” the **Taylor series** expansion at **order 4**

$$y_{i+1} = y_i + \frac{dy}{dx}\Big|_{x_i, y_i} (x_{i+1} - x_i) + \frac{1}{2!} \frac{d^2 y}{dx^2}\Big|_{x_i, y_i} (x_{i+1} - x_i)^2 + \frac{1}{3!} \frac{d^3 y}{dx^3}\Big|_{x_i, y_i} (x_{i+1} - x_i)^3 + \frac{1}{4!} \frac{d^4 y}{dx^4}\Big|_{x_i, y_i} (x_{i+1} - x_i)^4$$

due to supernumerary (*free*) variables,
several possible set of coefficients **k1**, ..., **k4**
are acceptable solutions

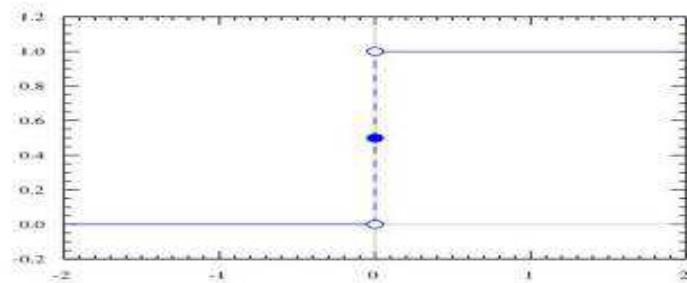
**The following choice is generally
very convenient**

$$\begin{aligned} k_1 &= \Delta t f(t_i, y_i) \\ k_2 &= \Delta t f(t_i + \frac{1}{2}\Delta t, y_i + \frac{1}{2}k_1) \\ k_3 &= \Delta t f(t_i + \frac{1}{2}\Delta t, y_i + \frac{1}{2}k_2) \\ k_4 &= \Delta t f(t_i + \Delta t, y_i + k_3) \\ y_{i+1} &= y_i + \frac{1}{6}(k_1 + k_2 + k_3 + k_4) \end{aligned}$$

DISTRIBUTIONS

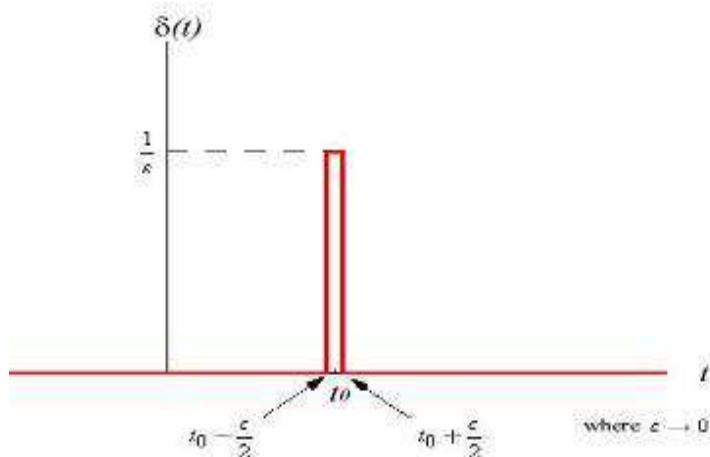
< T , φ >

HEAVISIDE “generalized function”



$$H(x) = \begin{cases} 0 & x < 0 \\ 1 & x > 0 \end{cases}$$

DIRAC “generalized function”



$$\delta(x) \equiv \frac{dH}{dx} = \begin{cases} 0 & x < 0 \\ \infty & x = 0 \\ 0 & x > 0 \end{cases}$$

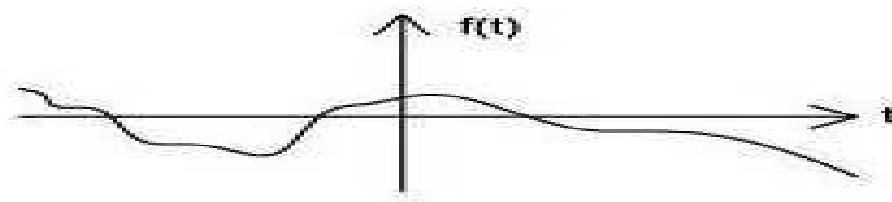
$$x\delta'(x) = -\delta(x)$$

$$\int_{-\infty}^{\infty} \delta(x) f(x) dx = f(0)$$

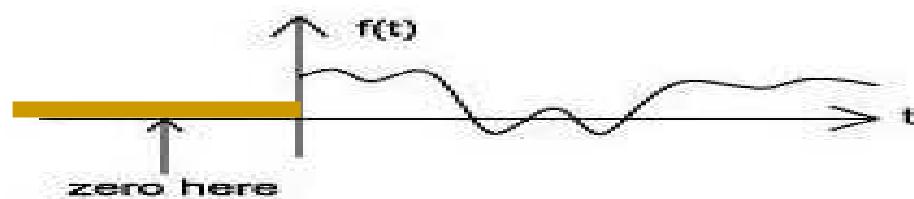
$$\int_{-\infty}^{\infty} a\delta(x) dx = a$$

★ Laurent Schwartz established it as a “mesure” (a particular subspecies of his “distributions” theory)
(Fields medal in 1950)

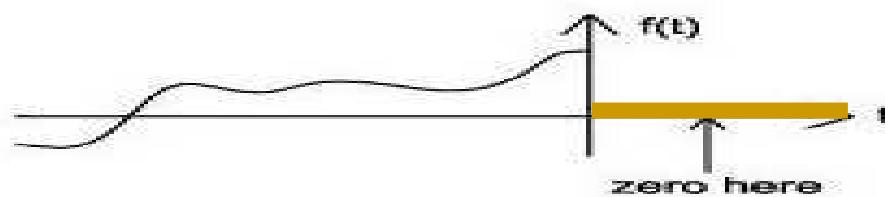
DEFINITIONS



noncausal signal



causal signal



anticausal signal

important relations

$$H'(t) = \delta(t)$$

$$f(t)\delta(t) = f(0)\delta(t)$$



$$H'(t-a) = \delta(t-a)$$

$$f(t)\delta(t-a) = f(a)\delta(t-a)$$

Let us consider a signal $y(t)$, its “CAUSAL representation” is $Y(t)$

$$Y(t) = y(t) H(t)$$

$$\text{Note : } (uv)' = u'v + uv'$$

$$Y'(t) = \underbrace{y'(t)H(t)}_{y(t)H'(t)} + \underbrace{y(t)H'(t)}_{y(t)H(t)} = y'(t)H(t) + y(0)\delta(t)$$

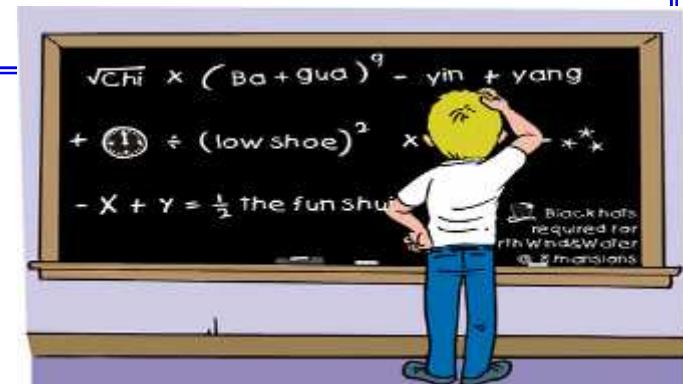
the expression of next derivative will simplify as

$$Y''(t) = y''(t)H(t) + y'(0)\delta(t) + y(0)\delta'(t)$$

etc . . .

The next thing to do would be to show what is happening when we inject those definitions into the differential equations of the circuit

maybe another day . . .



STABILITY

★ Poles and stability of a system (of any order)

➤ *Look at the real part of each pole:*

- If the **real part of ALL poles is negative** then all components of the system are **stable** and the **overall system is stable**
- If **ONE pole (or more)** has a **zero real part** then that component is critically/marginally stable and the **overall system is critically stable**
- If **ONE pole (or more)** has a **positive real part** then that component lead the **overall system to instability**

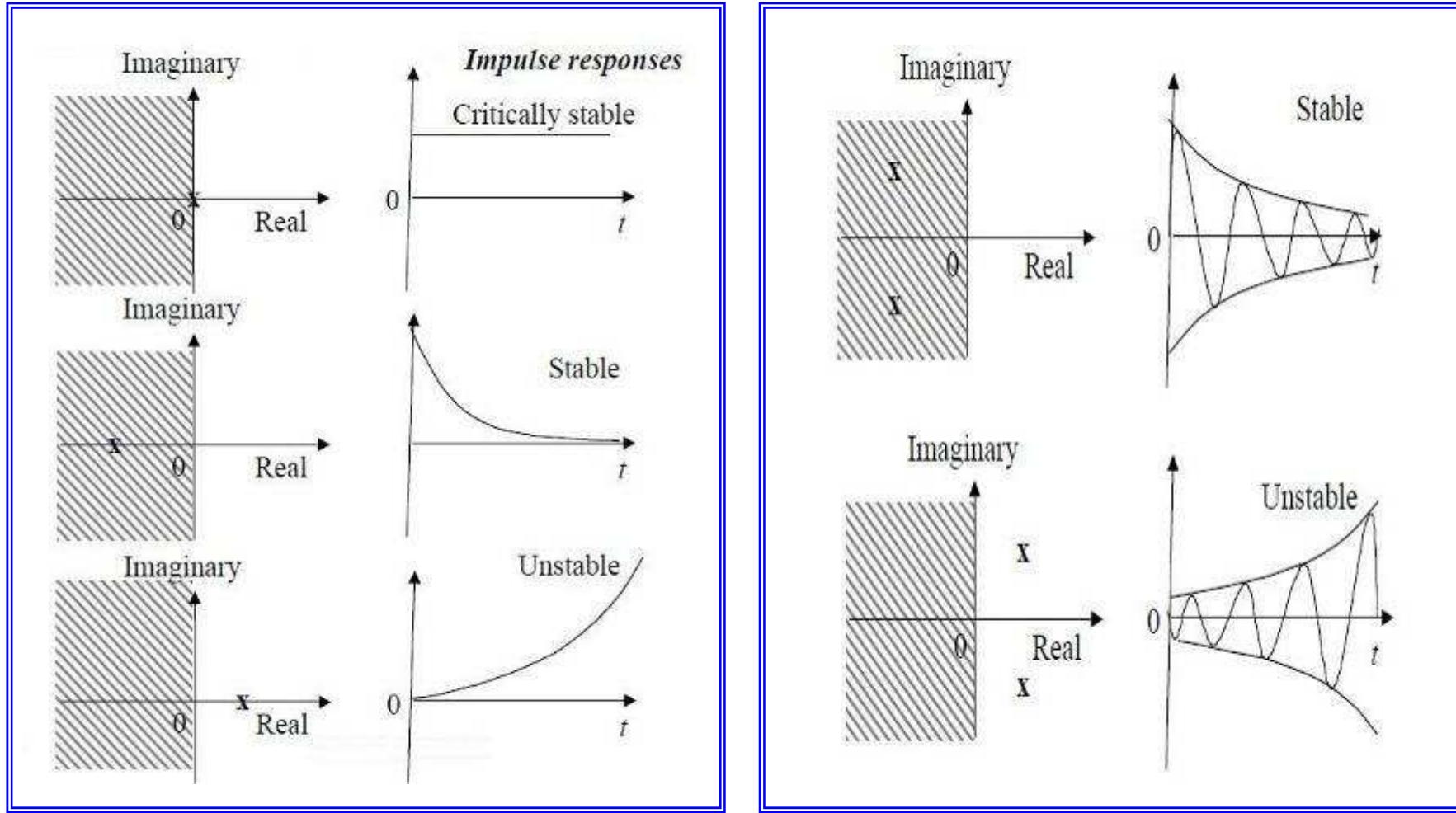


➤ *Look at the imaginary part for the presence of oscillations:*

- If the imaginary part of a pole is zero then that component does not have any oscillatory contribution
- If the imaginary part is not zero then its value is the frequency of oscillation of the corresponding component of the system

★ The zeros of a system do not affect the stability

- The zeros affect the transient response of the system



For **higher order** systems that have many poles,
the poles that are **further to the right** of the **s-plan** will **dominate** the response

Numerically, there is another issue

pole order



$$\begin{aligned} f &= 100 \text{ MHz} \implies \omega &= 2\pi f &= 6.28 \cdot 10^8 \\ \omega^2 & & &= 3.94 \cdot 10^{17} \\ \omega^3 & & &= 2.48 \cdot 10^{26} \end{aligned}$$

...

Such very large \pm exponent are also found in the $Y(t)$ high order derivatives

($f \uparrow$ \rightarrow step $h \downarrow$ very small and powered as h^2, h^3, h^4, \dots)

quantum slice of the period
must be SHANNON compliant to avoid aliasing

accordingly to TAYLOR expansion

NOTE: To minimize this issue the **ILT package** is coded “internally” with floating “long double” (x87 standard : 80 bits / 19 digits / $\approx \pm 4932$ exponent)

None of **3 methods** presented **structurally handle percussions**
since, *by hypothesis*, the circuit is assumed **Linear Time Invariant**

Nevertheless,

simulating the original circuit “up to” t^-

then, “starting at” t^+ , simulating a **modified circuit**
having as “**initial conditions**” the **final values** reached
just **before** the **percussion**

does play the trick !!!

This is OK !!!, but the making of such **simulator** changing the **user defined circuit**
during the runtime is likely to be **clumsy**

The aim of this presentation is precisely to claim that

an **equational solution exist** to stick, **all along the simulation,**
with a more seamless approach

NAPA



15 years ago,

Yves Leduc who is the **owner, developer** and **enhancer** of an "**in-house**" **analog functional simulator** (named **NAPA**) has asked me to develop a **plug-in** performing the **Inverse Laplace Transform**

Along the years, collecting new algorithms in the **literature** and on the **web** **NAPA** has succeed to become a unique collection of **state of the art mathematical library** feature (sophisticated **FFT** windows, **noise** analysis, etc ...)

NAPA “ILT” plug-in hypothesis

For the **ILT package**, the output **Y[t]** of a **transfer function** is the same

if INPUT[t] is changing while INPUT[t - h] did also change value

or if INPUT[t - h] was maintaining a steady value

such **heuristic** (that neglects step “**input variation**”) enables the **ILT package** to behave independently of the **INPUT derivatives**
(in line with the **STATE MODEL** activated by a **steady input** during a **step interval**)

Therefore, such **ILT processing can be viewed as “Markovian”** ★

This is particularly well adapted to **non-vectorial** simulator such as **NAPA** that can **consequently** consider **any filter** as an **ordinary function** to which it has trivially to provide just a simple “**scalar**” **INPUT** “**sample**” at each call

★ the state of such process at time **t** is independent of the history of the process before its previous state

$$\frac{Y(s)}{X(s)} = \frac{s^2 N2 + s N1 + N0}{s^2 D2 + s D1 + D0}$$



**N2
N1
N0**

$$\left(\frac{d^2}{dt^2} Y(t) \right) D2 + \left(\frac{d}{dt} Y(t) \right) D1 + Y(t) D0 = \left(\frac{d^2}{dt^2} X(t) \right) N2 + \left(\frac{d}{dt} X(t) \right) N1 + X(t) N0$$

$\downarrow = 0 \quad \downarrow = 0$

X(t) causal $X(t) = x(t) H(t)$

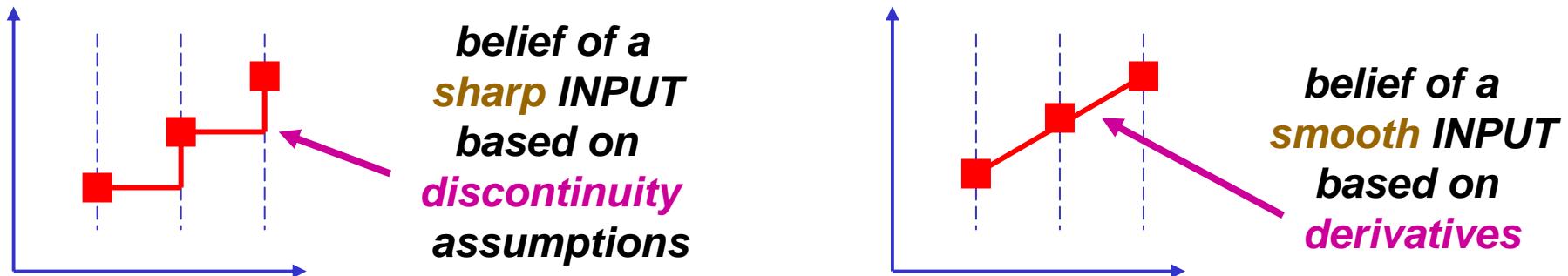
developing this *differential equation* even with input $x''(t) = 0$ and $x'(t) = 0$
 still leave the opportunity for the **numerator coefficients N2 and N1**
 to contribute to the **Y(t) output** response equation

The “**Markovian**” approach of the **ILT package** is motivated by another **tactical choice**

- ➡ To assume that the **INPUT** potentially may contain some **sharp discontinuities** (**Initialization** being one of them ➡ **ILT curves** at “ t_{0+} ” **match the theory**)
- ➡ and consequently, **not to assume**, estimate (nor “*interpolate*”), any **smoothness** of the **INPUT**

This would **cost** to store (to an *arbitrary chosen depth*) a vector of the previous **INPUTs** and to *interpolate* them to compute **associated derivatives**

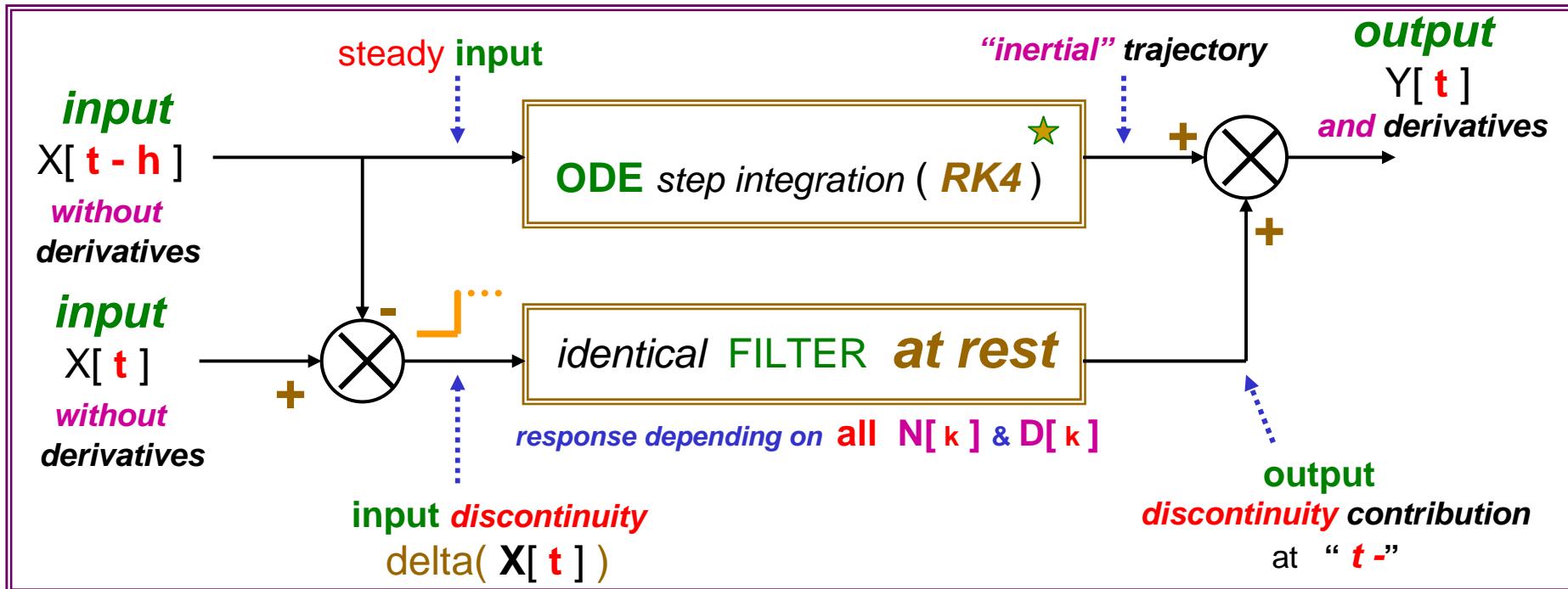
This would also add some extra pain in case of **STEP** changes (*re-sampling*)



In other words, the **expectation** to accurately calculate the response to some **sharp** signal is somewhat tight to the **necessity** of a **brut force** processing of the **INPUT**

Quality of the **LINEAR** response ???

So, the **output** of the **ILT package** is calculated as



In summary, the hypothesis of sharp INPUT based on discontinuity assumptions enables accurate responses of “switched” type of circuits (in particular, the proper calculation of filter “impulse response”)

However, for “smooth” input, the error is closely tight to the user STEP granularity

★ the **ILT package** uses **RK4** for historical reasons, a simpler integration should be OK

ILT plug-in package functions :

ILT_init_LAPLACE() to define the *transfer function* **NUMERATOR**
DENOMINATOR
and the “***time step***”

ILT_change_h() enable ‘*on the fly*’ change of the “***time step***”

ILT_step_LAPLACE() **input(k + 1) == >> output(k + 1)**

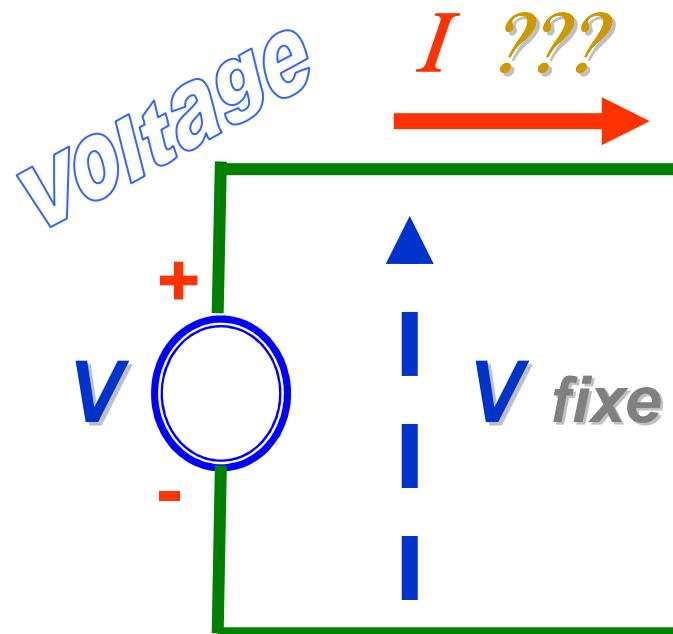
After all those fascinating (?) slides of introduction,

Now, we can go back to our frightening question :

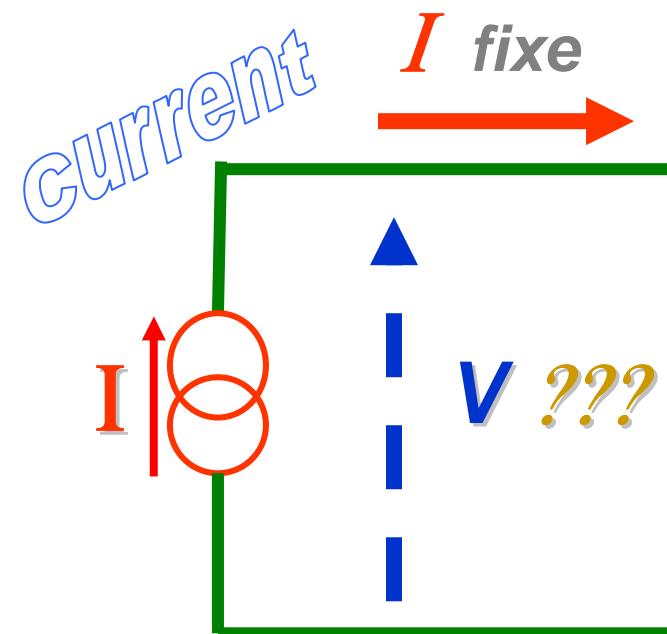
Component PERCUSSION of Linear Time Invariant analog circuit

What is happening ???

“IDEAL” SOURCES

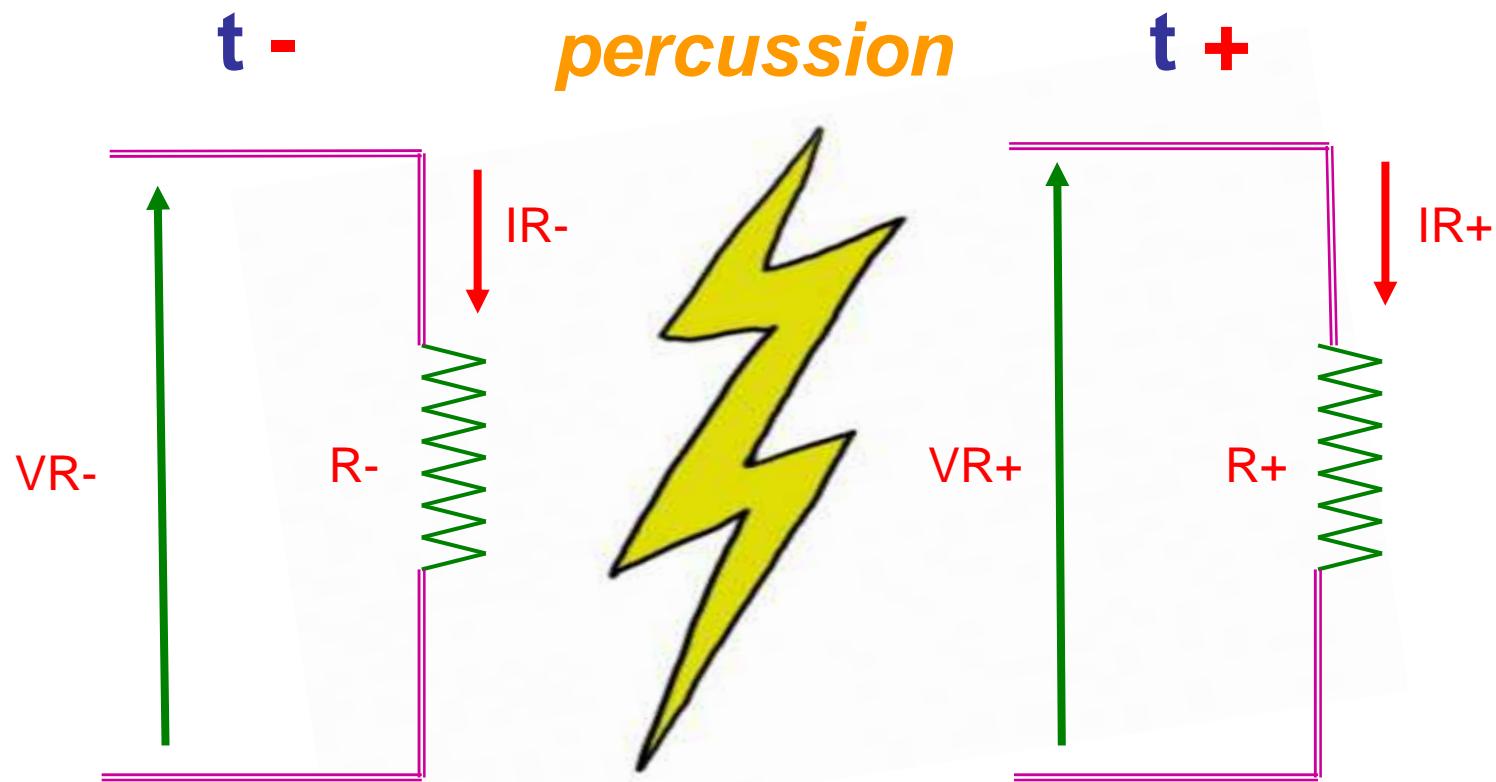


*the current
will take whatever value
as a consequence of the
connected circuit*



*the voltage
will take whatever value
as a consequence of the
connected circuit*

In both cases, the VOLTAGE and CURRENT are not “interdependent”



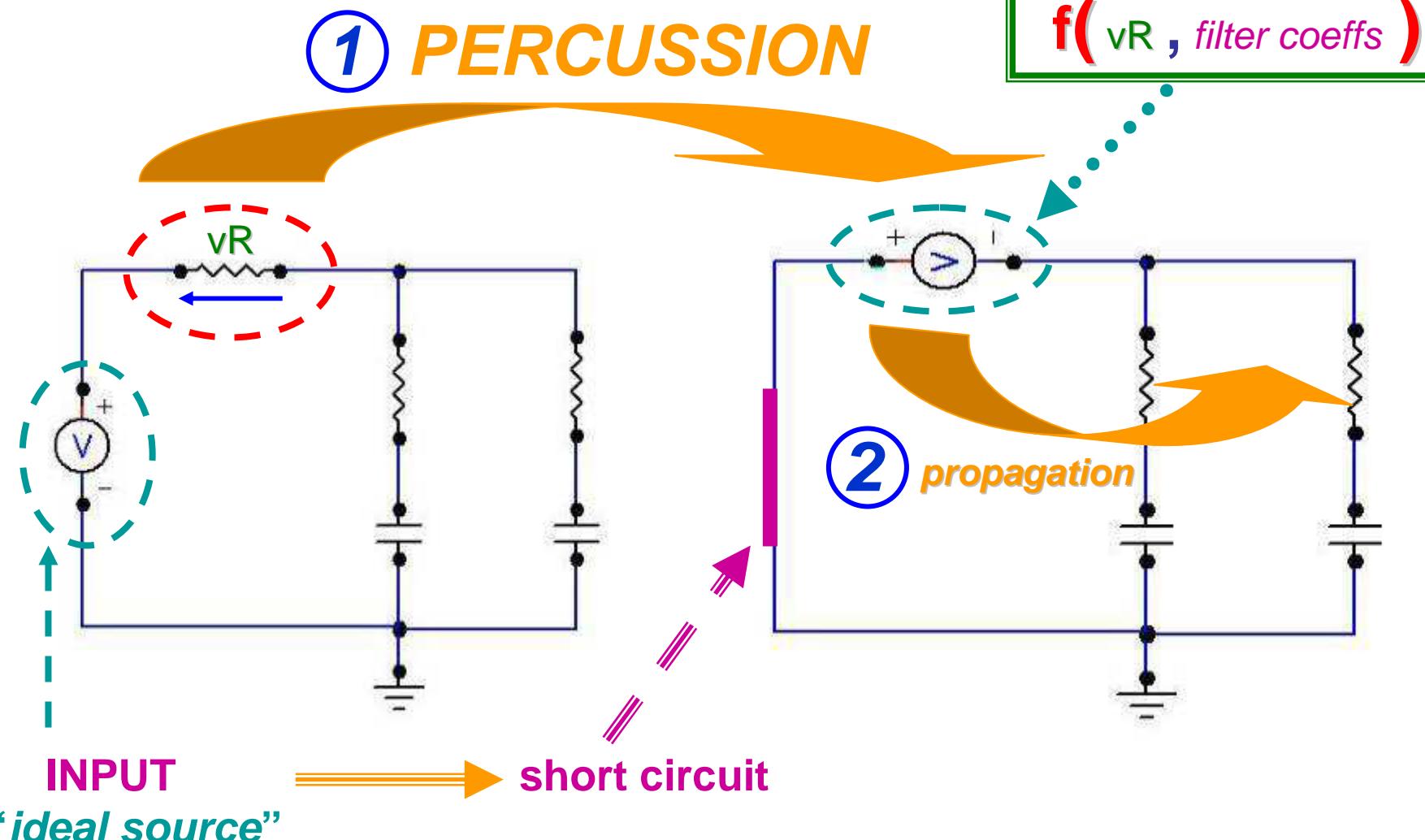
*At the precise instant of the **percussion**,*

VR δ and **IR δ** are paradoxically no longer “**interdependent**”

since the “**percuted**” component can be viewed (during such **null time**) as an “**ideal source**”

2 stages mystery . . .

"ideal source"
but for a null time !!!



ILT PERCUSSION processing

FIRST stage :

"COMPONENT" PERCUSSION

ILT plug-in package function: *ILT_percussion()*

Please remember that the **ILT** package **initialization** has only defined the **coefficients Numerator / Denominator** of some **transfer functions**

for example, **the numerical values
of the polynomial coefficients**
 $N = [n_1, 0]$ and $D = [d_2, d_1, 1]$

$$\frac{v_{R1}}{V} = \frac{C1 R1 s}{C1 L1 s^2 + C1 R1 s + 1}$$

Can we just make a call to a function with an **updated** set of polynomial coefficients (computed, for example, using a new value of “**R1**”) **without passing**★ the information that a **R**, a **L**, or a **C** has changed ???

★ This is a very ambitious **assumption** regarding the assumed impact of the coefficients on the behavior of the **targeted differential equation variables**, demonstration is beyond this presentation

Yes !!!, *equations* can be found using “**DISTRIBUTION**” formalism

(as expected, they predominantly depend on the **old & new denominator coefficients**)

but . . .

[it is OK to “**percute the Transfer Function**” representative of the **voltage**, or the **current**, across a **RESISTOR**]



for a **CAPACITANCE** value change, it is mandatory to “**percute**” the **voltage** since the **current** and the charge **Q** will remain unchanged by the **percussion**
It is the **voltage** that will **evolve**, according to the new **C**, to maintain **$Q = CV$**

for an **INDUCTANCE** value change, it is mandatory to “**percute**” the **current** since the **voltage** and the **flux ϕ** will remain unchanged by the **percussion**
It is the **current** that will **evolve**, according to the new **L**, to maintain **$\phi = L I$**



[If several components have **simultaneously** to change their values they must do it through **several consecutive (2 stages) percussions**]

VALID SIMULATIONS



| | <i>never percuted</i> | <i>sometimes percuted</i> | | |
|---|----------------------------------|----------------------------------|--|--|
| R | $V_R(s) = H(s) \text{ input}(s)$ | $I_R(s) = H(s) \text{ input}(s)$ | $V_R(s) = H(s) \text{ input}(s)$ | $I_R(s) = H(s) \text{ input}(s)$ |
| L | $V_L(s) = H(s) \text{ input}(s)$ | $I_L(s) = H(s) \text{ input}(s)$ | $V_L(s) = H(s) \text{ input}(s)$ | $I_L(s) = H(s) \text{ input}(s)$ |
| C | $V_C(s) = H(s) \text{ input}(s)$ | $I_C(s) = H(s) \text{ input}(s)$ | $V_C(s) = H(s) \text{ input}(s)$ | $I_C(s) = H(s) \text{ input}(s)$ |

ILT PERCUSSION processing

SECOND stage :

PERCUSSION “PROPAGATION”

ILT plug-in package function: ILT_export_percussion()

Now that an **updated value** of the **voltage** (or **current**) across the “**percuted component** has been computed, **we must update** the **voltage** (or **current**) across **any other output** for which the user has defined, **at initialization**, an **observable transfer function**

That is, for our **RLC** example, ***iL1*** & ***vC1***

$$\frac{i_{L1}}{V} = \frac{C_1 s}{C_1 L_1 s^2 + C_1 R_1 s + 1}$$
$$\frac{v_{C1}}{V} = \frac{1}{C_1 L_1 s^2 + C_1 R_1 s + 1}$$

HOW TO DO IT ???



The “percuted” component **will act** (but for a **null time**)
as an **ideal INPUT source** within a **virtual circuit** having
a **short circuit** in place of the **regular INPUT !!!**



So, we have to predefine, for those **OUTPUTs** requiring a **side effect** of the **main percussion**, a set of **null execution time** (!!!) **transfer functions** modeling those **OUTPUTs** in the **virtual circuit**



equations can be found using **“DISTRIBUTION” formalism**
to **superpose** to their current **OUTPUT** value at “ t^- ” **the effect** of the **main discontinuity** propagated through the **virtual circuit**

simulation model:

ALL percuted component ILTs

+ user observed signal ILTs

INIT

```
ILT_init_LAPLACE( pILT_vr0_x , h , N_vr0_x , D_vr0_x , 2 , 2 , 0 ) ;  
ILT_init_LAPLACE( pILT_vc1_x , h , N_vc1_x , D_vc1_x , 1 , 2 , 0 ) ;  
ILT_init_LAPLACE( pILT_vc2_x , h , N_vc2_x , D_vc2_x , 1 , 2 , 0 ) ;
```

```
vr0 = ILT_step_LAPLACE( pILT_vr0_x , x ) ;  
vc1 = ILT_step_LAPLACE( pILT_vc1_x , x ) ;  
vc2 = ILT_step_LAPLACE( pILT_vc2_x , x ) ;  
  
if ( percussion_R0 ) {  
    UPDATE_ALL_COEFFS( new_R0 ) ;
```

*step from
previous t to t-*

*virtual circuit
vc1 = H(vr0)*

*stage #1 R0 percussion
stage #2 propagation*

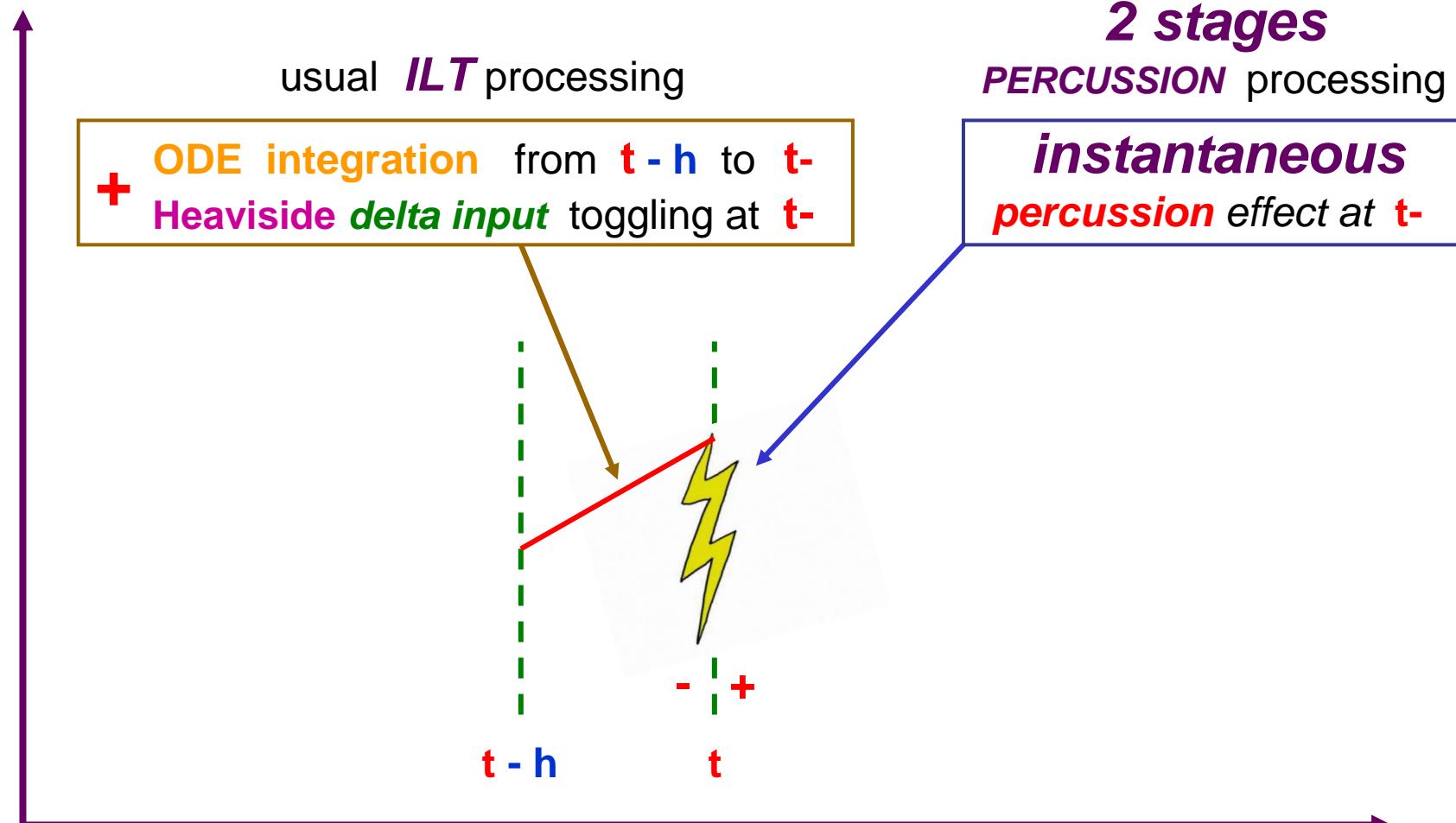
```
vr0 = ILT_percussion( pILT_vr0_x ) ;  
vc1 = ILT_export_percussion( pILT_vr0_x , pILT_vc1_x , pILT_vc1_$vr0 ) ;  
vc2 = ILT_export_percussion( pILT_vr0_x , pILT_vc2_x , pILT_vc2_$vr0 ) ;
```

```
}  
  
if ( percussion_C1 ) {  
    UPDATE_ALL_COEFFS( new_C1 ) ;
```

*stage #1 C1 percussion
stage #2 propagation*

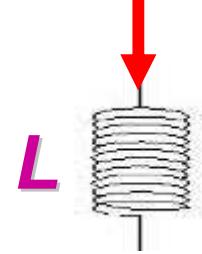
```
vc1 = ILT_percussion( pILT_vc1_x ) ;  
vr0 = ILT_export_percussion( pILT_vc1_x , pILT_vr0_x , pILT_vr0_$vc1 ) ;  
vc2 = ILT_export_percussion( pILT_vc1_x , pILT_vc2_x , pILT_vc2_$vc1 ) ;
```

ILT PERCUSSION processing summary

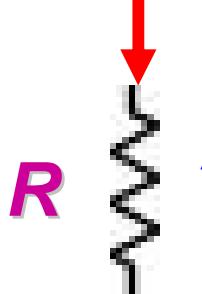


COMPONENT

PERCUSSION



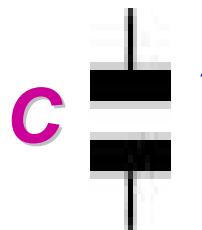
$$I_L = H_I(s) \text{ input}(s)$$



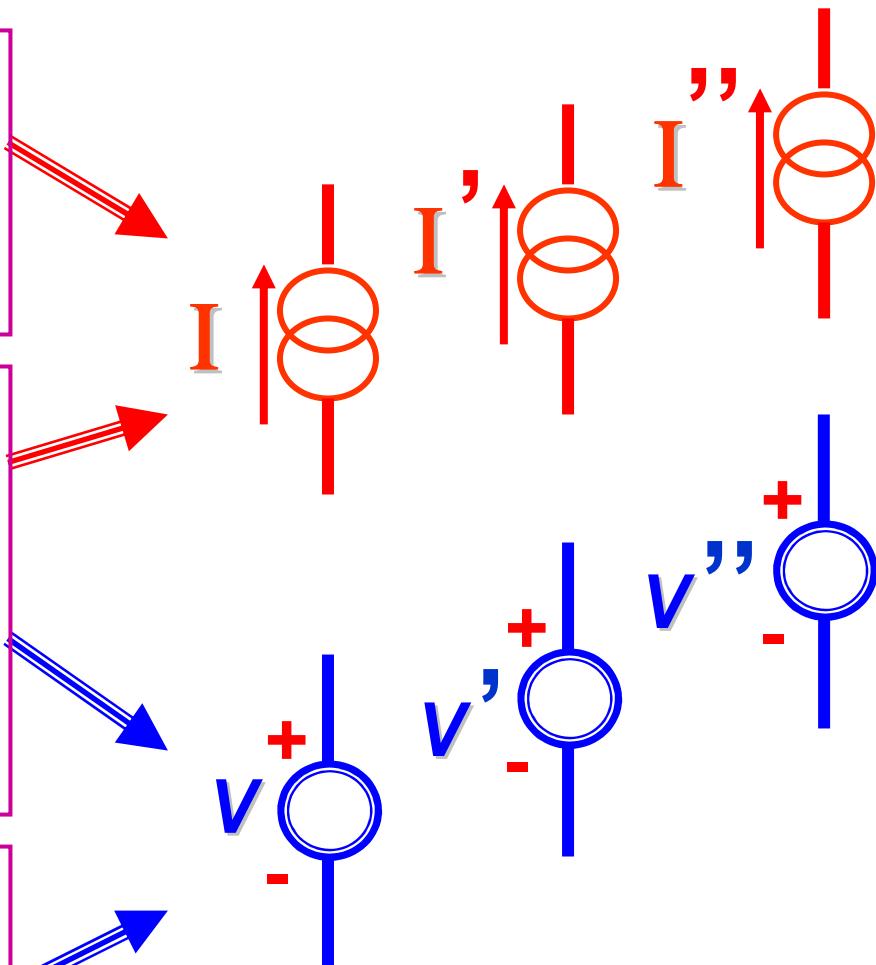
$$I_R = H_I(s) \text{ input}(s)$$

$$V_R = H_V(s) \text{ input}(s)$$

WARNING: V_R and I_R are no longer trivially related



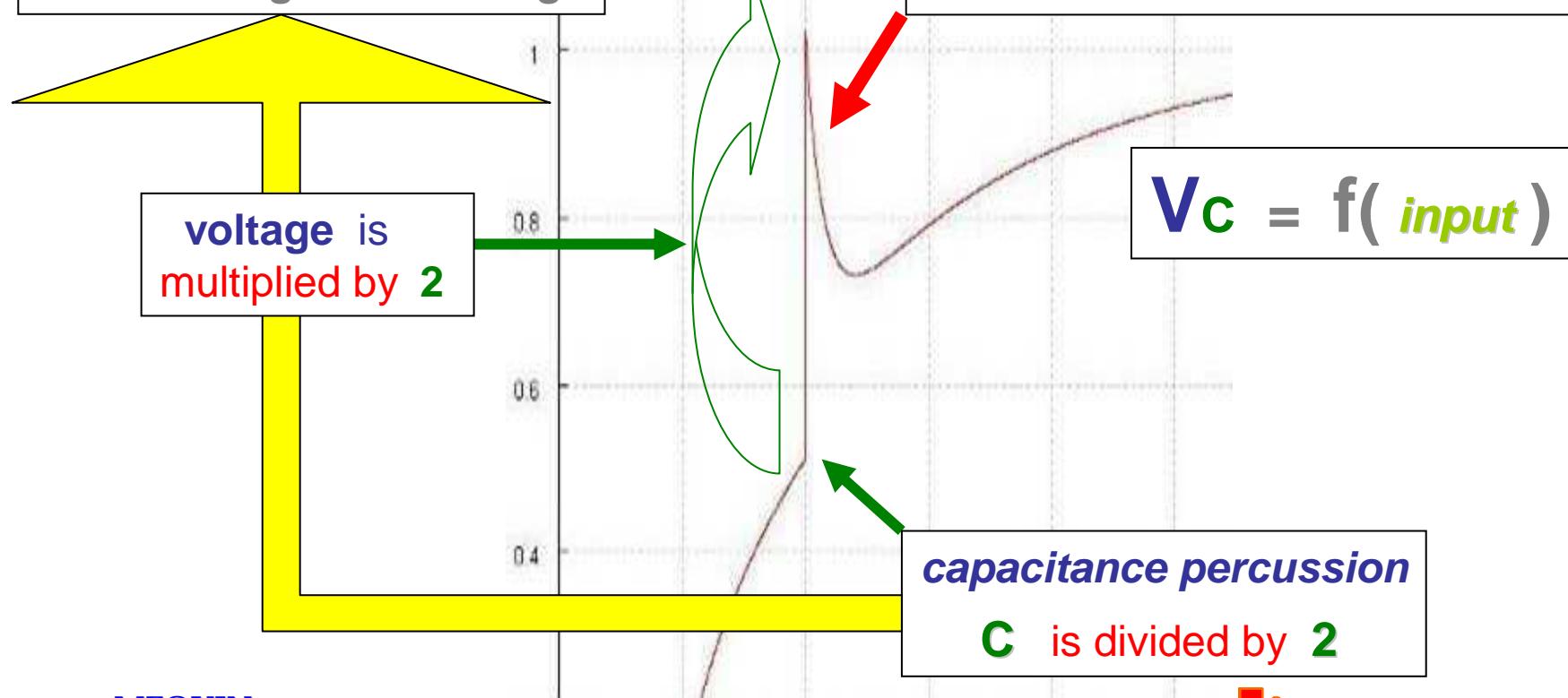
$$V_C = H_V(s) \text{ input}(s)$$



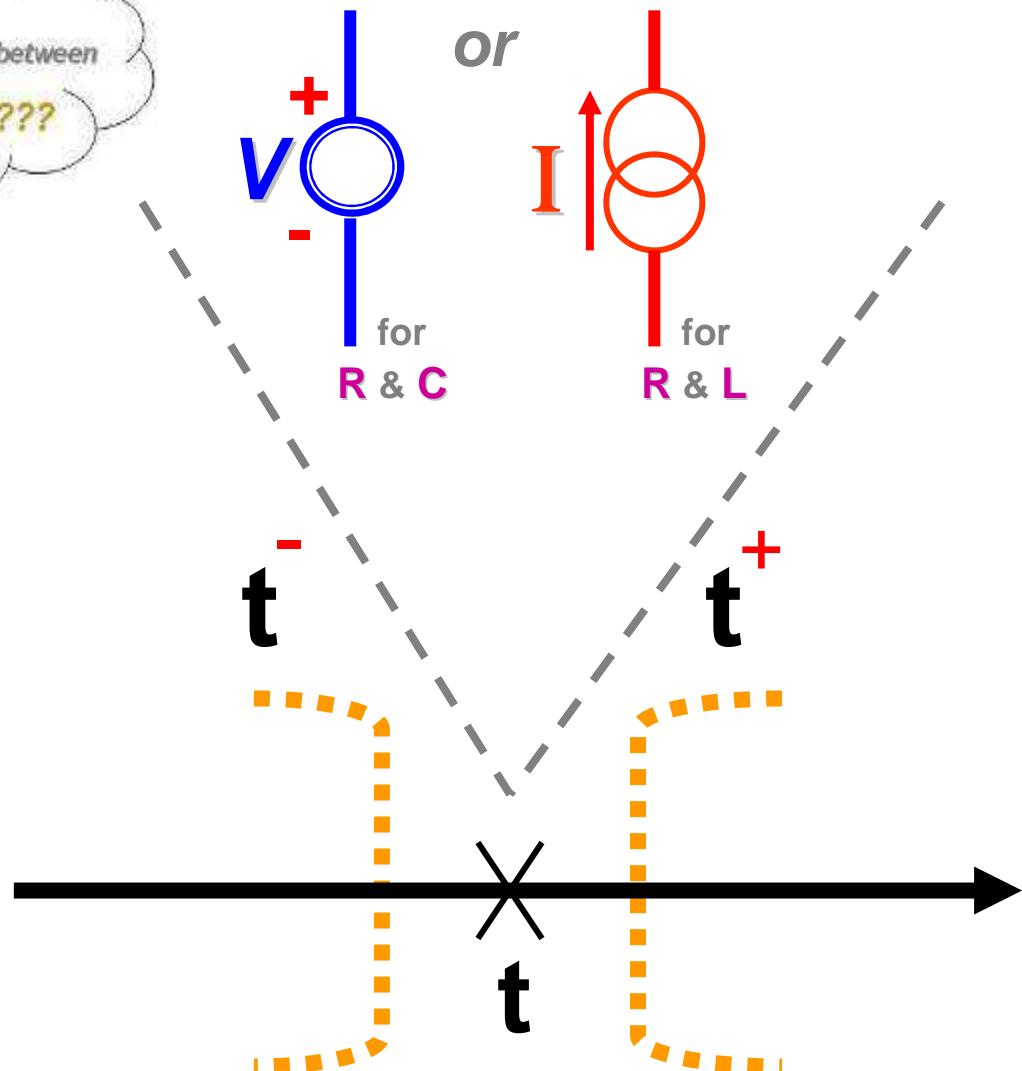
*derivatives of those IDEAL SOURCES
do also have their own equations*

from now on the curve becomes **the resultant** of a new **free-mode + forced-mode**

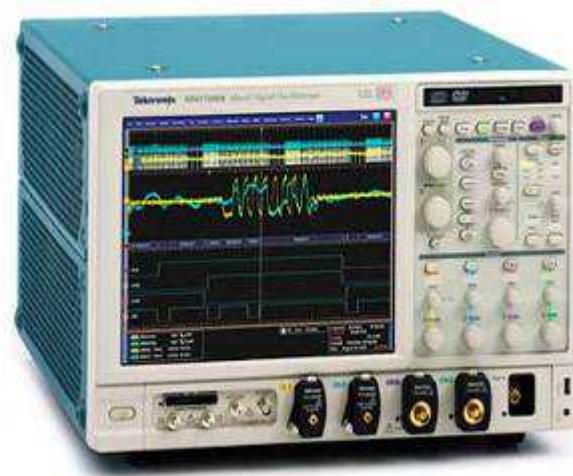
*at the precise instant
of the percussion
there is nothing very
impressive in the fact that
the equations predict
the doubling of the voltage*



percuted component

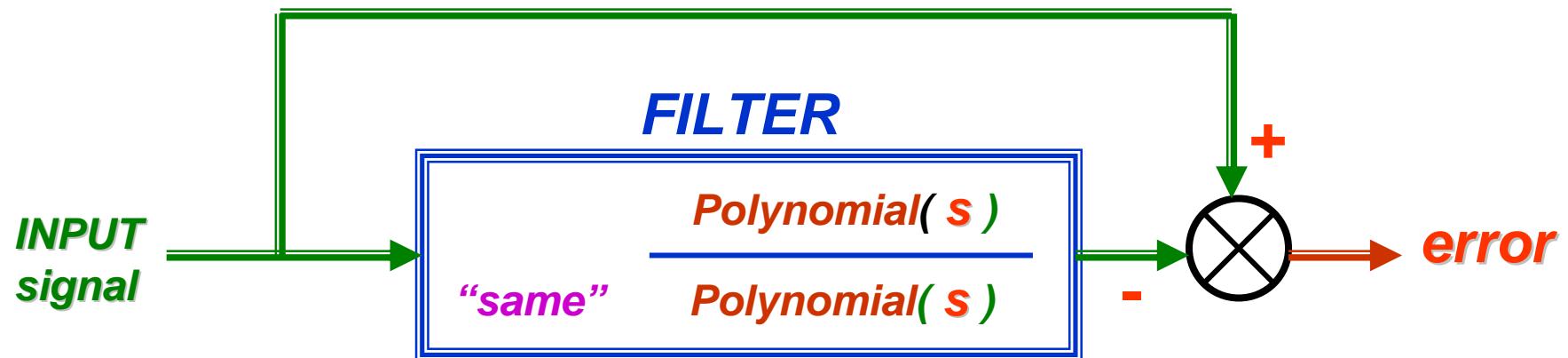
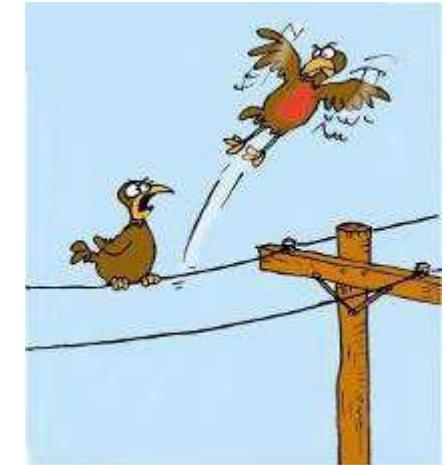


TEST



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NICE EWTBU / DTE
2012 j-mequin@ti.com

“ WIRE ” test



Assuming that the **FILTER** is **stable**, this is expected to behave as an **ideal “wire”**

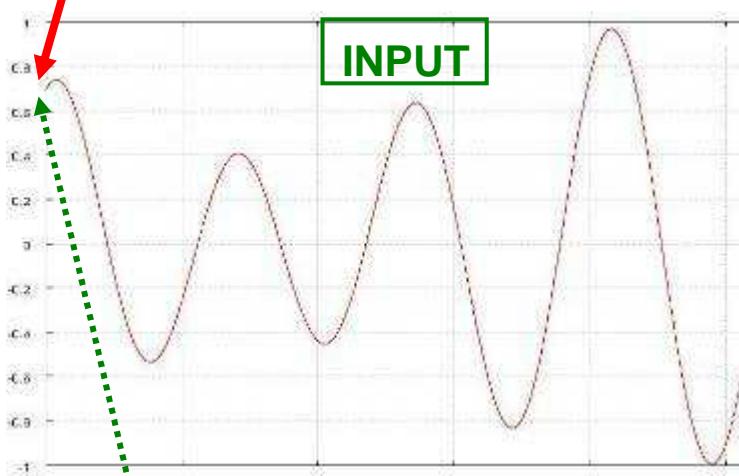
“ WIRE ” test

this filter is “**stable**” by construction

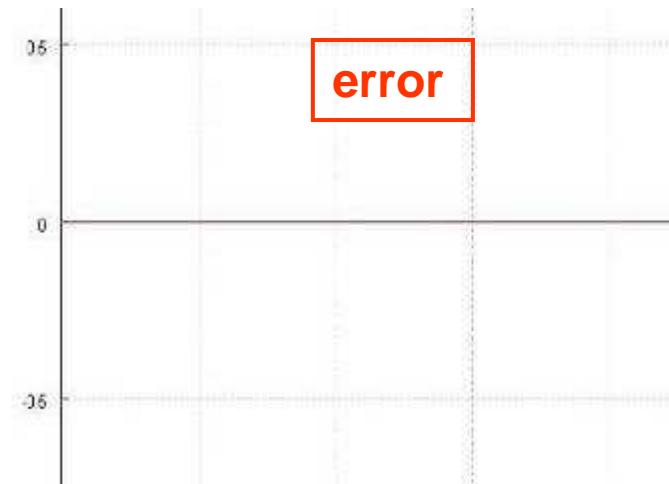
$$ka * \cos(wa * t) + kb * \sin(wb * t)$$

output(0) will only depends on
this first input sample
and the *filter coefficients*

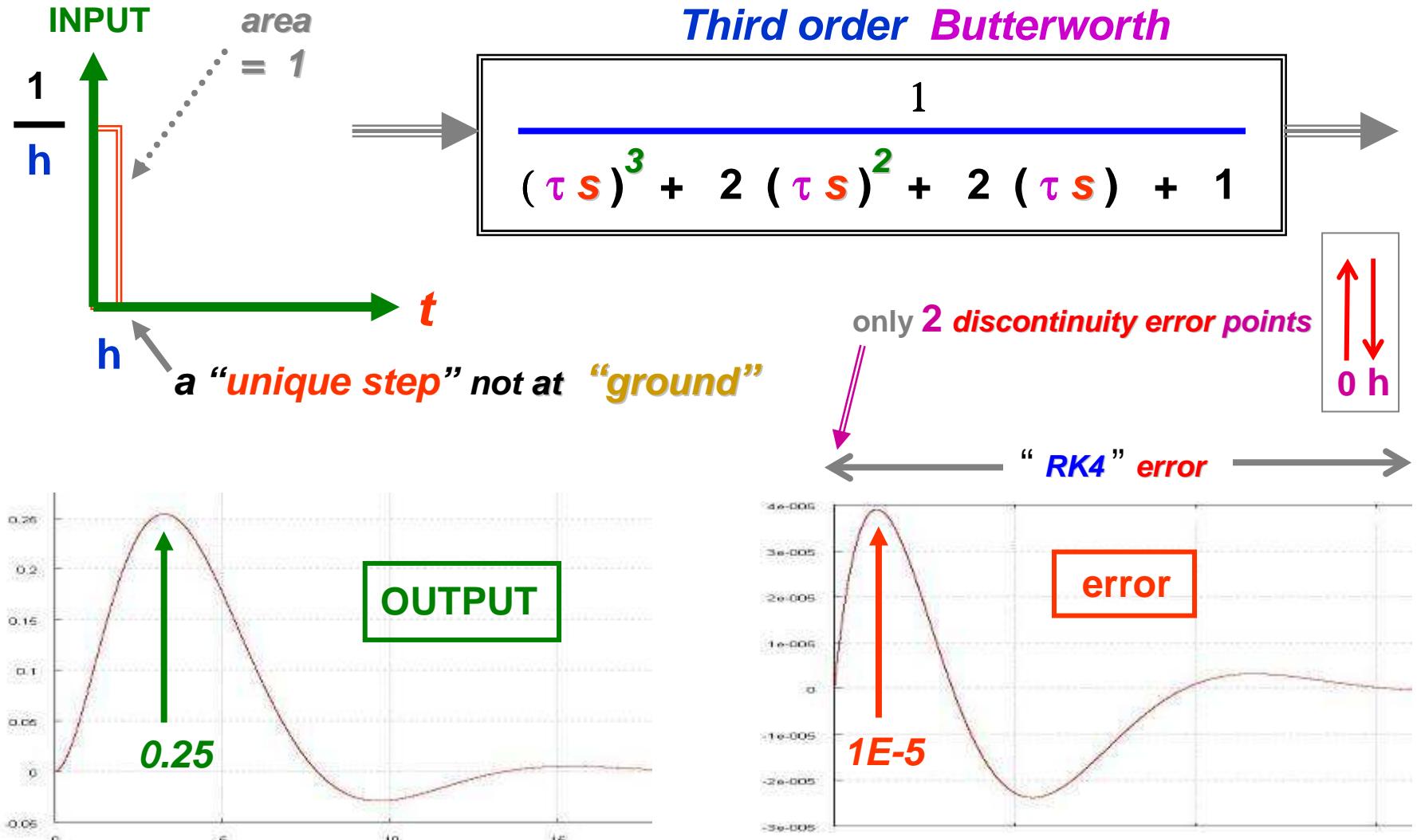
$$\frac{(s+1)(s+2)(s+3)(s+4)(s+5)}{(s+1)(s+2)(s+3)(s+4)(s+5)}$$



to increase **stress** this **activation** does not start at 0

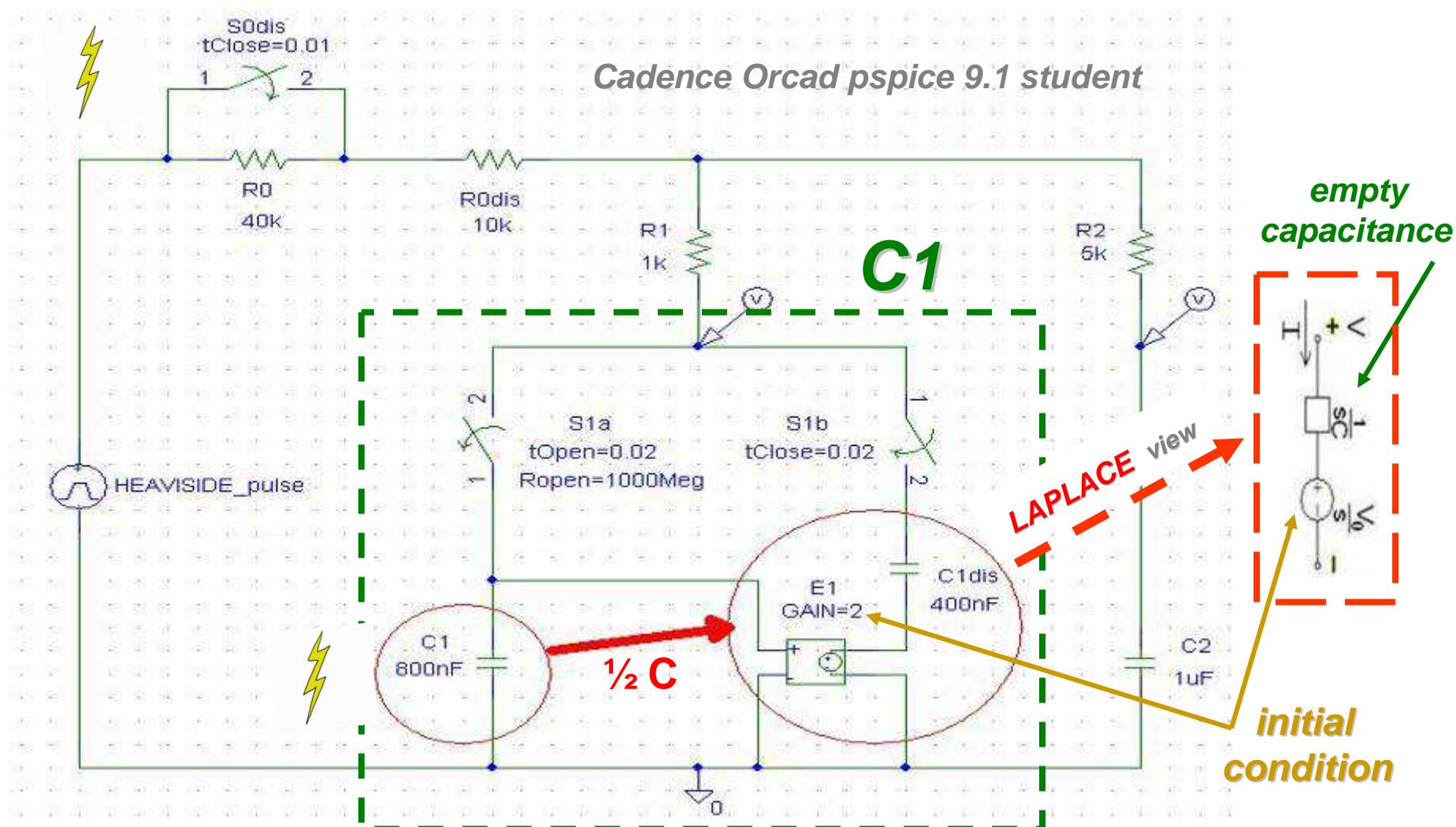


“IMPULSIONAL” response

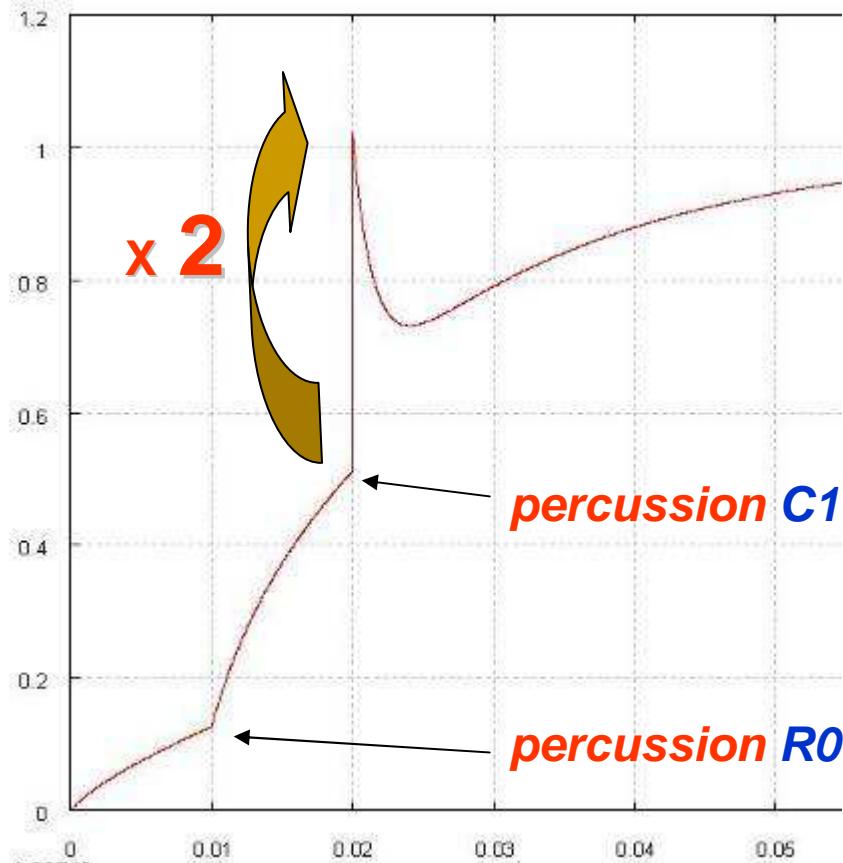


PERCUSSION

at $t = 0.01$ $R_0 = 50 \text{ Kohm} \Rightarrow 10 \text{ Kohm}$
 at $t = 0.02$ $C_1 = 800 \text{ nF} \Rightarrow 400 \text{ nF}$



VC1 output



error on VC1 output



The INPUT of this test case is a trivial **HEAVISIDE**

Therefore, there is **no contribution** of the "**ILT package**" INPUT **intrinsic sampling error method** to this filter **output response**

This is a favorable situation to measure on its own the percussion processing quality



SPICE is a *matrix based solver*

computing the behavior of some (*linear* and/or *non-linear*) *global circuit* by the *full evaluation* of its *possibly very (very) large* number of individual *NODES*

INTEGRATION (and *NON-LINEARITY*) are handled through “*iterative*” *heuristics*

for *Linear Time Invariant (LTI)* circuit,

TRANSFER FUNCTION models the behavior of a *single observed signal*

The *DENOMINATOR* reflects the complexity of the *circuit STATE “as a whole”*

Therefore, to limit the *polynomial order* this method is reserved to *small circuit*

Nevertheless, the advantage of a *TRANSFER FUNCTION* is that the associated *DIFFERENTIAL EQUATION* allows the manipulation of the overall *circuit STATE* through “*algorithmic equations*”

This is why this facilitates the modeling of **PERCUSSIONS**

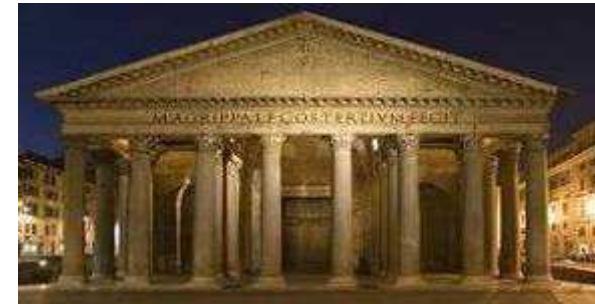
PERCUSSIONS also occur in a “pinball”

By analogy with a **Differential Equation**, the metal ball has a **trajectory**

An **instantaneous shock** on an **obstacle** is enough to change it **forever . . .**



HALL of Fame



| | | |
|-----------------------------------|---|-------------------|
| | Brook Taylor | 1685 - 1731 |
| | Marquis Pierre-Simon Laplace | 1749 - 1827 |
| | Joseph Fourier | 1768 - 1830 |
| | Gustav Robert Kirchhoff | 1824 - 1887 |
| | Oliver Heaviside | 1850 - 1925 |
| | Carl Runge & Martin Wilhelm Kutta method | 1901 |
| | Paul Adrien Maurice Dirac | 1902 - 1984 |
| | Laurent Schwartz | 1915 - 2002 |
| <i>MIT Macsyma (now Maxima)</i> | | <i>late 1960s</i> |
| <i>SPICE 1 public domain</i> | | 1972 |

QUESTIONS ?

